

*Correlation between electronic
structure, magnetism and physical
properties of Fe-Cr alloys: ab initio
modeling*

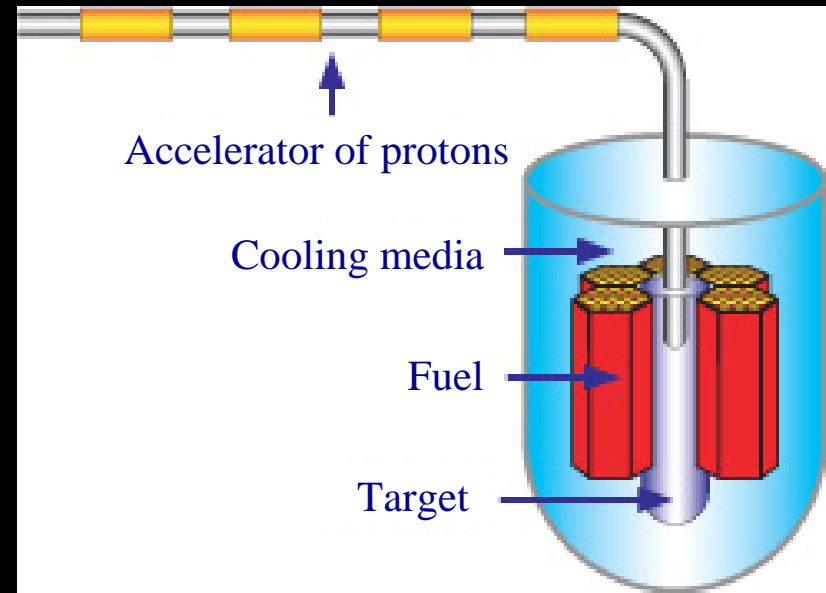
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and Biology (IFM),**

Linköping University, Sweden

Fe-Cr alloys

- Are the base for many important industrial steels
- Used as cladding material in fast neutron reactors
- Low Cr steels, up to 10 % Cr, show:
 - anomalous stability
 - resistance to neutron radiation induced swelling
 - corrosion resistance
 - increased ductile to brittle transition temperature



In collaboration with:



P. Olsson, EDF R&D, France

A. V. Ponomareva, MIS&A, Russia

A. V. Ruban, KTH, Sweden

L. Vitos, KTH, Sweden

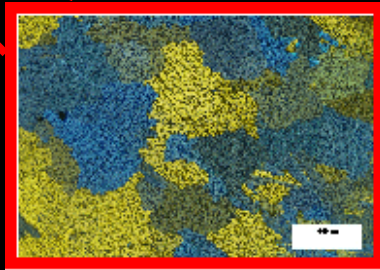
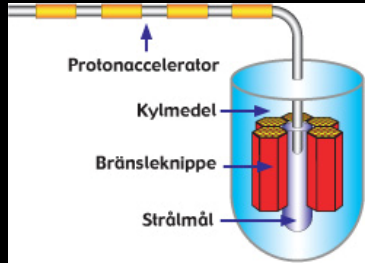
J. Wallenius, KTH, Sweden

CONTENTS :



- Phase equilibria in solid solutions: theoretical background
- First-principles methods: recent developments
- Fe-Cr alloys: first-principles results
- Discussion: understanding of first-principles results
- Conclusions

Microstructure – Property Relationships



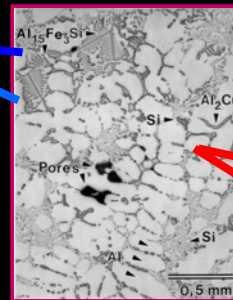
Microstructure

- Grains

≅ 1 – 10 mm

Properties

- High cycle fatigue
- Ductility



Microstructure

- Phases

≅ 100 – 500 microns

Properties

- Yield strength
- Ultimate tensile strength
- High cycle fatigue
- Low cycle fatigue
- Thermal Growth
- Ductility



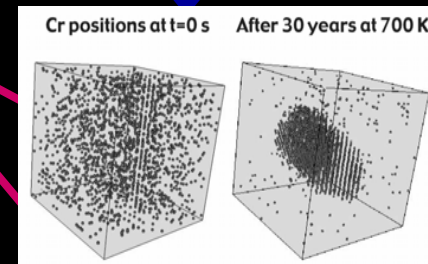
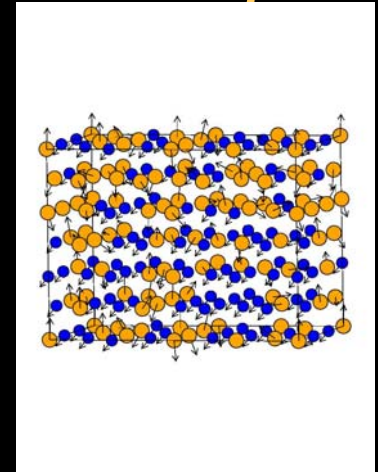
Microstructure

- Phases

≅ 3-100 nanometers

Properties

- Yield strength
- Ultimate tensile strength
- Low cycle fatigue
- Ductility



Atoms

≅ 10-100 Angstroms

Properties

- Thermal Growth

Original idea for this figure belongs to Chris Wolverton Ford Motor Company

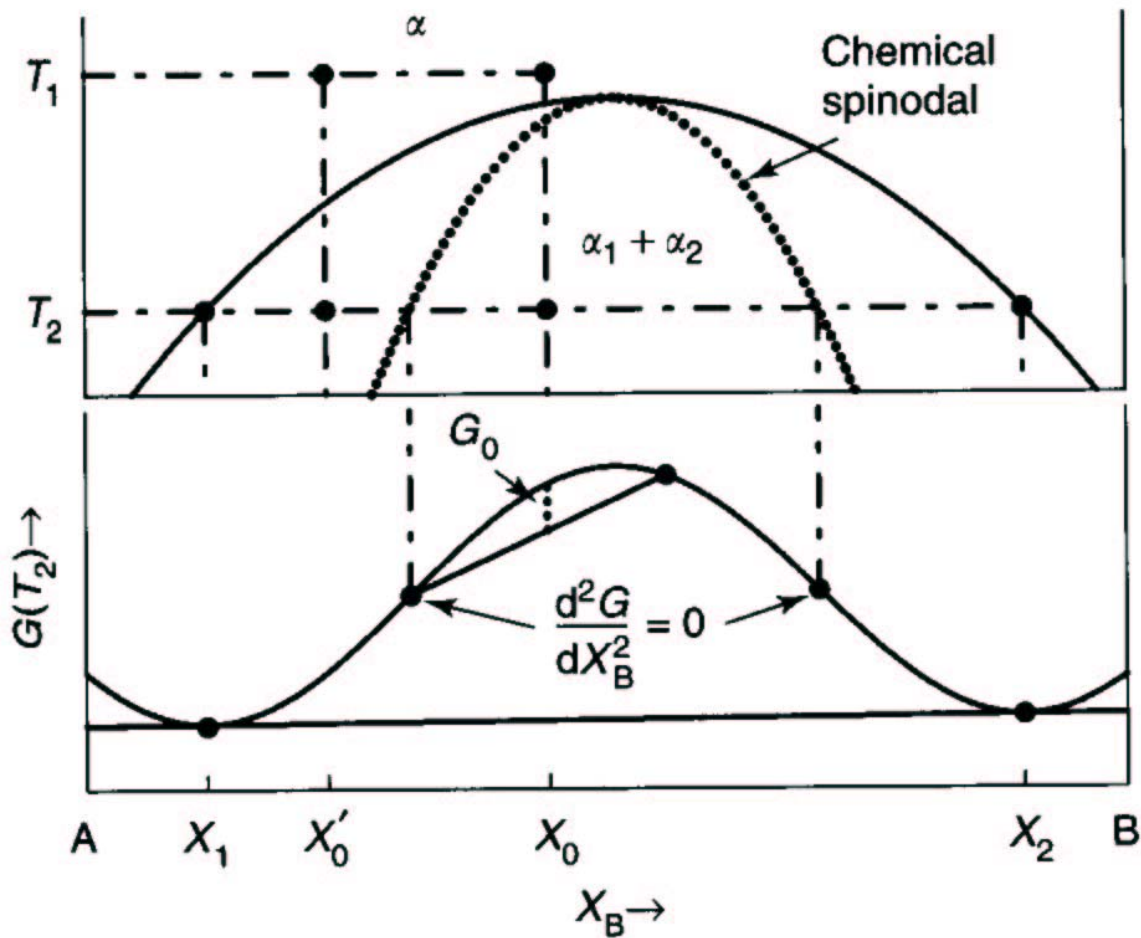
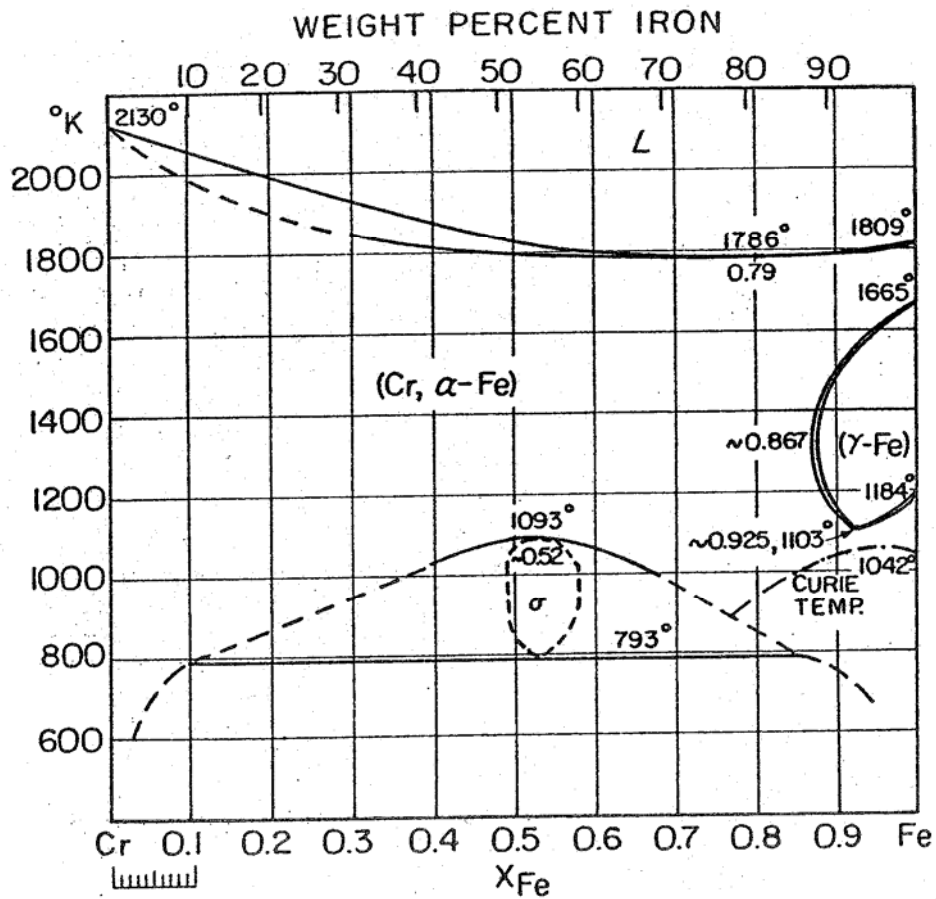
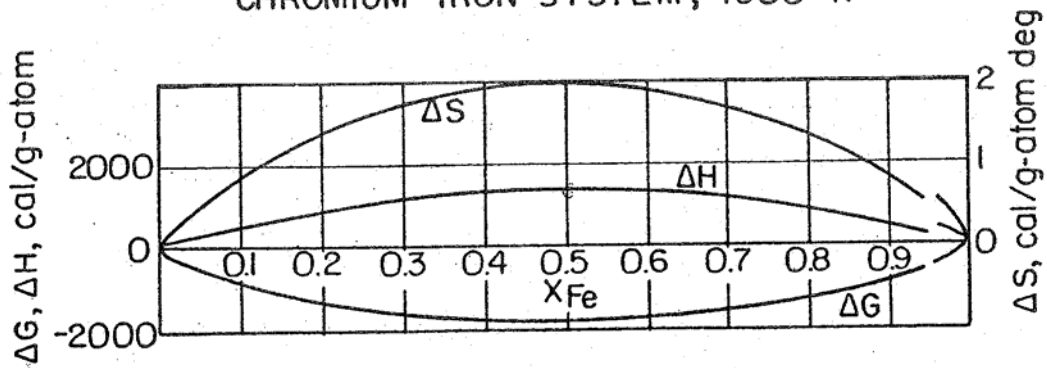


Figure 9 (a) Schematic phase diagram and (b) free energy vs. composition diagram for alloys between the spinodal points, which are unstable and can decompose into two coherent phases α_1 and α_2 without overcoming an activation energy barrier. Alloys between the miscibility gap and the spinodal are metastable and can decompose only after nucleation of the other phase.



CHROMIUM - IRON SYSTEM, 1600 °K

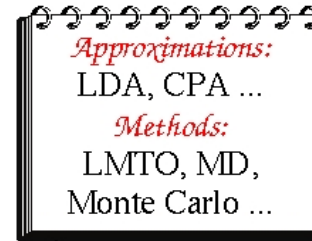
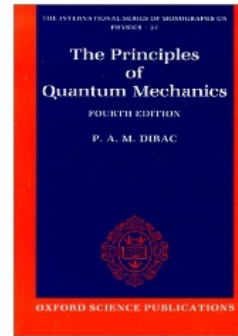


First-principles calculations

Periodic Table

Lanthanides: La, Ce, Pr, Nd, Pm, Sm, Eu, Gd, Tb, Dy, Ho, Er, Tm, Yb, Lu

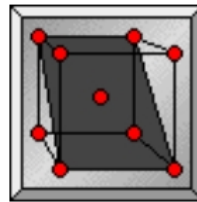
Actinides: Ac, Th, Pa, U, Np, Pu, Am, Cm, Bk, Cf, Es, Fm, Md, No, Lr



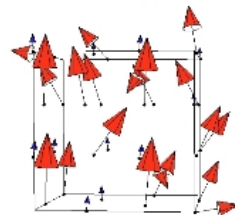
Experimental data



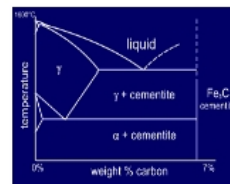
Adjustable parameters



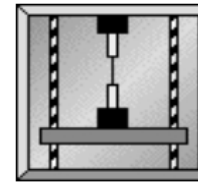
Structure



Magnetism



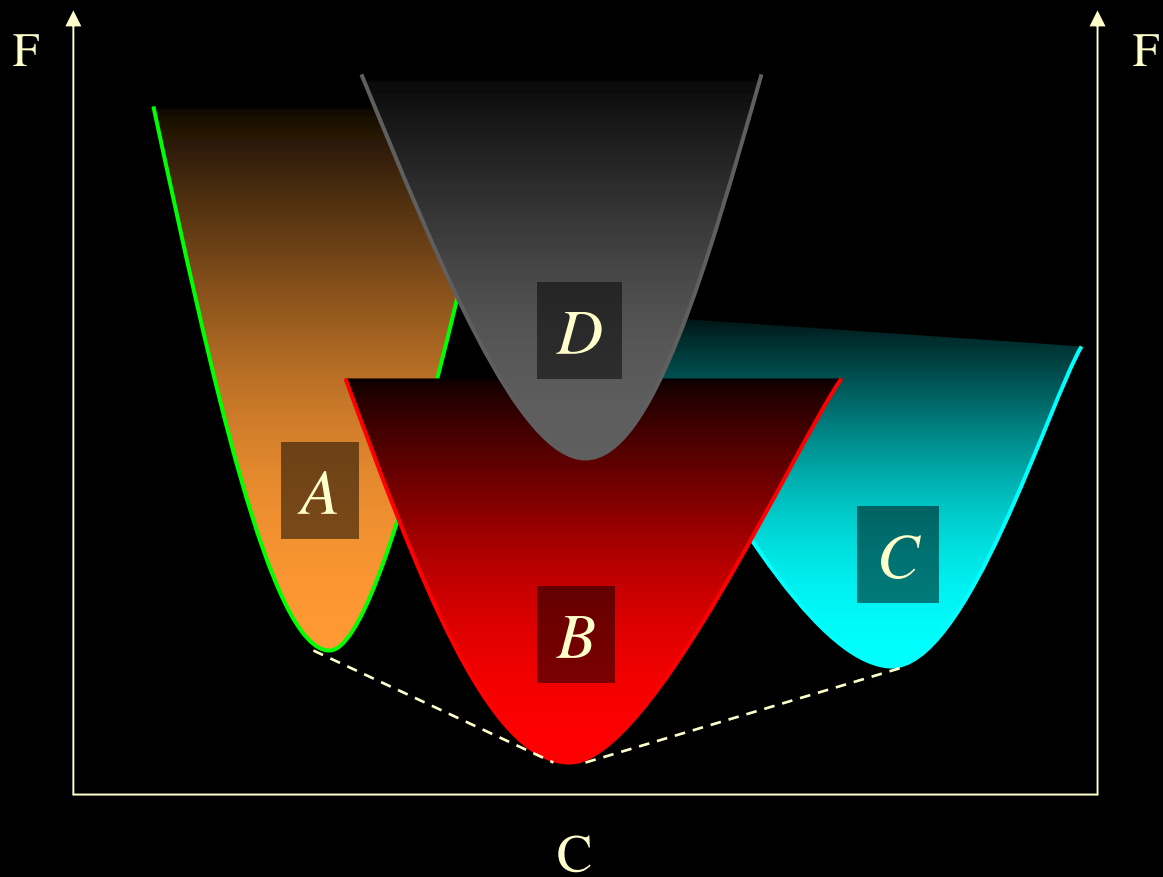
Phase equilibria



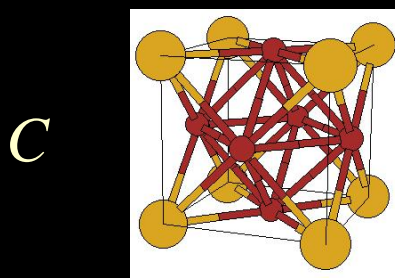
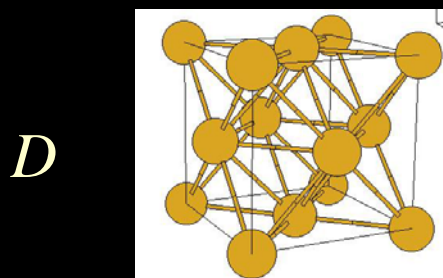
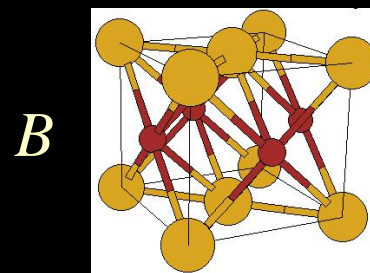
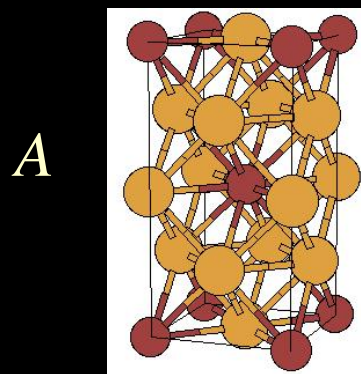
Mechanical properties



Fundamental understanding



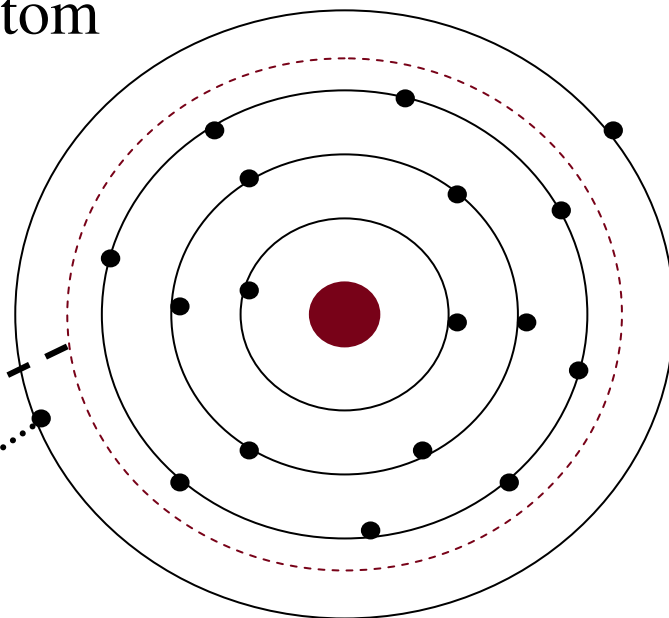
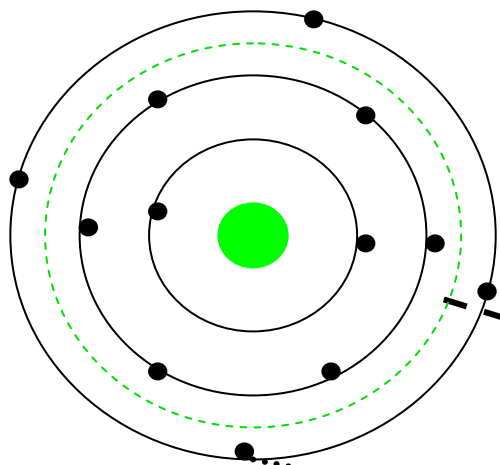
Structures:



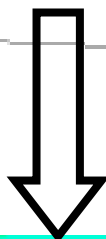
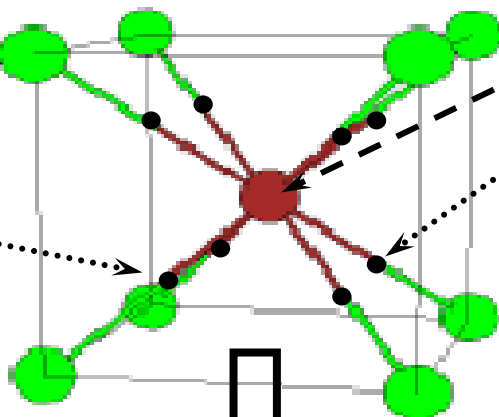
$$\underline{F = -k_B T \ln Z} \quad \underline{Z = \sum_s \exp\left(-\frac{E_s}{kT}\right)}$$

A atom

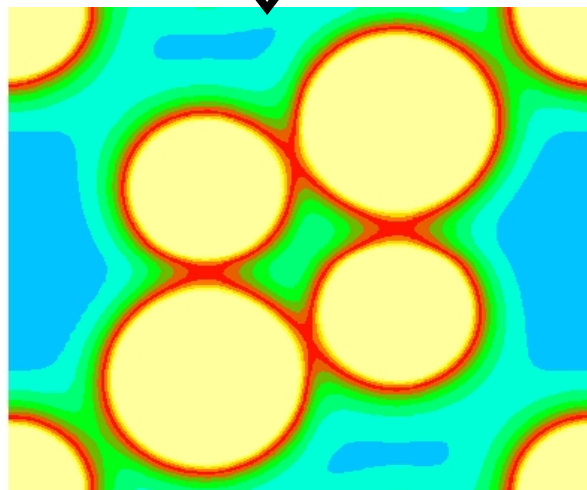
B atom



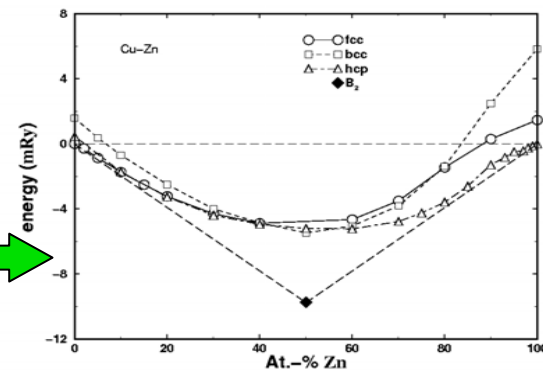
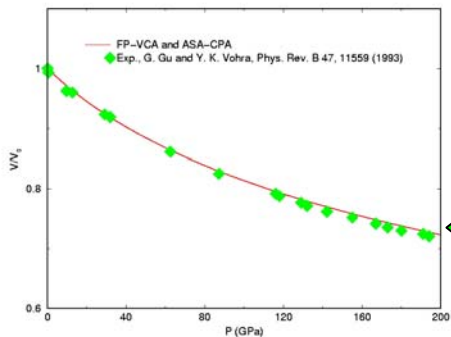
$A_x B_{1-x}$ alloy

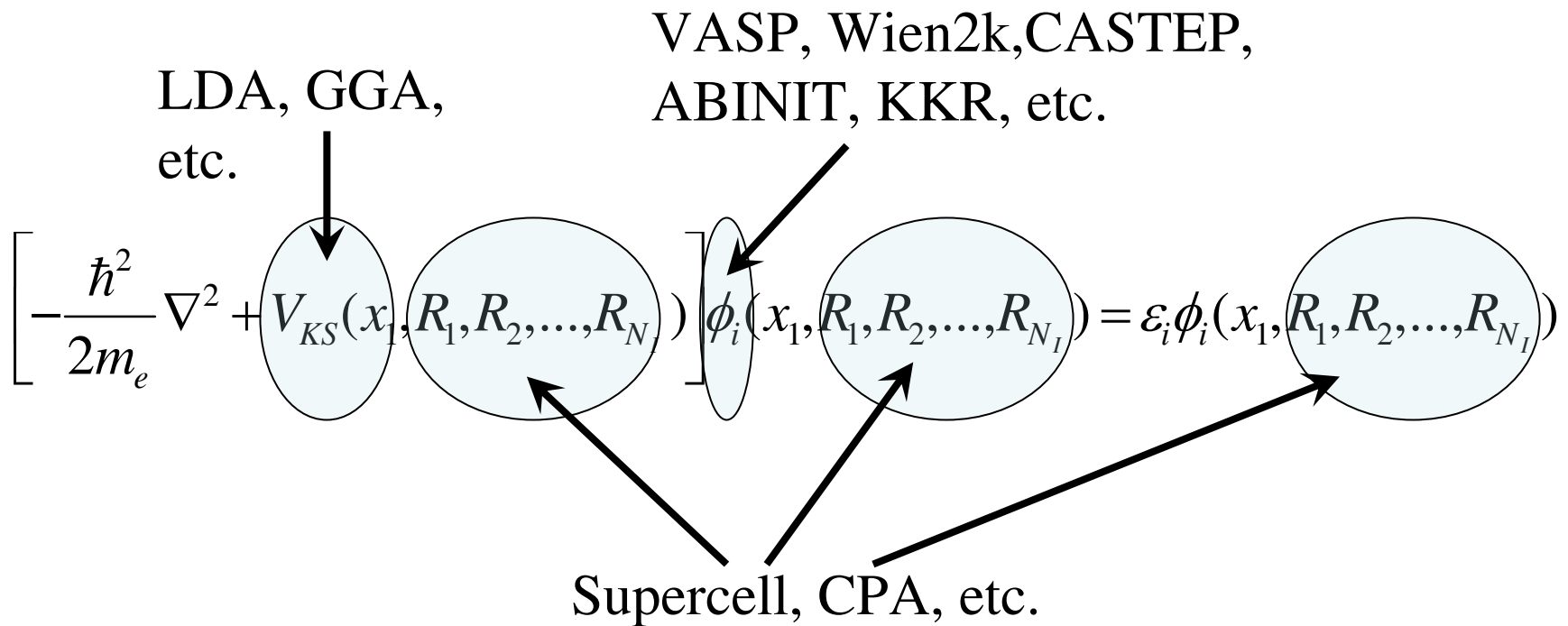


$n(\mathbf{r})$, the electron density



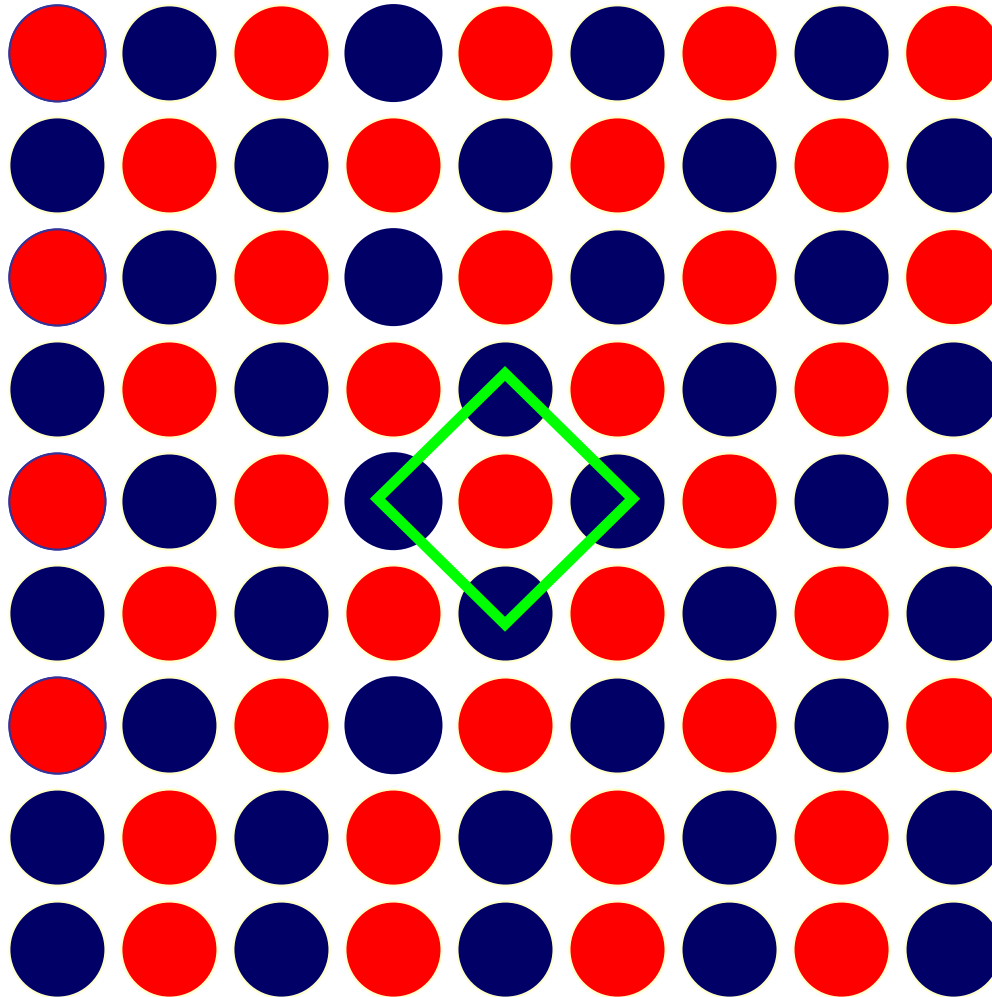
The EOS for bcc $Mo_{69}Re_{32}$ random alloy

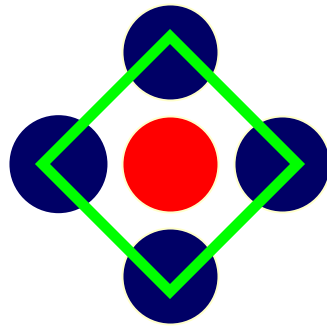




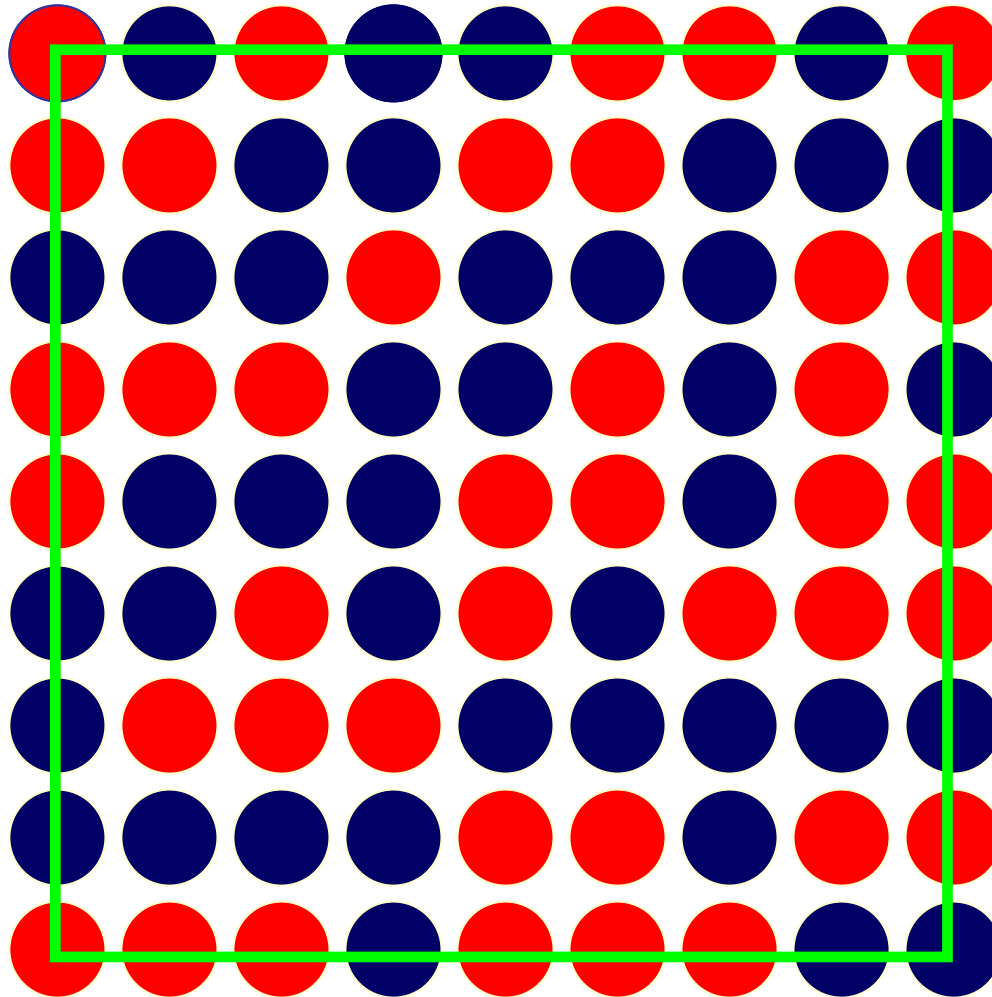
**It is NOT a trivial task
 to run *ab initio* software!**

Ordered compounds

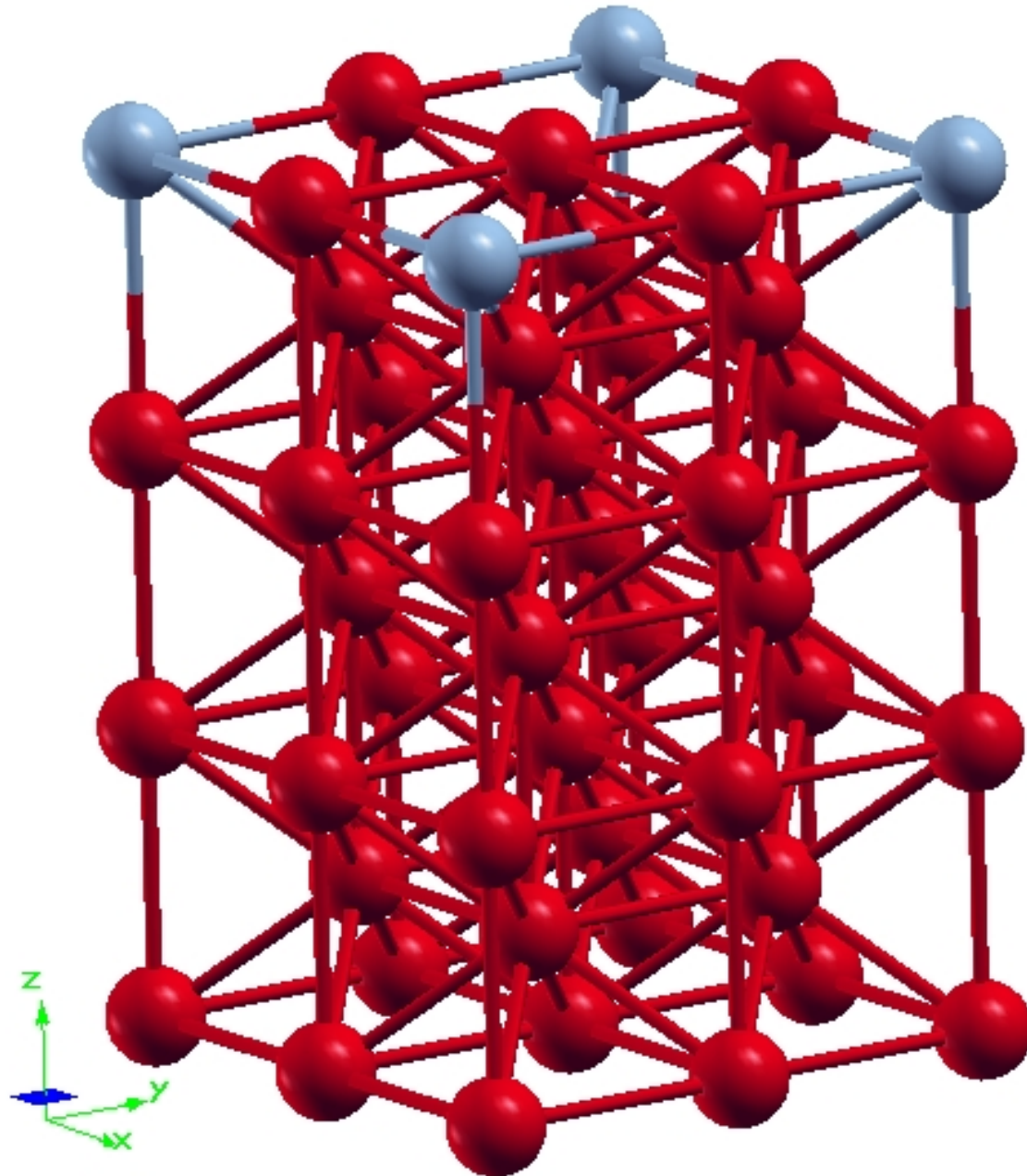




Solution phases: supercell method

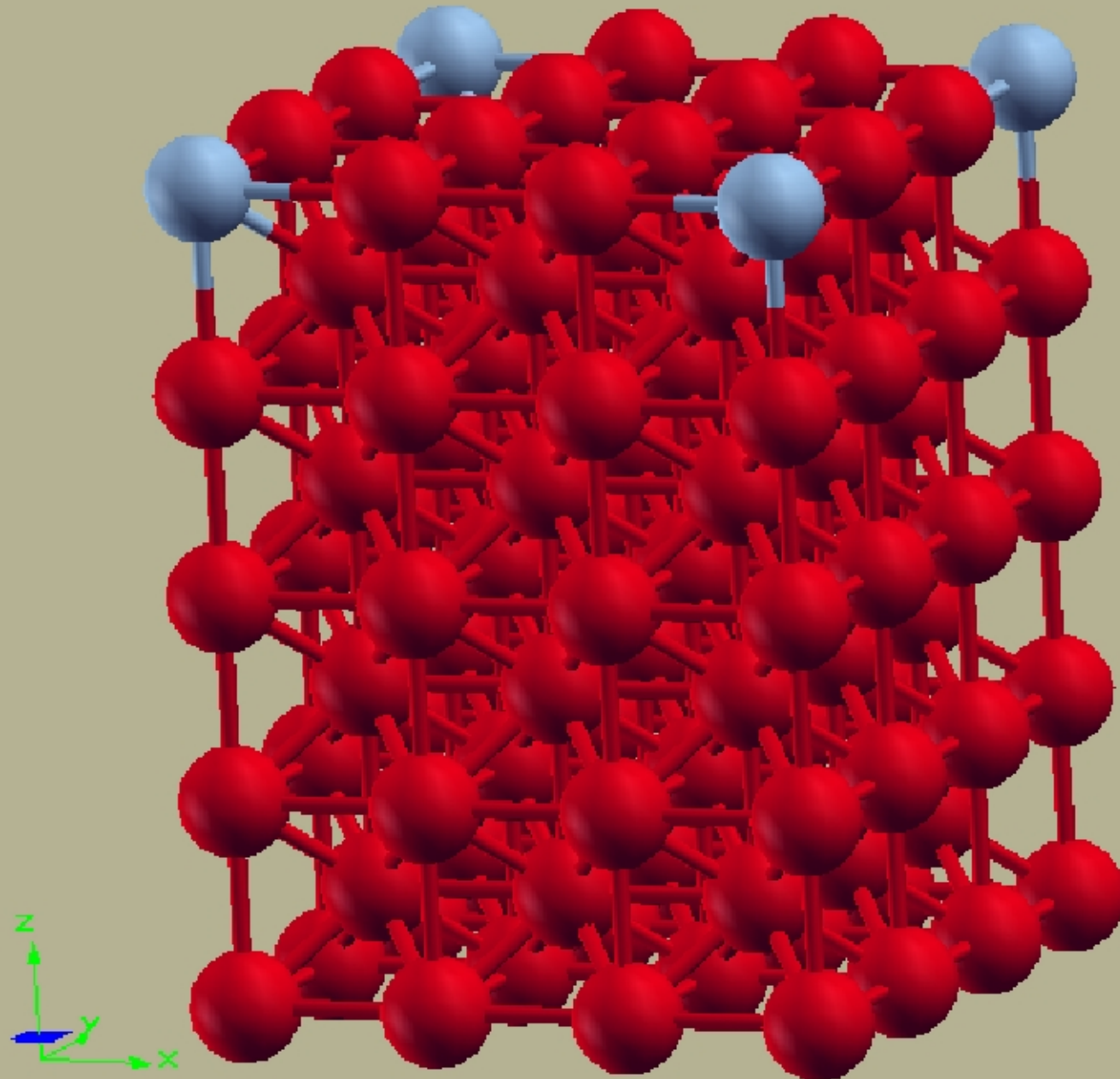


2X2X7, segregation energy -0.043 eV



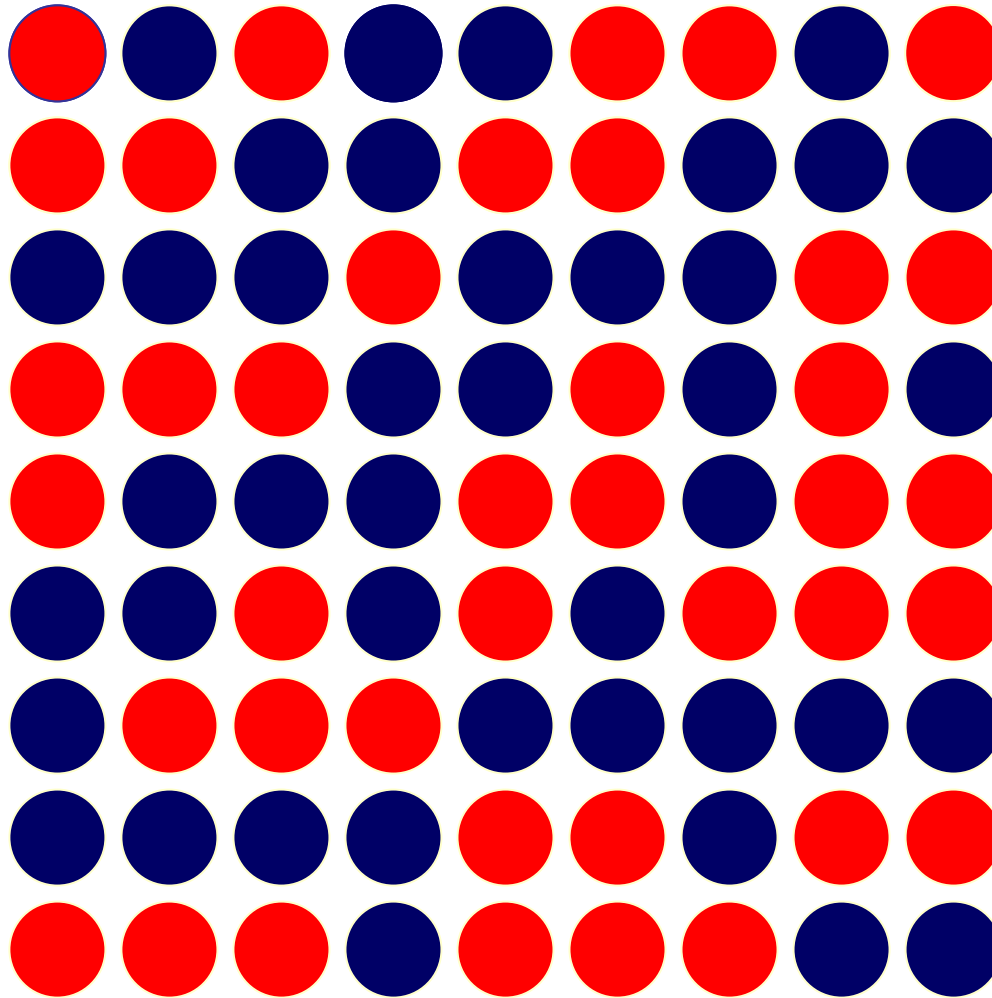
A. V. Ponomareva *et al.*, Phys. Rev. B 75, 245406 (2007)

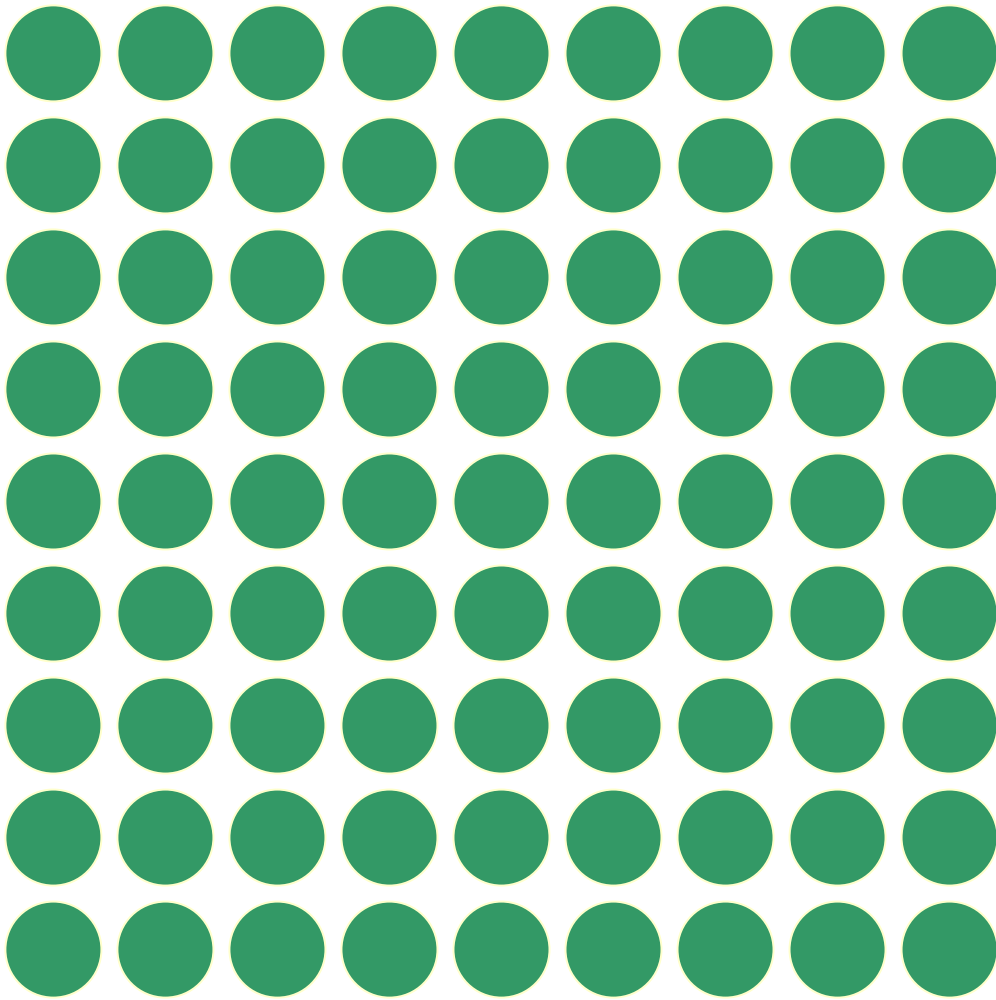
3X3X9, segregation energy 0.090 eV

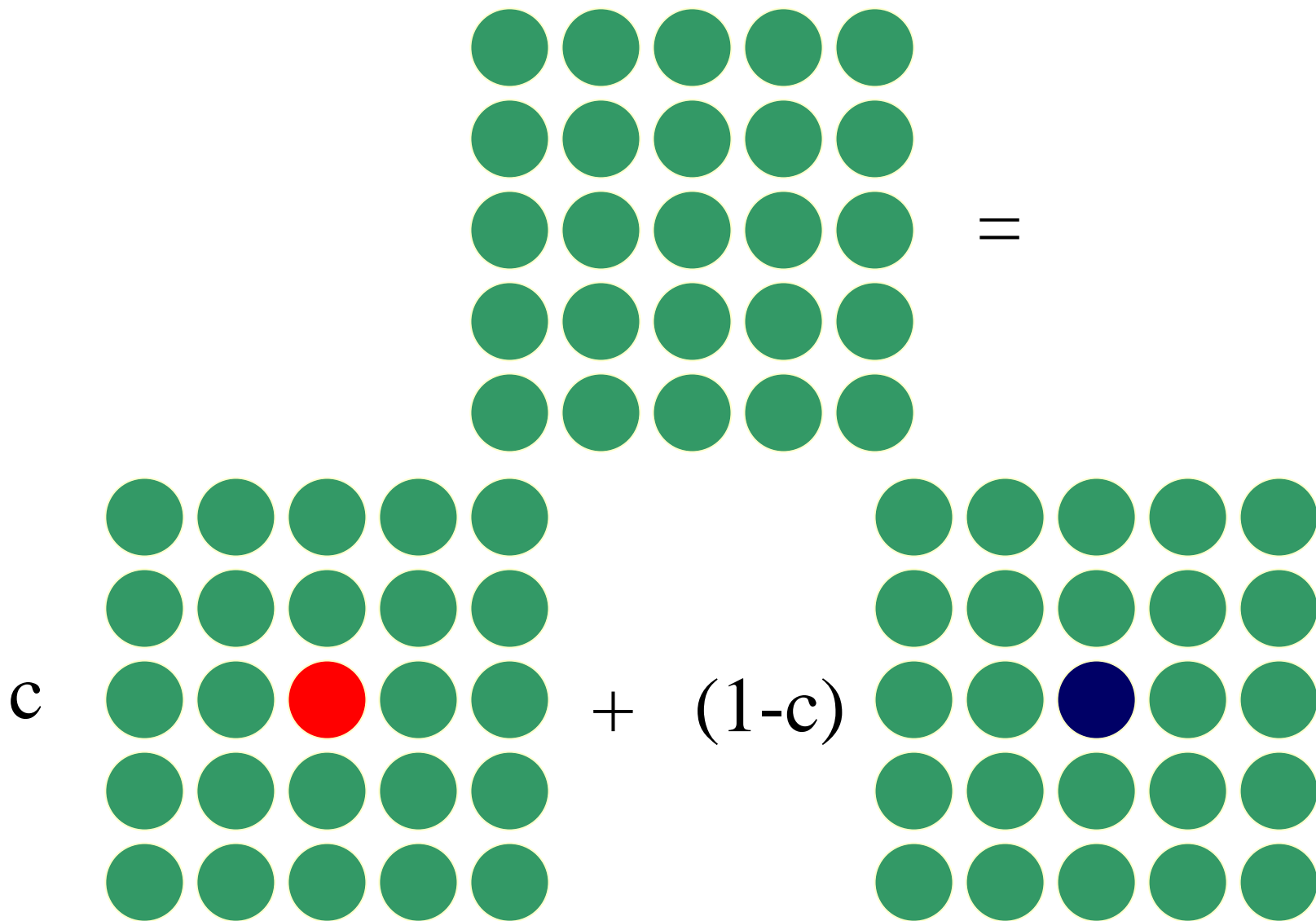


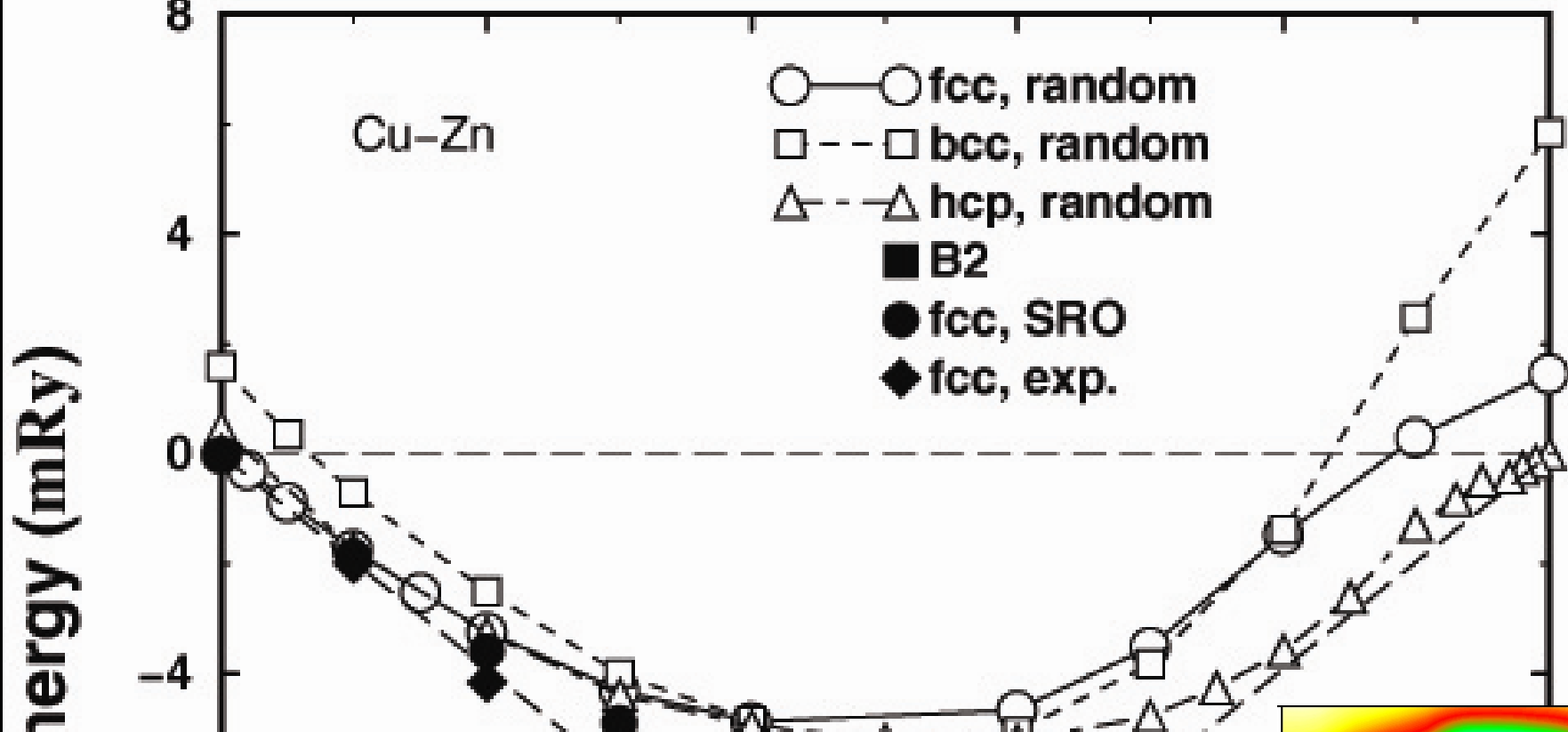
A. V. Ponomareva *et al.*, Phys. Rev. B 75, 245406 (2007)

Solution phases: coherent potential

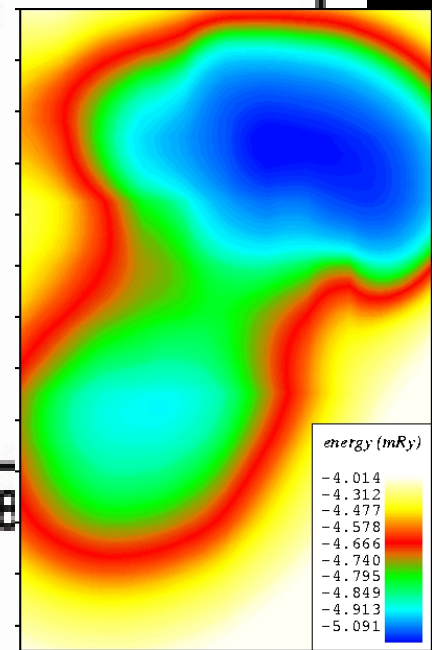
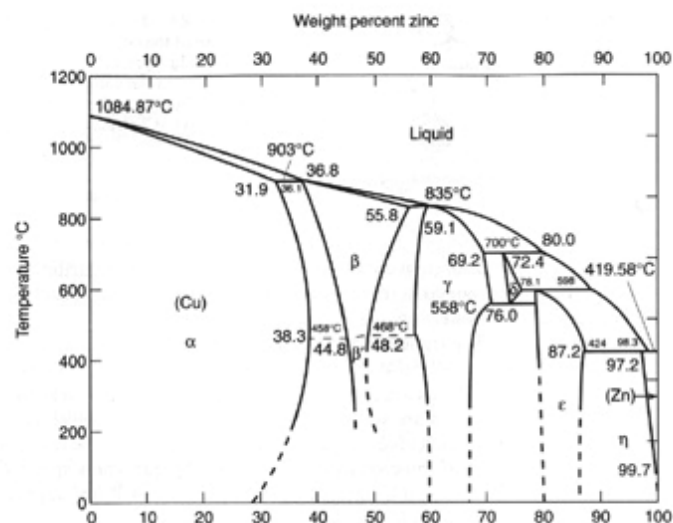


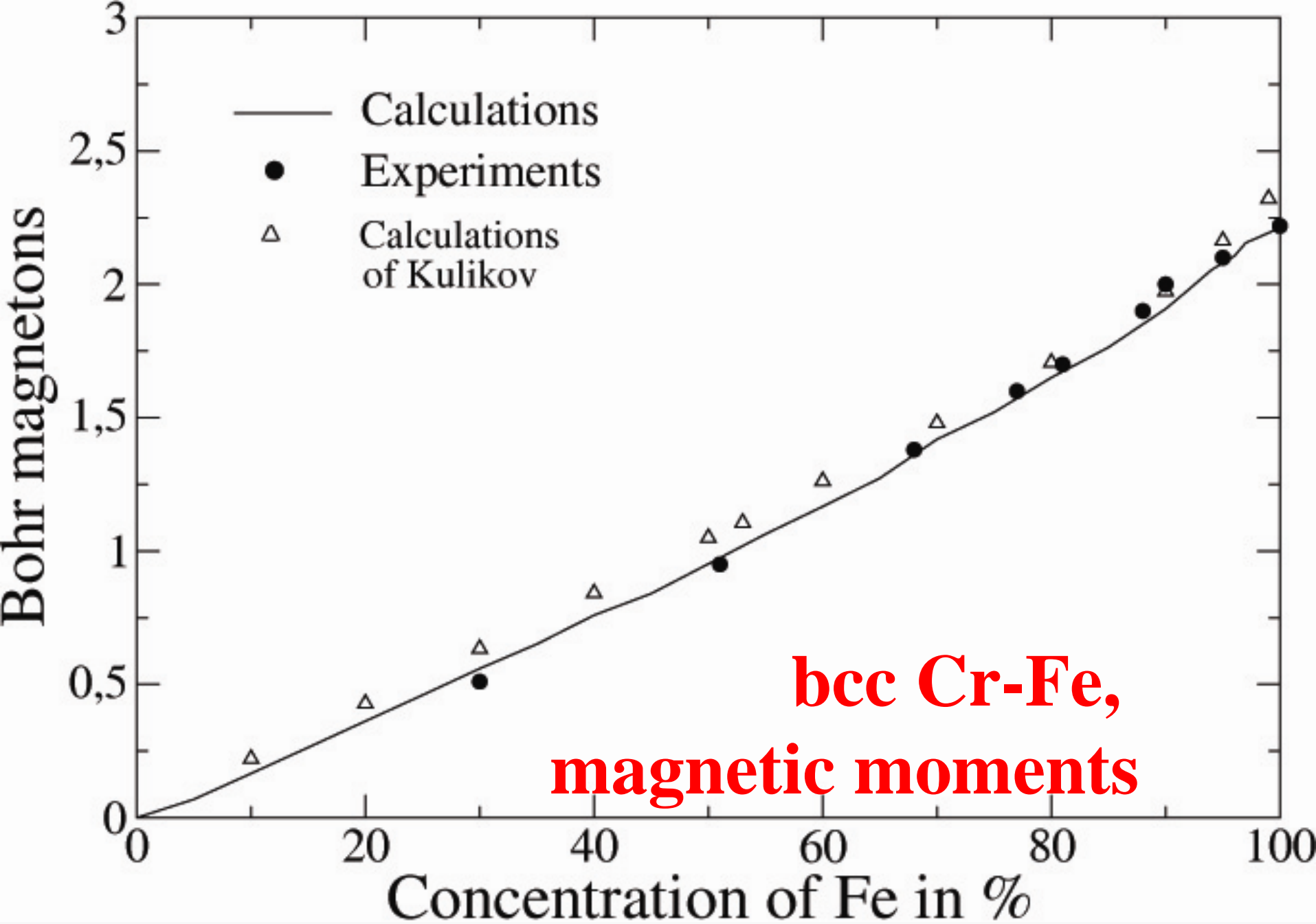


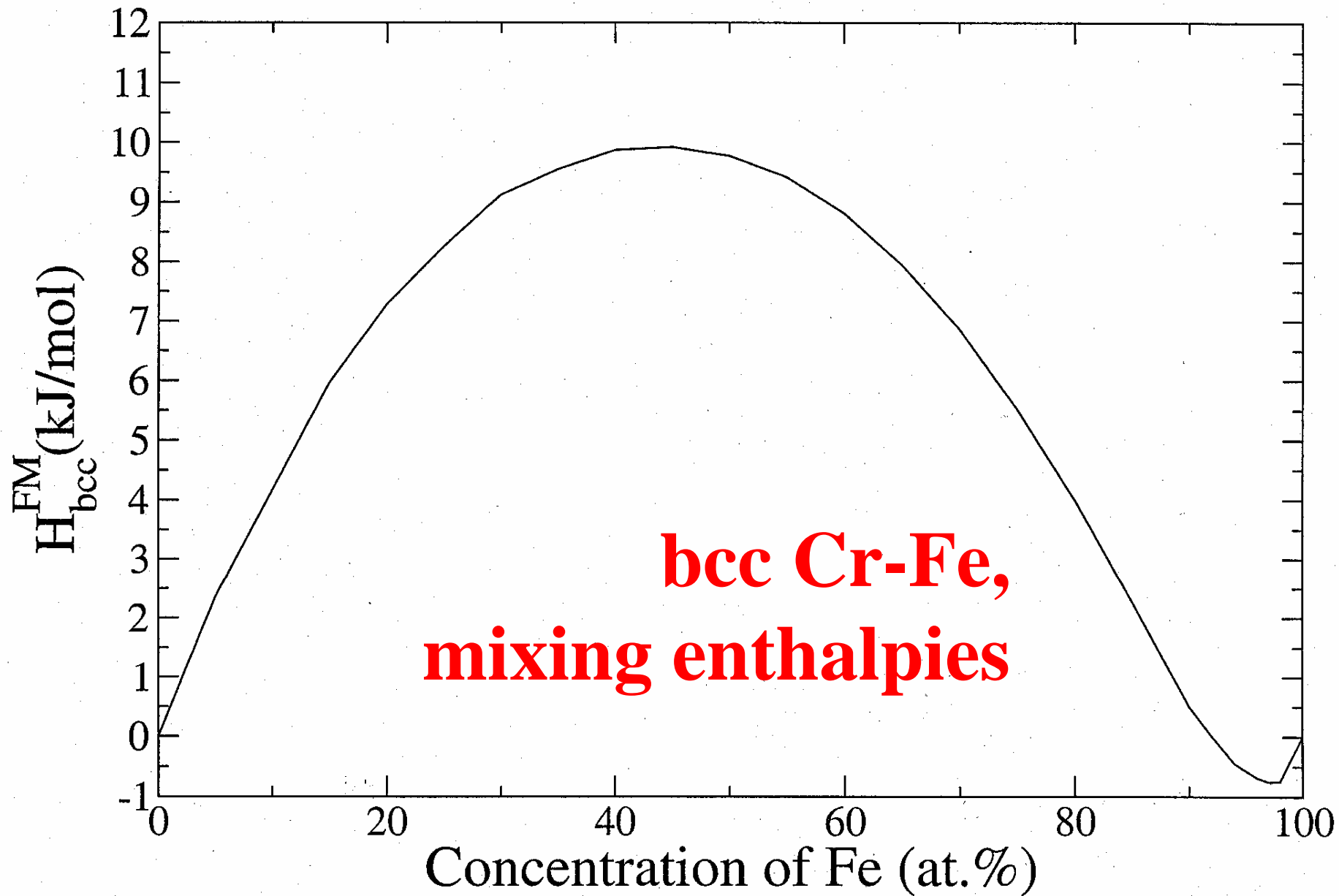




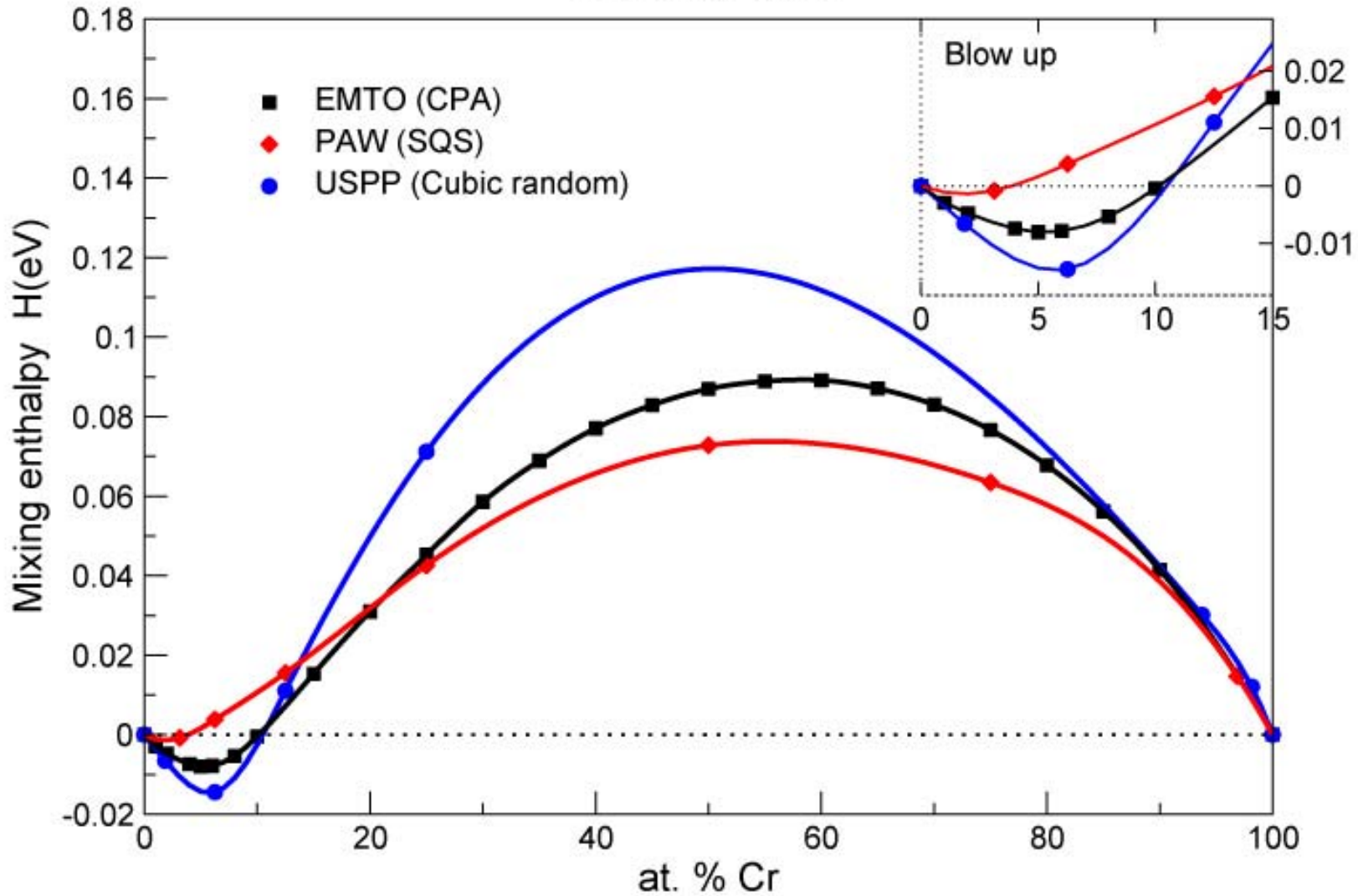
The Copper-Zinc Phase Diagram







FM bcc Fe-Cr



FM Fe-Cr bcc

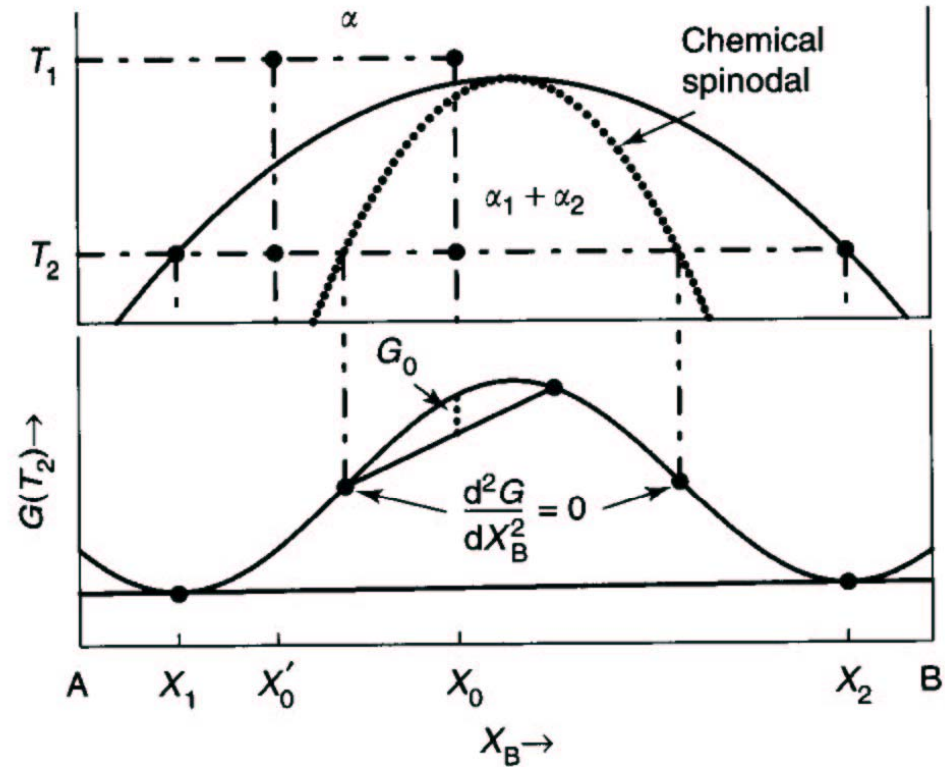
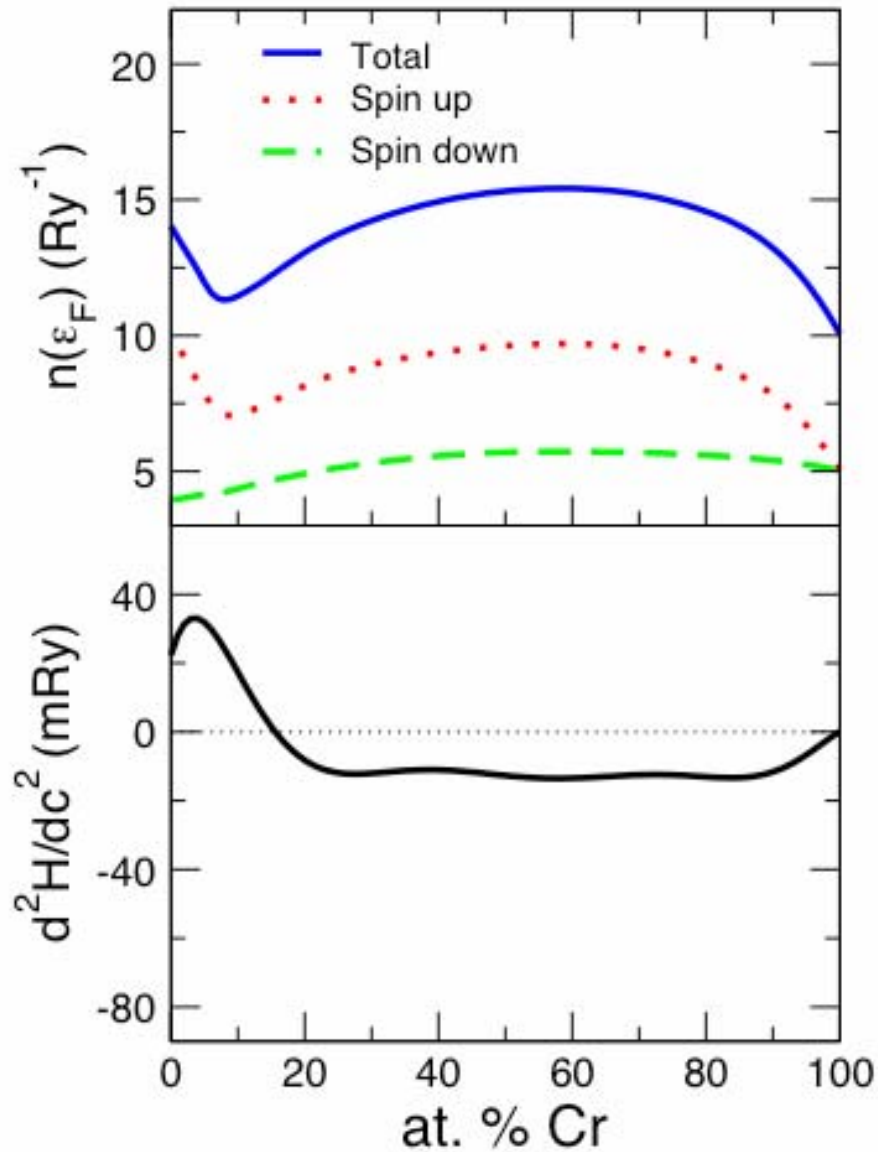
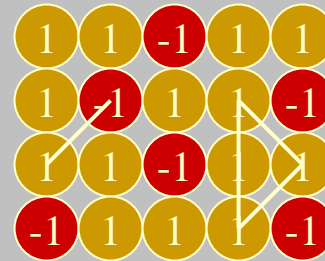


Figure 9 (a) Schematic phase diagram and (b) free energy vs composition diagram for alloys between the spinodal points which are unstable and can decompose into two coherent phases α_1 and α_2 without overcoming an activation energy barrier. Alloys between the miscibility gap and the spinodal are metastable and can decompose only after nucleation of the other phase.

$$\underline{F = -k_B T \ln Z} \quad \underline{Z = \sum_s \exp\left(-\frac{E_s}{kT}\right)}$$

$$\vec{\sigma}_s = \{\sigma_i\} \quad \sigma_i = \begin{cases} 1 & \text{if site } i \text{ is occupied by atom A} \\ -1 & \text{otherwise} \end{cases}$$

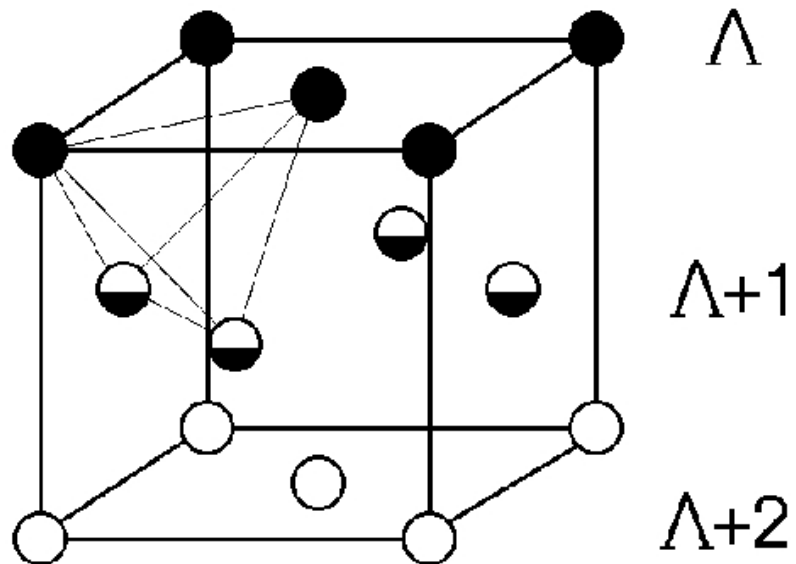
$$\langle \sigma_i \sigma_j \rangle^s, \quad \langle \sigma_i \sigma_j \sigma_k \rangle^s, \dots$$



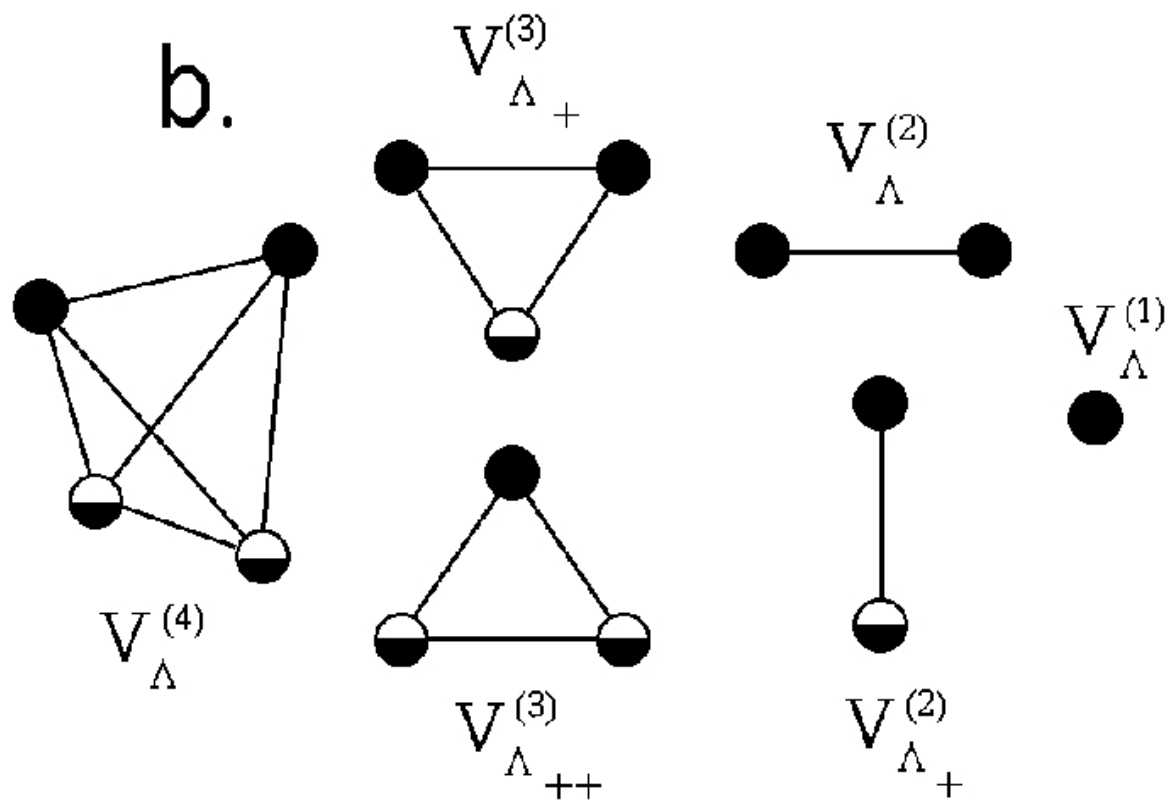
$$\underline{E_{tot}} = \underline{V^{(0)}} + \underline{V^{(1)}} \langle \sigma \rangle +$$

$$\sum_s \underline{V^{(2,s)}} \langle \sigma_i \sigma_j \rangle + \sum_s \underline{V^{(3,s)}} \langle \sigma_i \sigma_j \sigma_k \rangle + \dots$$

a.



b.



Calculations of effective interatomic potentials

The generalized perturbation method

1. Calculate electronic structure of a random alloy (for example, use the CPA):

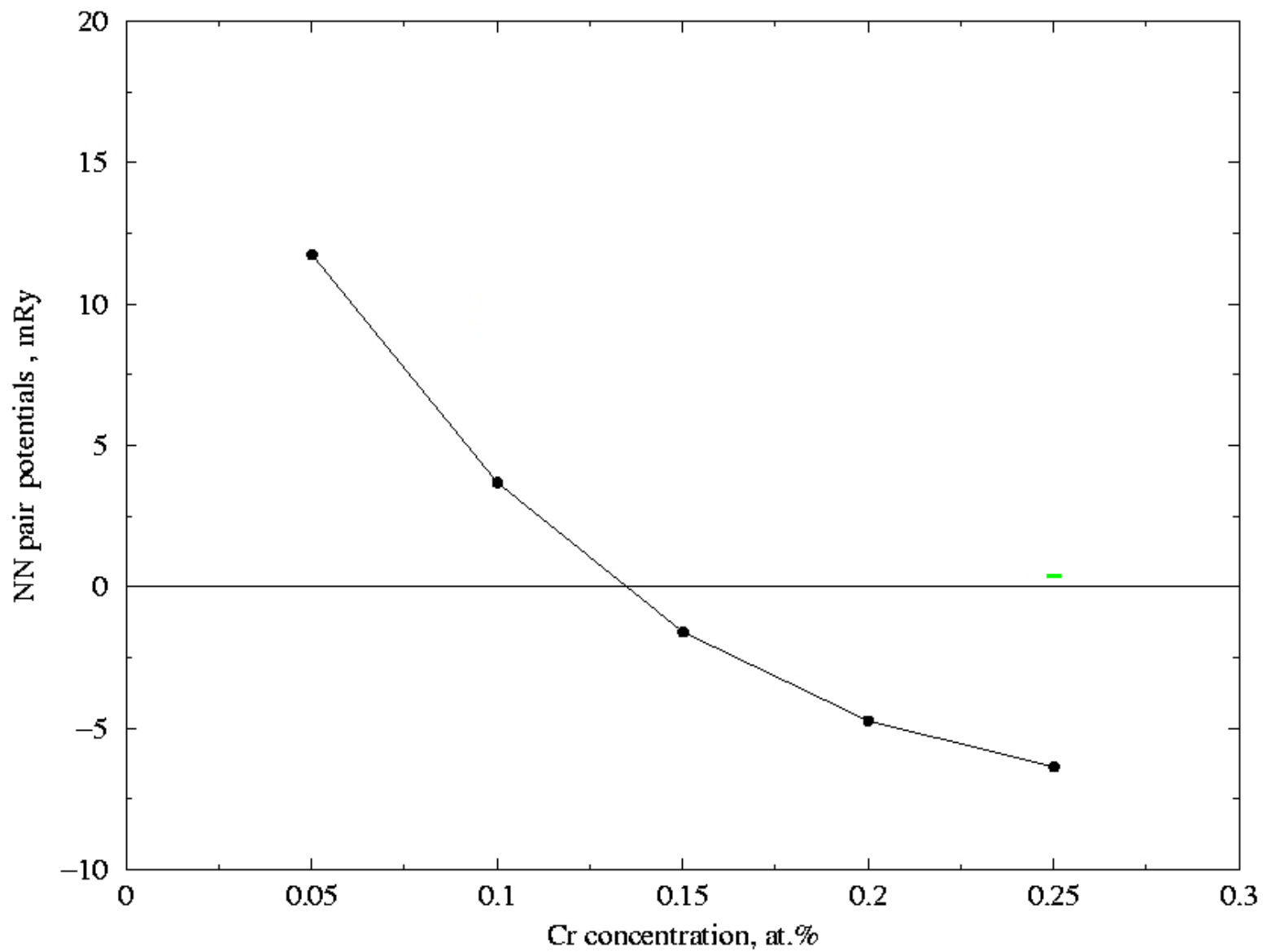
$$\underline{\tilde{g}, t^{A(B)}}$$

2.
$$E_{one} = E_{one}(c) + \frac{1}{2} \sum_{RR'} V_{RR'} \langle \sigma_i \sigma_j \rangle^{(RR')} \dots$$

-determine a perturbation of the band energy due to small variations of the correlation functions

where the effective interatomic interactions are given by an analytical formula:

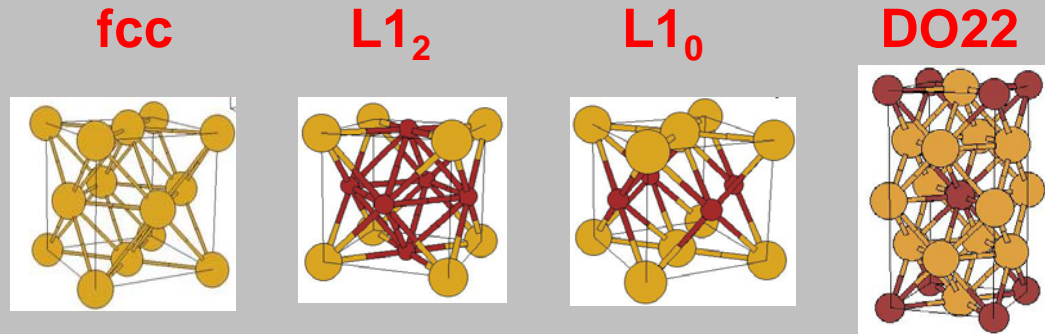
$$V_{RR'} = -\frac{1}{\pi} \int dE \left\{ t_R^A \gamma_{RR'} t_{R'}^B \gamma_{R'R} t_R^A \right\}$$



Calculations of effective interatomic potentials

The Connolly-Williams method

1. Choose structures



with predefined
correlation functions

$$\left[\langle \sigma_i \sigma_j \rangle^{(2,1)}, \langle \sigma_i \sigma_j \rangle^{(2,2)}, \langle \sigma_i \sigma_j \sigma_k \rangle^{(3,1)}, \dots \right]$$

2. Calculate E_{tot} :

E(fcc)

E(L1₂)

E(L1₀)

E(DO22)

$$E_{tot} = V^{(0)} + V^{(1)} \langle \sigma \rangle + \sum_s V^{(2,s)} \langle \sigma_i \sigma_j \rangle \dots$$

If $N_{pot} \leq N_{str}$ then $\{V\}$ are found by L.S.M:

$$\min_m \left[E_m - \sum_f \langle \dots \rangle^f V_f \right]^2 = \min$$

The Monte Carlo method

Calculations of averages at temperature T :

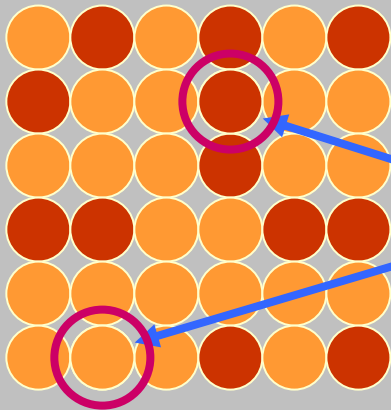
$$\langle A \rangle = \frac{\sum_s A_s \exp\left(-\frac{E_s}{k_B T}\right)}{Z}$$

Create the Markov chain of configurations:

$$P_s = \frac{1}{Z} \exp\left(-\frac{E_s}{k_B T}\right)$$

Balance at the equilibrium state:

$$W(s \rightarrow s') \exp\left(-\frac{E_s}{k_B T}\right) = W(s' \rightarrow s) \exp\left(-\frac{E_{s'}}{k_B T}\right)$$



ΔE

$$\left\{ \begin{array}{l} \Delta E \leq 0 \\ \Delta E > 0 \quad \exp\left[-\frac{\Delta E}{k_B T}\right] > r \quad (0 \leq r \leq 1) \end{array} \right\}$$



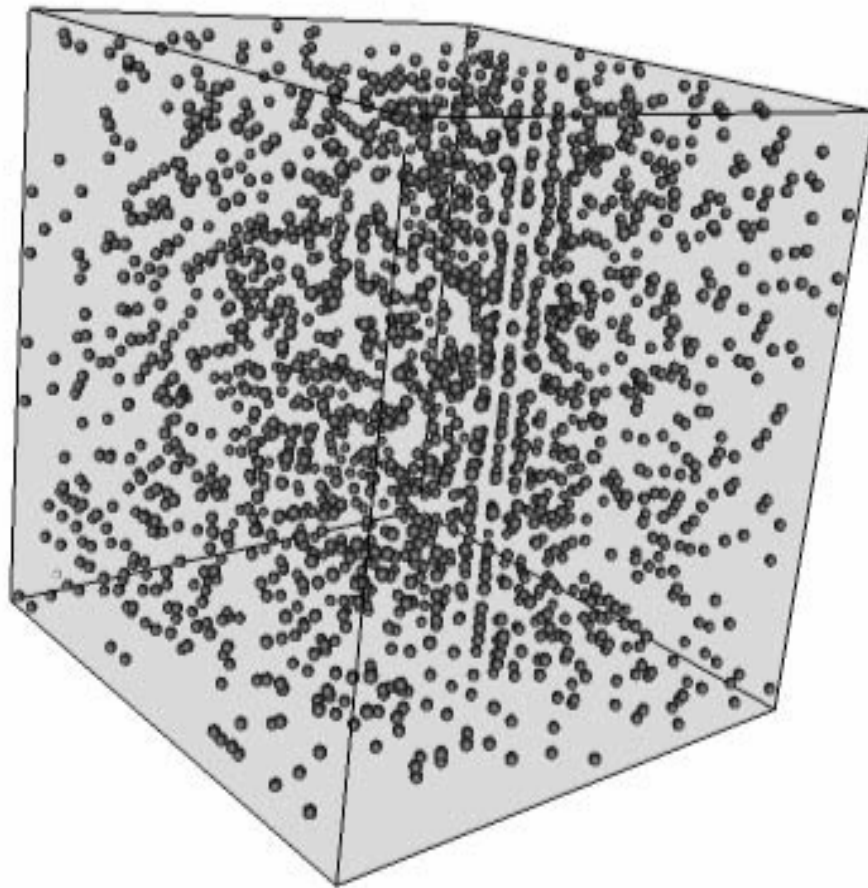
Atoms exchanged

Multiscale simulations of radiation damage in Fe-Cr alloys

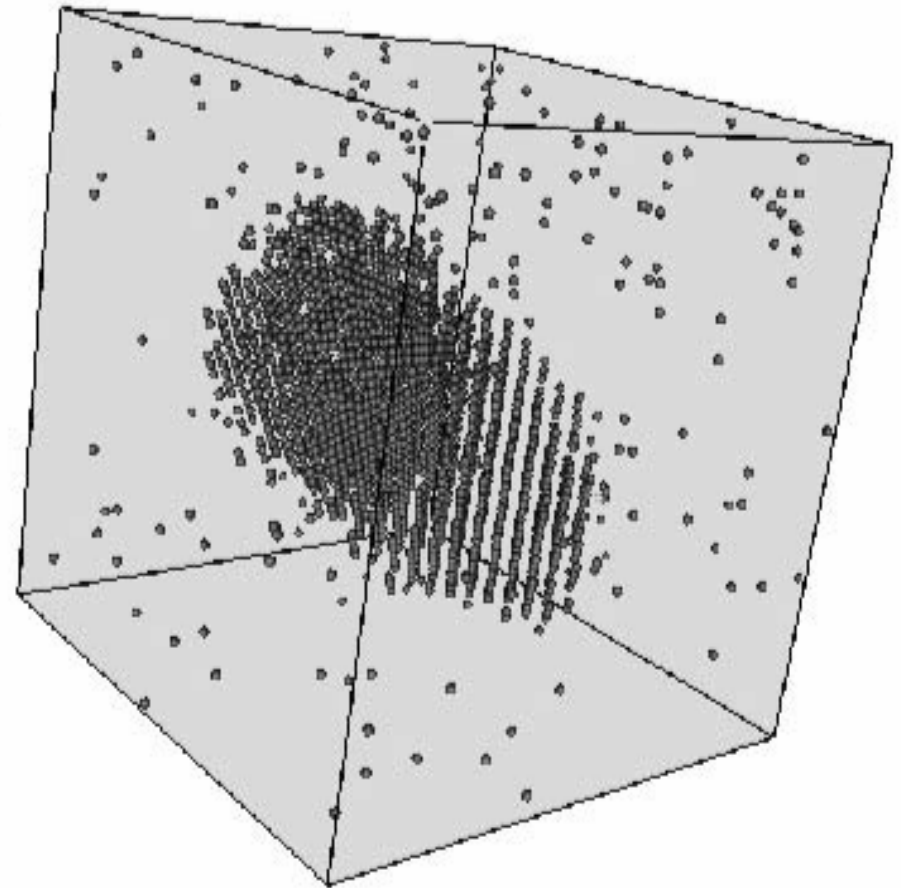
J. Wallenius, I. A. Abrikosov, P. Olsson *et al.*,

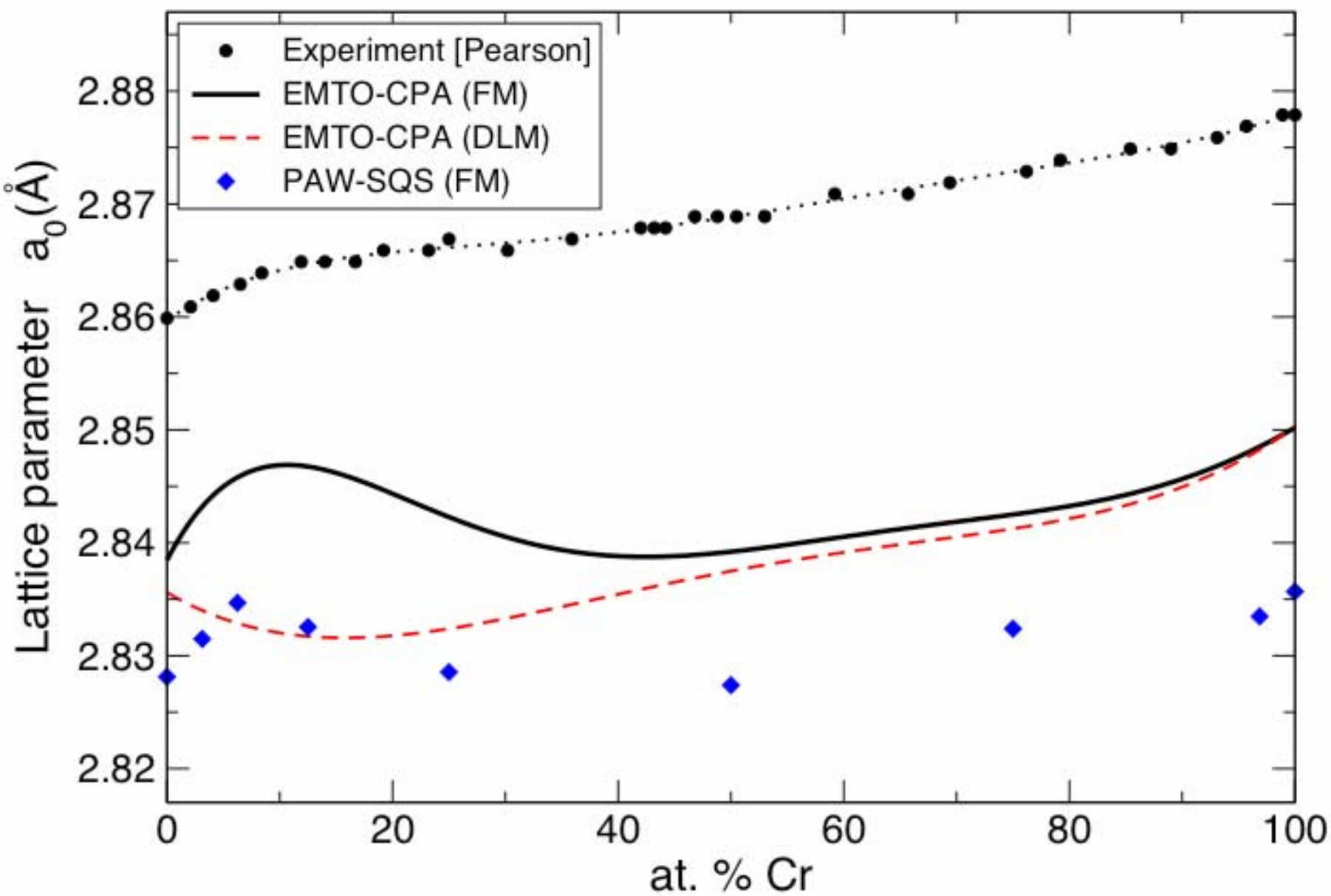
J. Nucl. Mater. **329-333**, 1175 (2004).

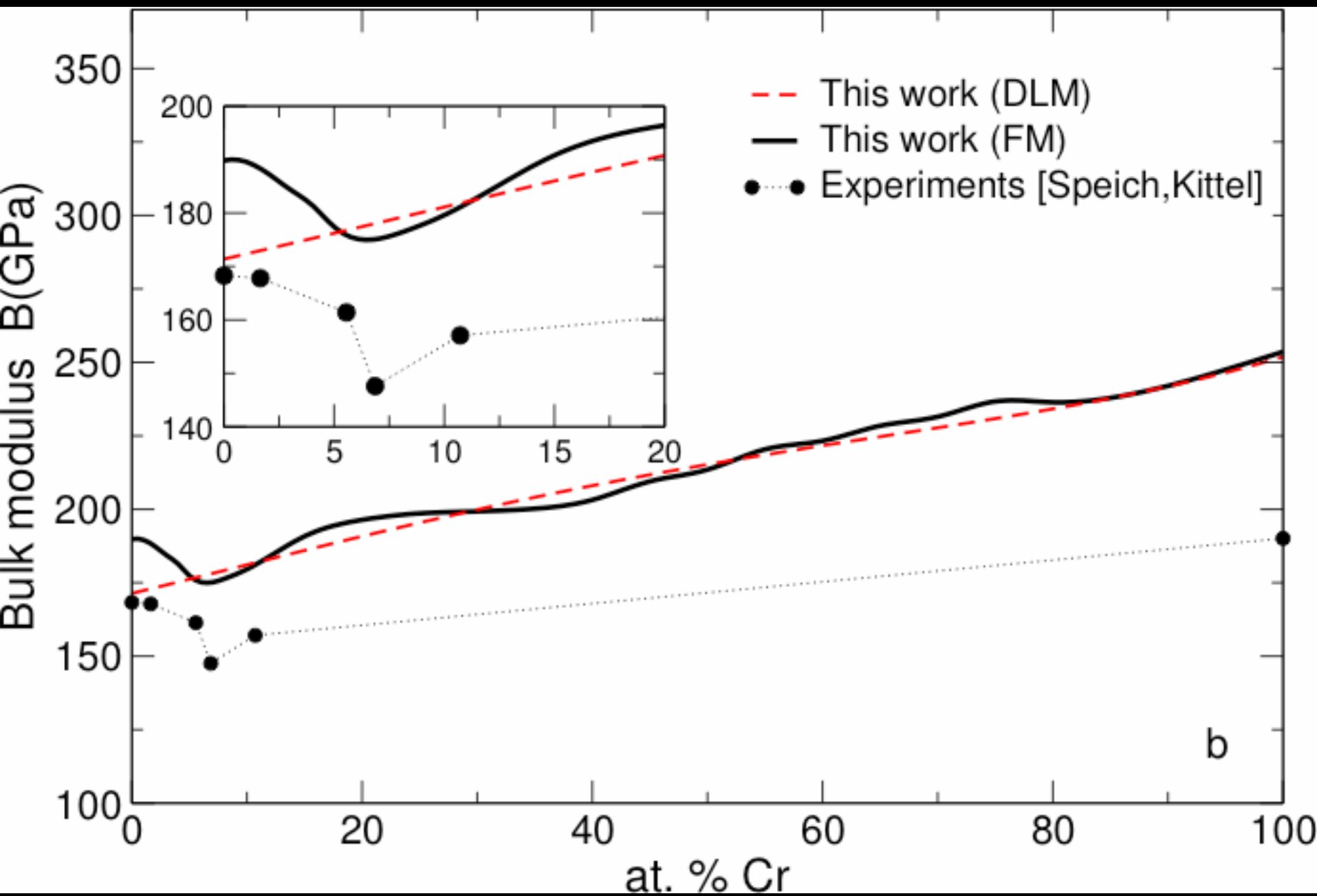
Cr positions at $t=0$ s

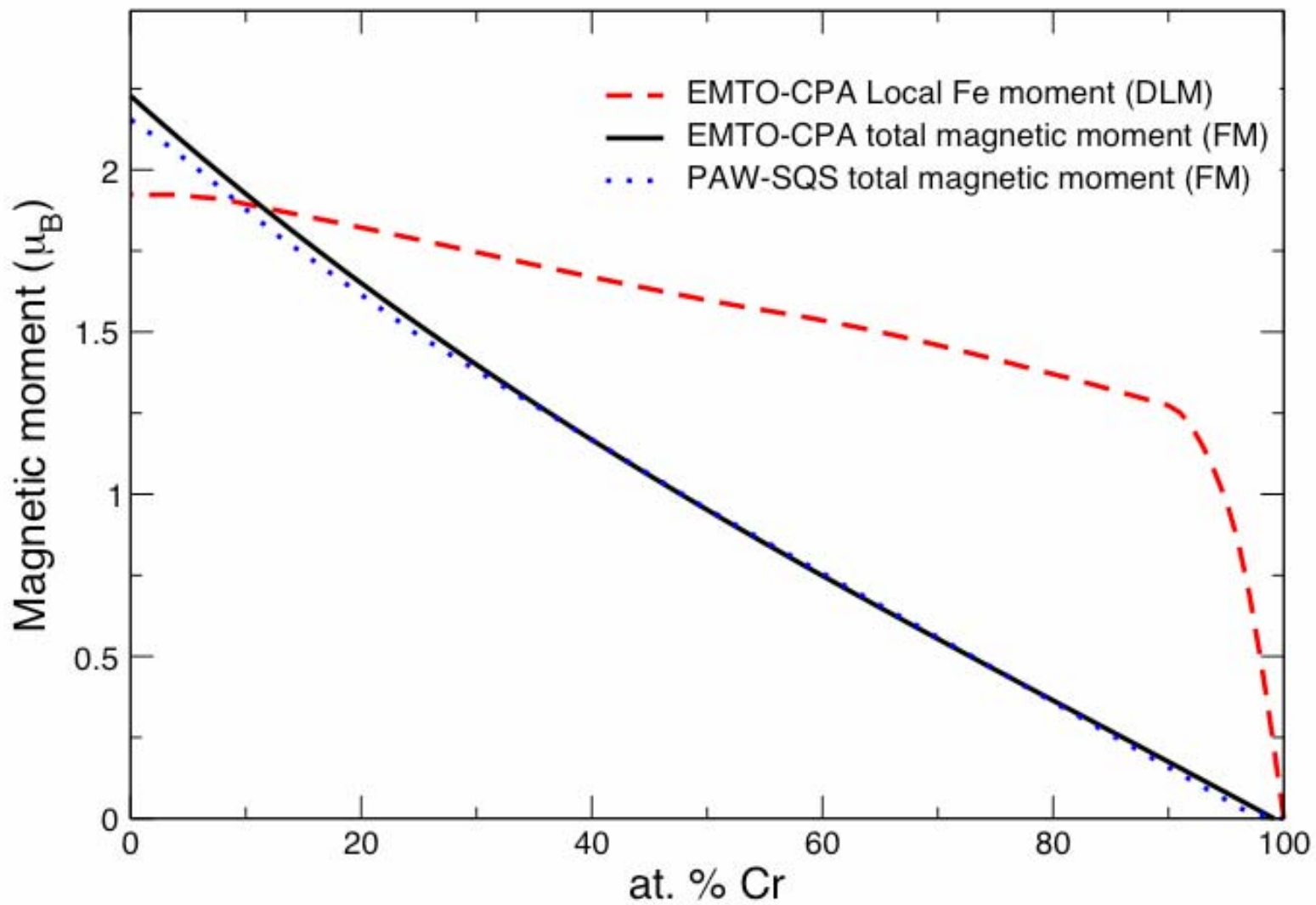


After 30 years at 700 K



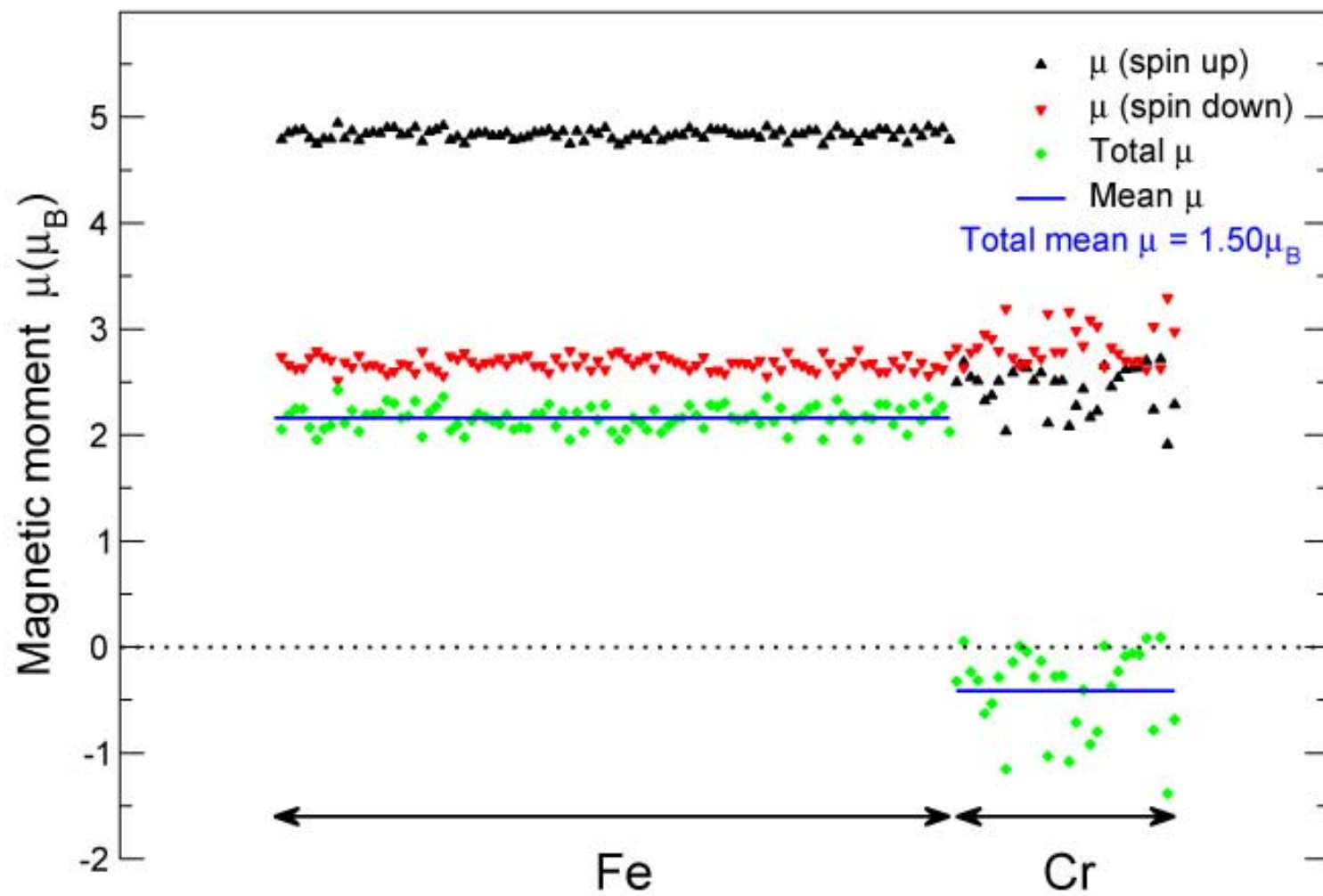


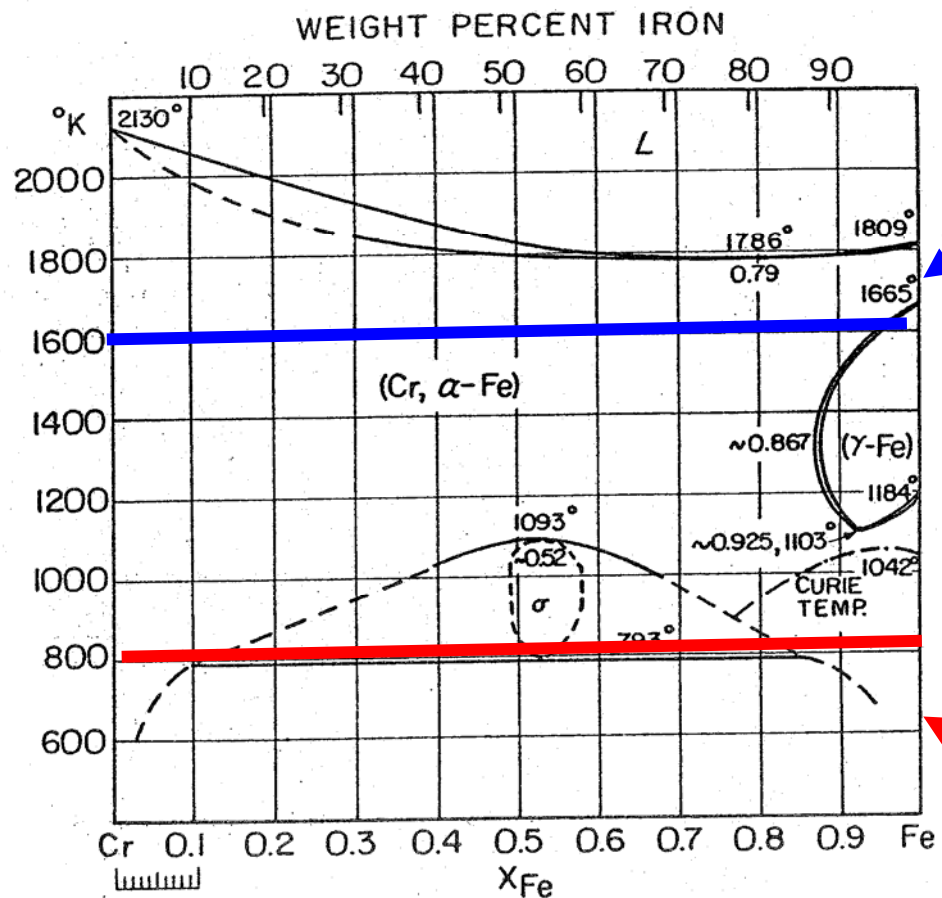




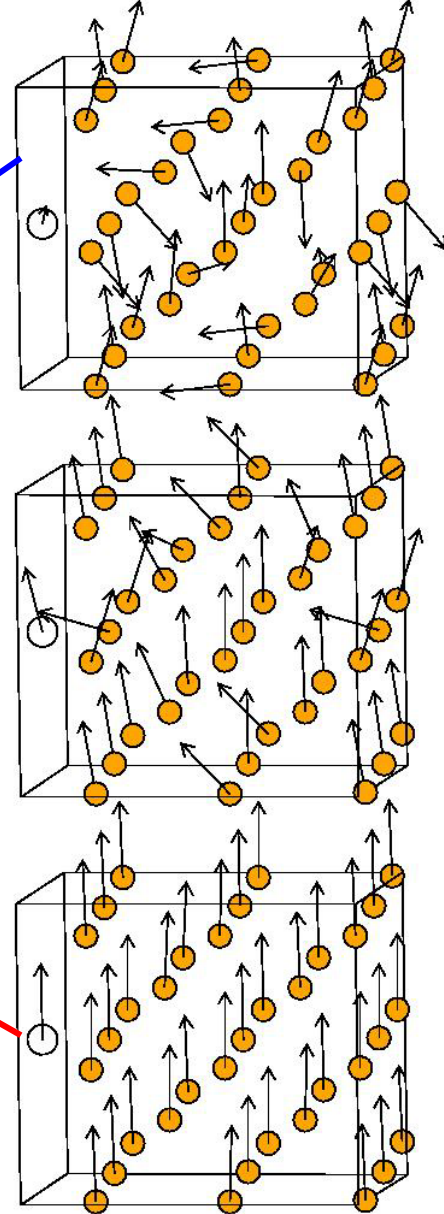
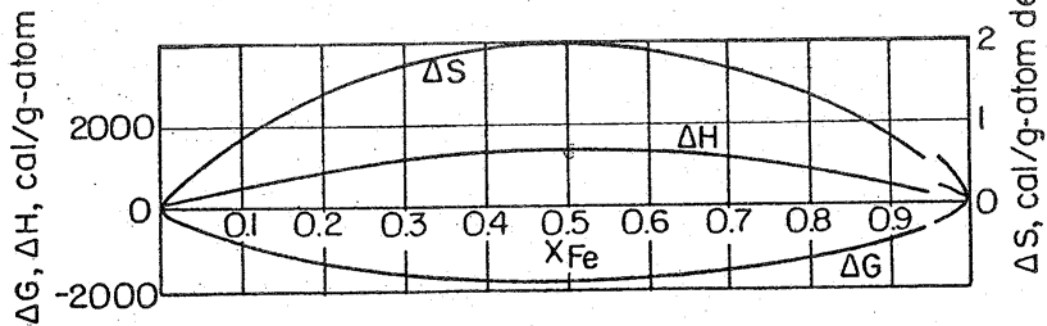
FM bcc Fe₇₅Cr₂₅

128 atom SQS

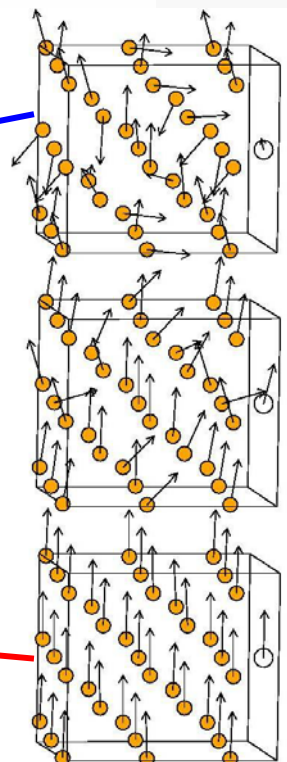
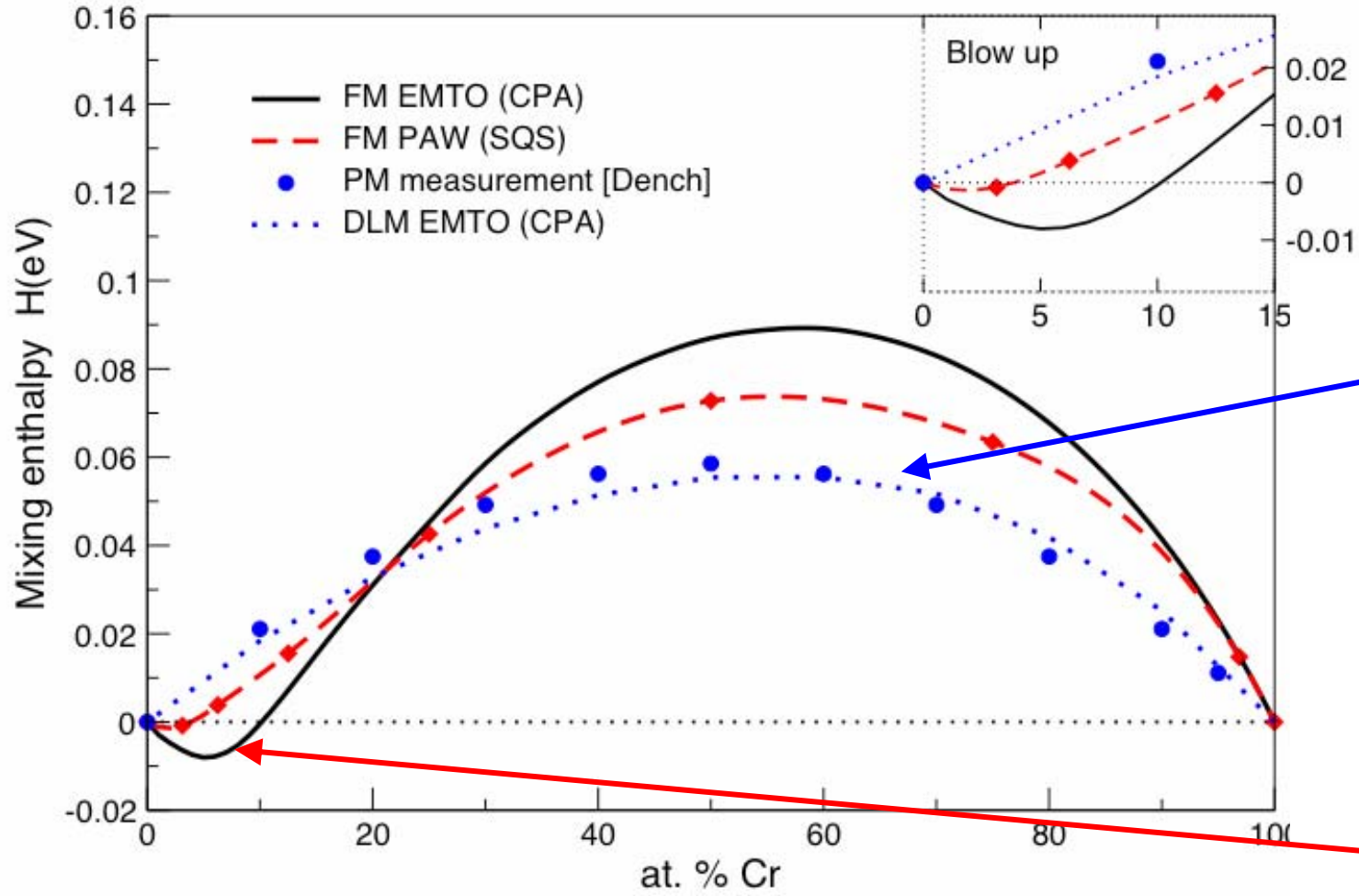


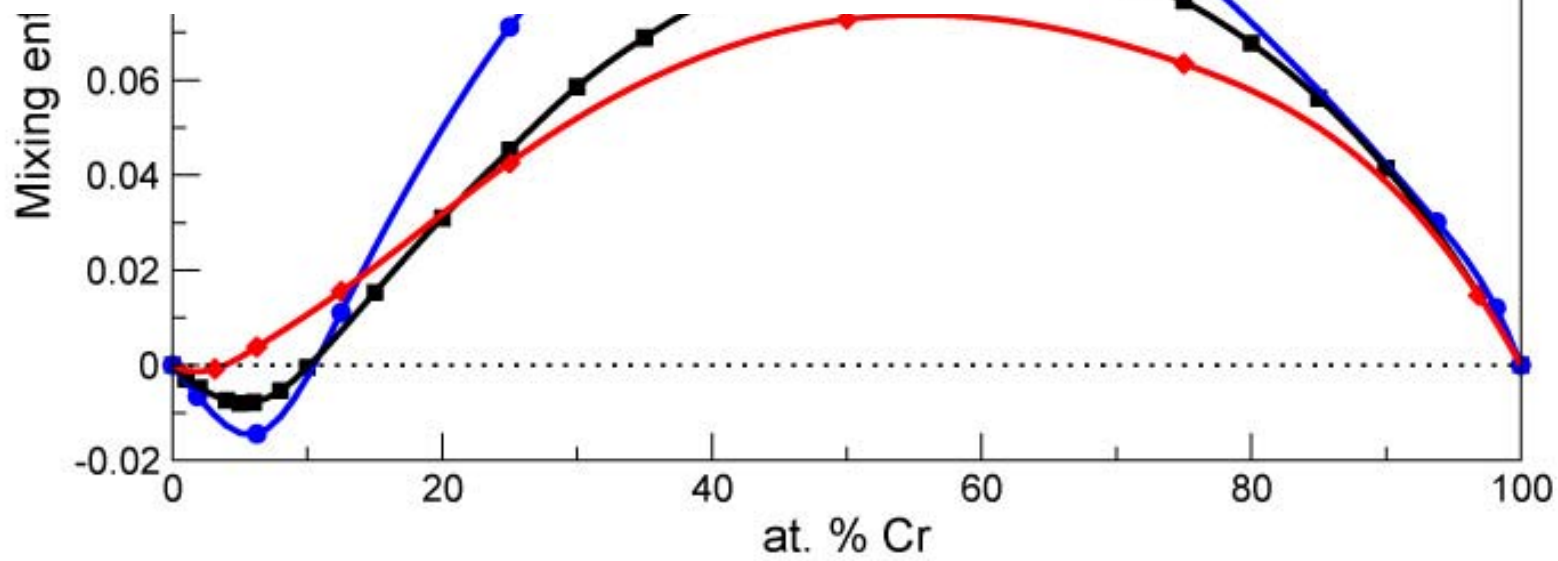
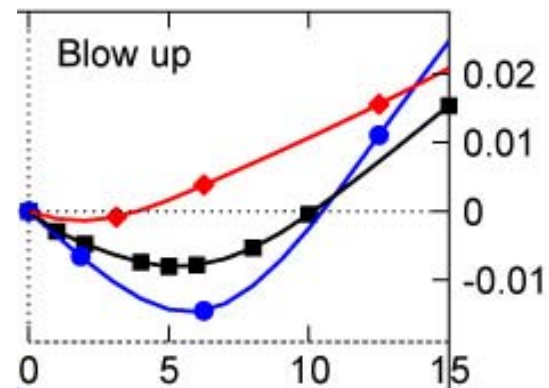
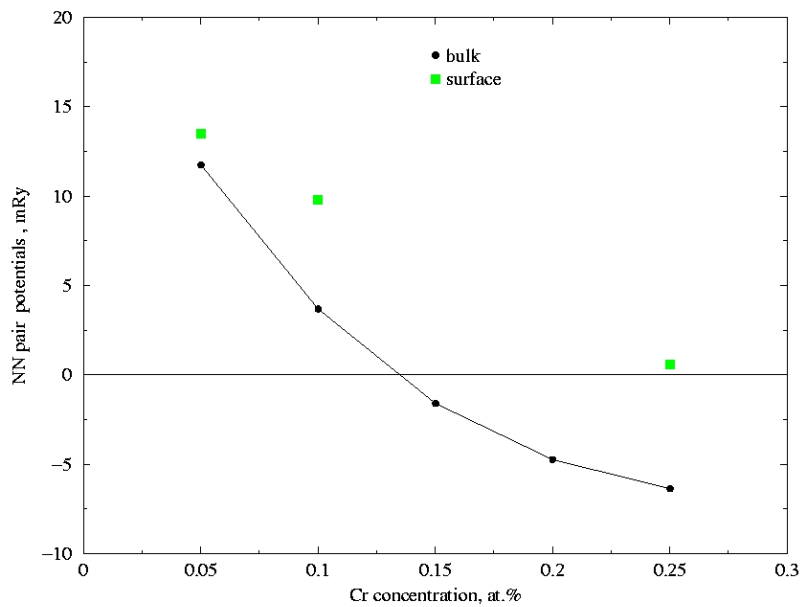


CHROMIUM - IRON SYSTEM, 1600 °K



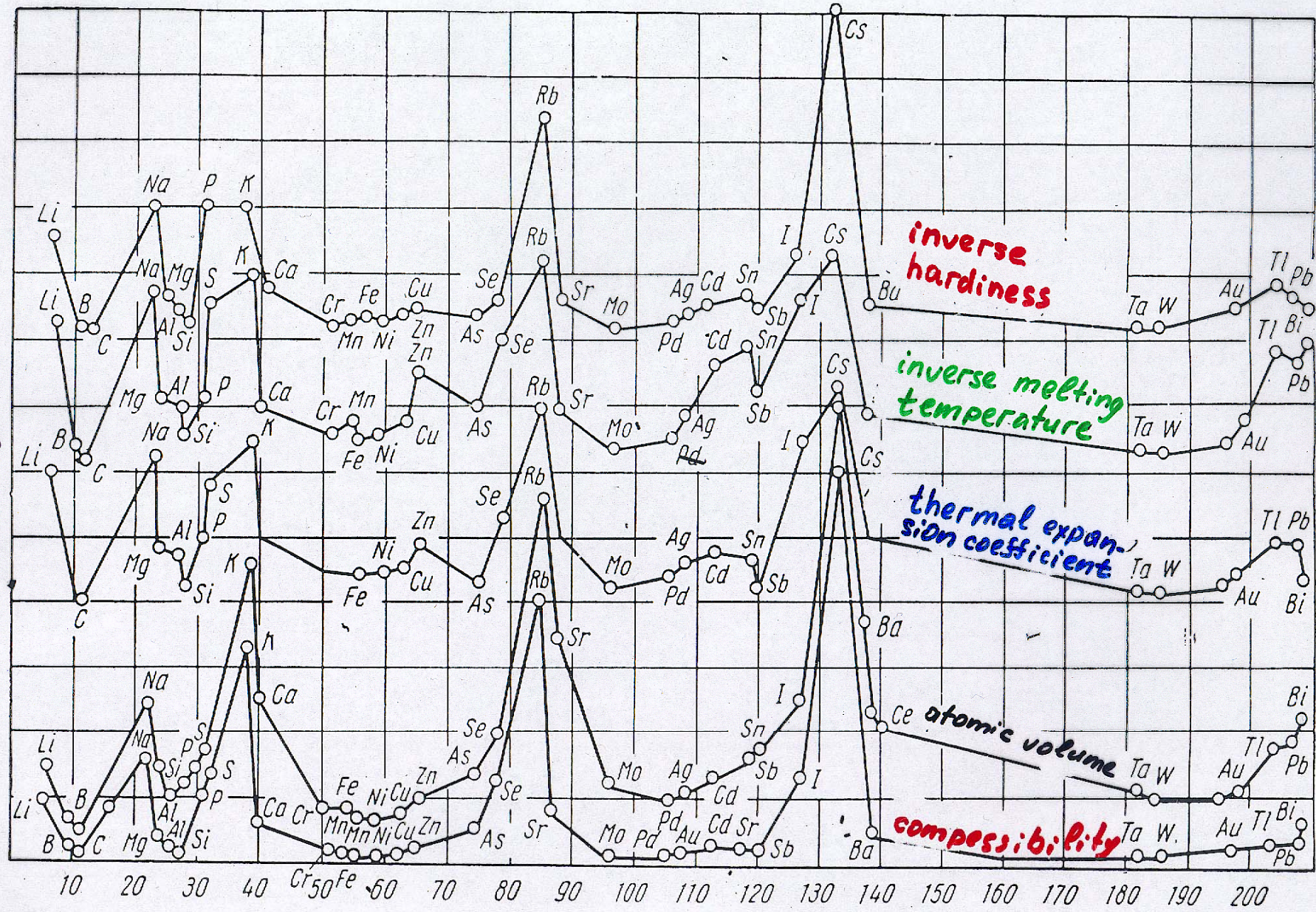
bcc Fe-Cr



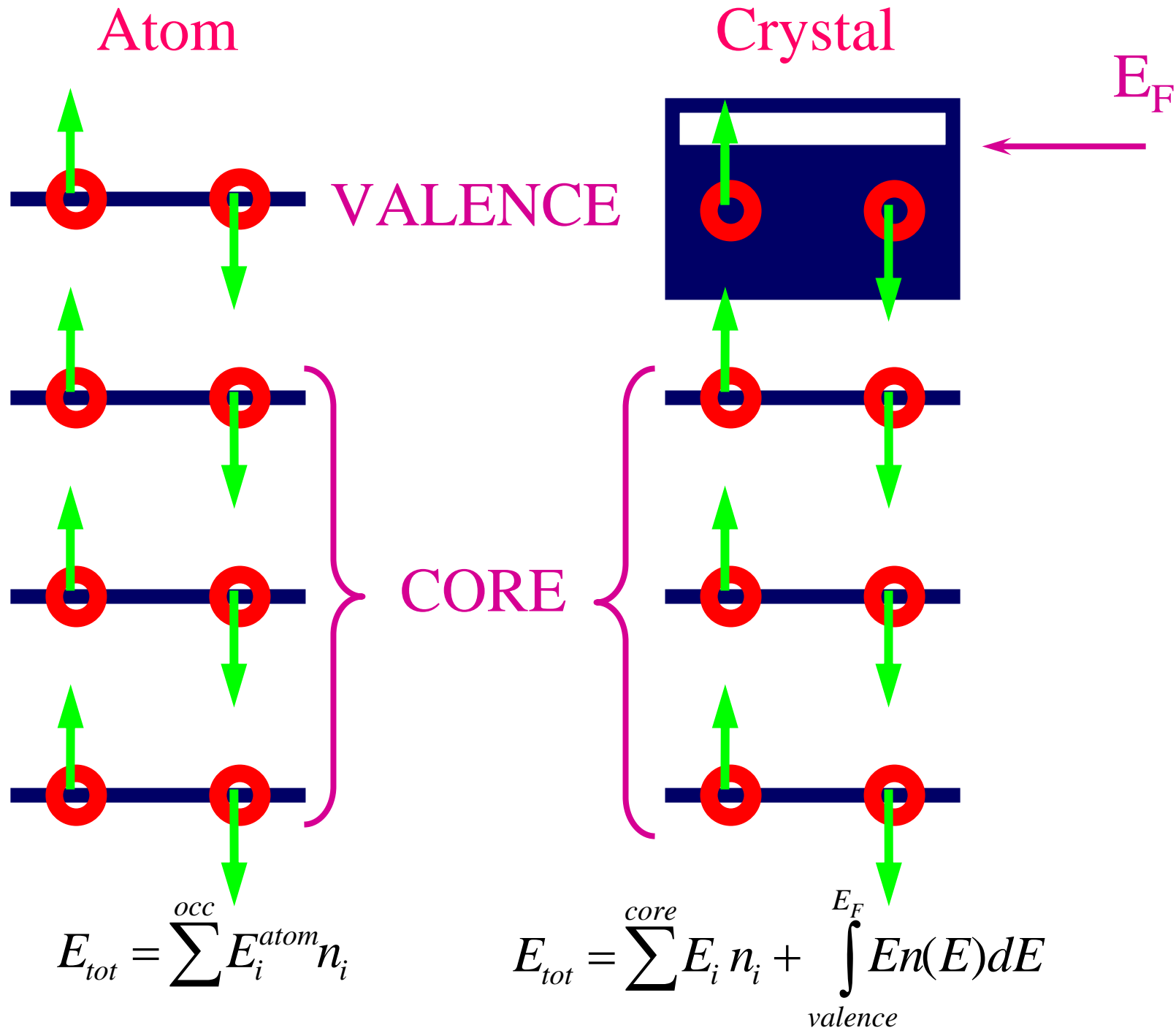


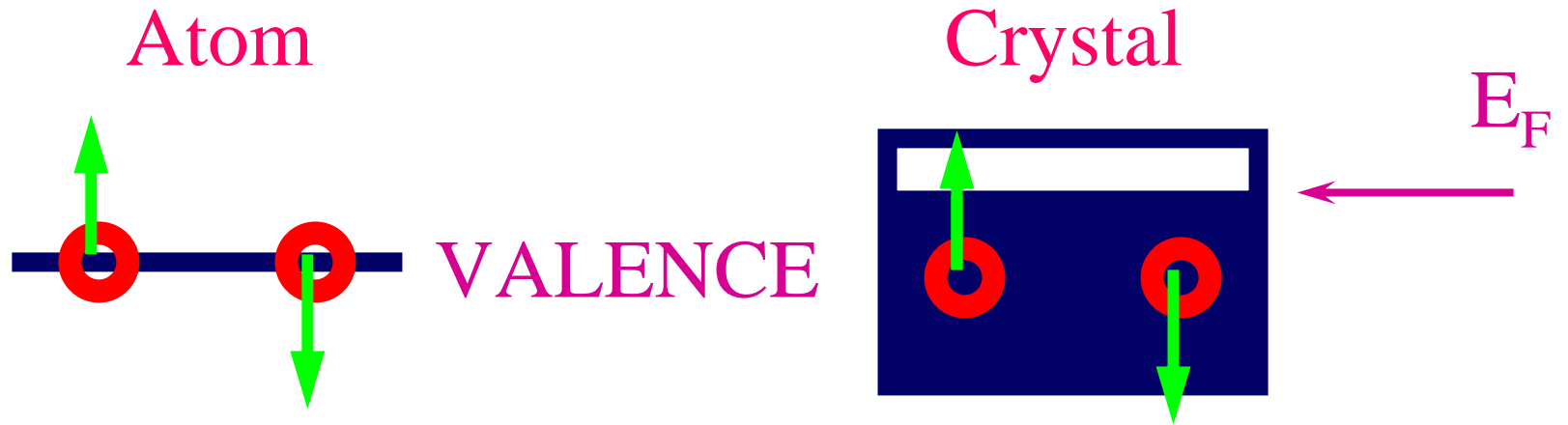
| | | | | | | | | | | | | | | | | | |
|----|----|----|----|----|----|----|----|----|-----|-----|-----|----|----|----|----|----|----|
| H | | | | | | | | | | | | | | | | | He |
| Li | Be | | | | | | | | | | | B | C | N | O | F | Ne |
| Na | Mg | | | | | | | | | | | Al | Si | P | S | Cl | Ar |
| K | Ca | Sc | Ti | V | Cr | Mn | Fe | Co | Ni | Cu | Zn | Ga | Ge | As | Se | Br | Kr |
| Rb | Sr | Y | Zr | Nb | Mo | Tc | Ru | Rh | Pd | Ag | Cd | In | Sn | Sb | Te | I | Xe |
| Cs | Ba | | Hf | Ta | W | Re | Os | Ir | Pt | Au | Hg | Tl | Pb | Bi | Po | At | Rn |
| Fr | Ra | | Rf | Db | Sg | Bh | Hs | Mt | Uun | Uuu | Uub | | | | | | |
| | | | La | Ce | Pr | Nd | Pm | Sm | Eu | Gd | Tb | Dy | Ho | Er | Tm | Yb | Lu |
| | | | Ac | Th | Pa | U | Np | Pu | Am | Cm | Bk | Cf | Es | Fm | Md | No | Lr |

Property



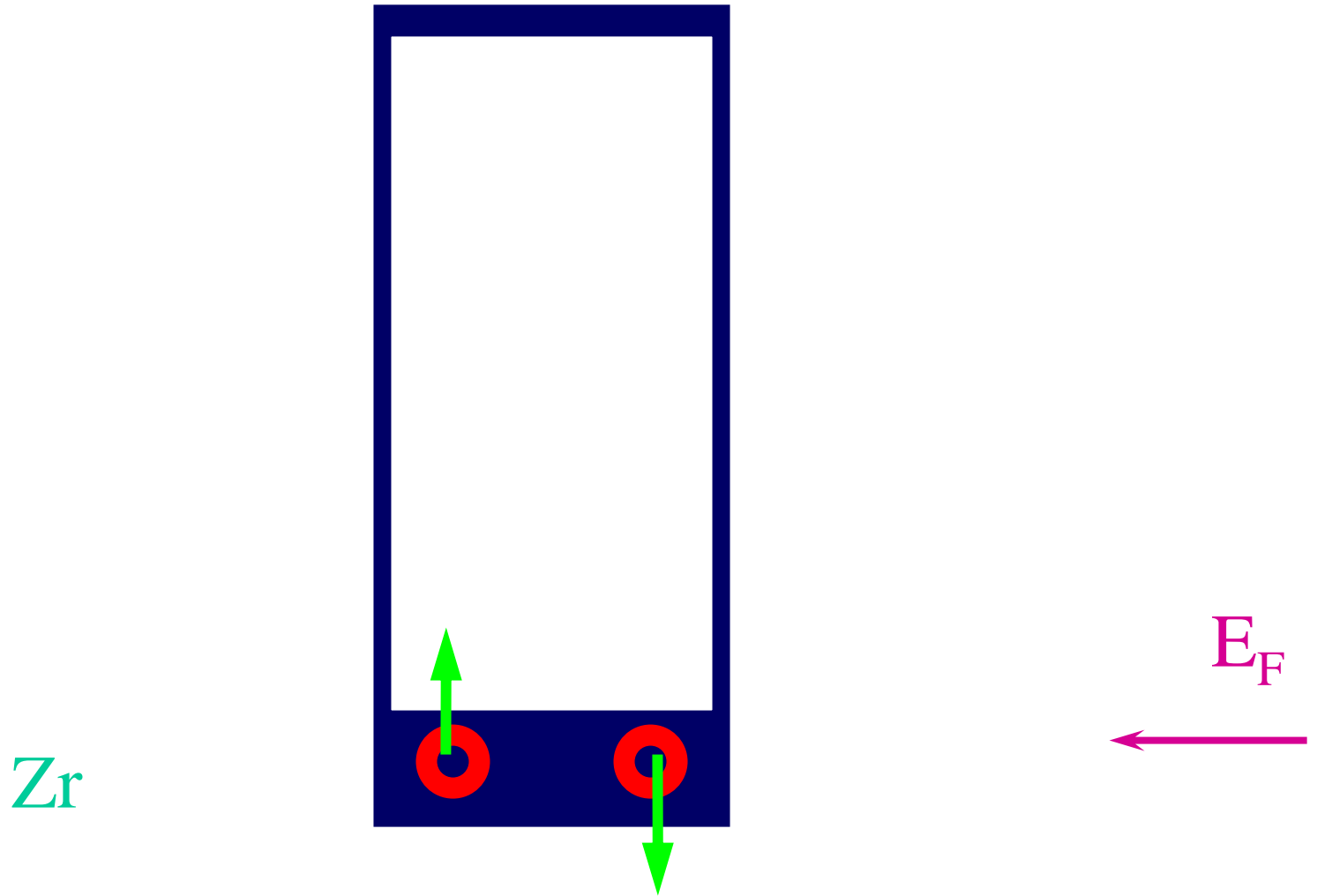
Atomic weight





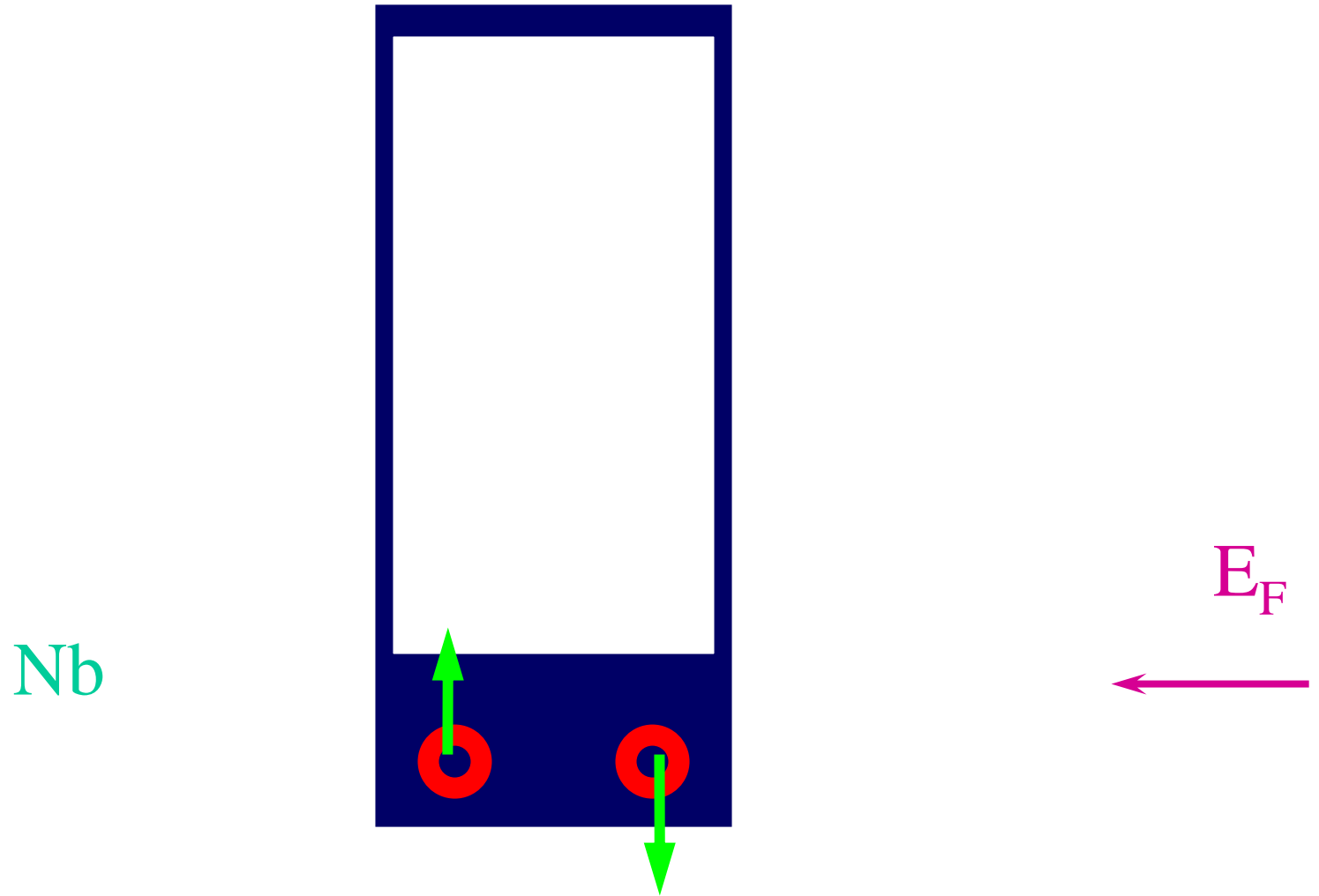
$$E_{BOND} = E_{tot}^{crystal} - E_{tot}^{atom} = \int_{valence}^{E_F} (E - E_{valence}^{atom}) n(E) dE$$

Transition metal



$$E_{BOND} = E_{tot}^{crystal} - E_{tot}^{atom} = \int_{valence}^{E_F} (E - E_{valence}^{atom}) n(E) dE$$

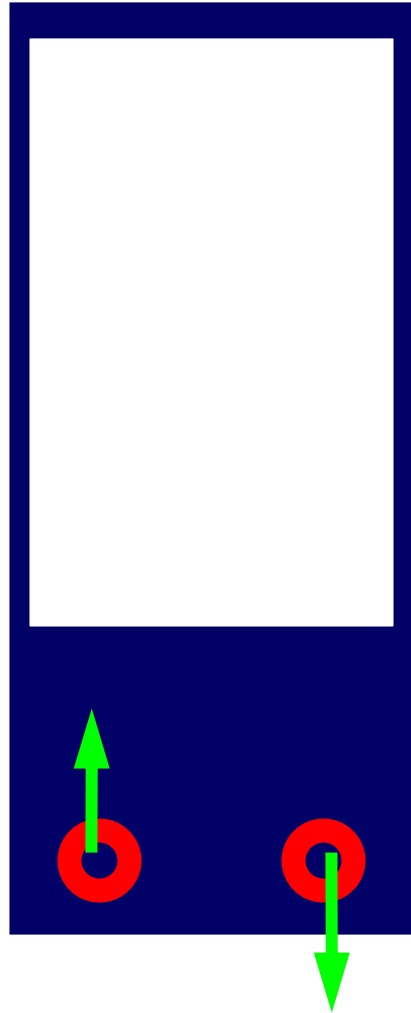
Transition metal



$$E_{BOND} = E_{tot}^{crystal} - E_{tot}^{atom} = \int_{valence}^{E_F} (E - E_{valence}^{atom}) n(E) dE$$

Transition metal

Mo

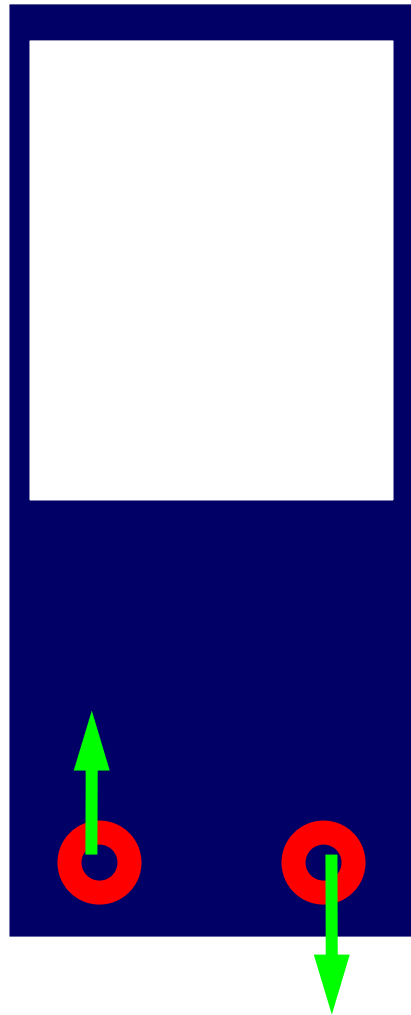


E_F

$$E_{BOND} = E_{tot}^{crystal} - E_{tot}^{atom} = \int_{valence}^{E_F} (E - E_{valence}^{atom}) n(E) dE$$

Transition metal

Tc

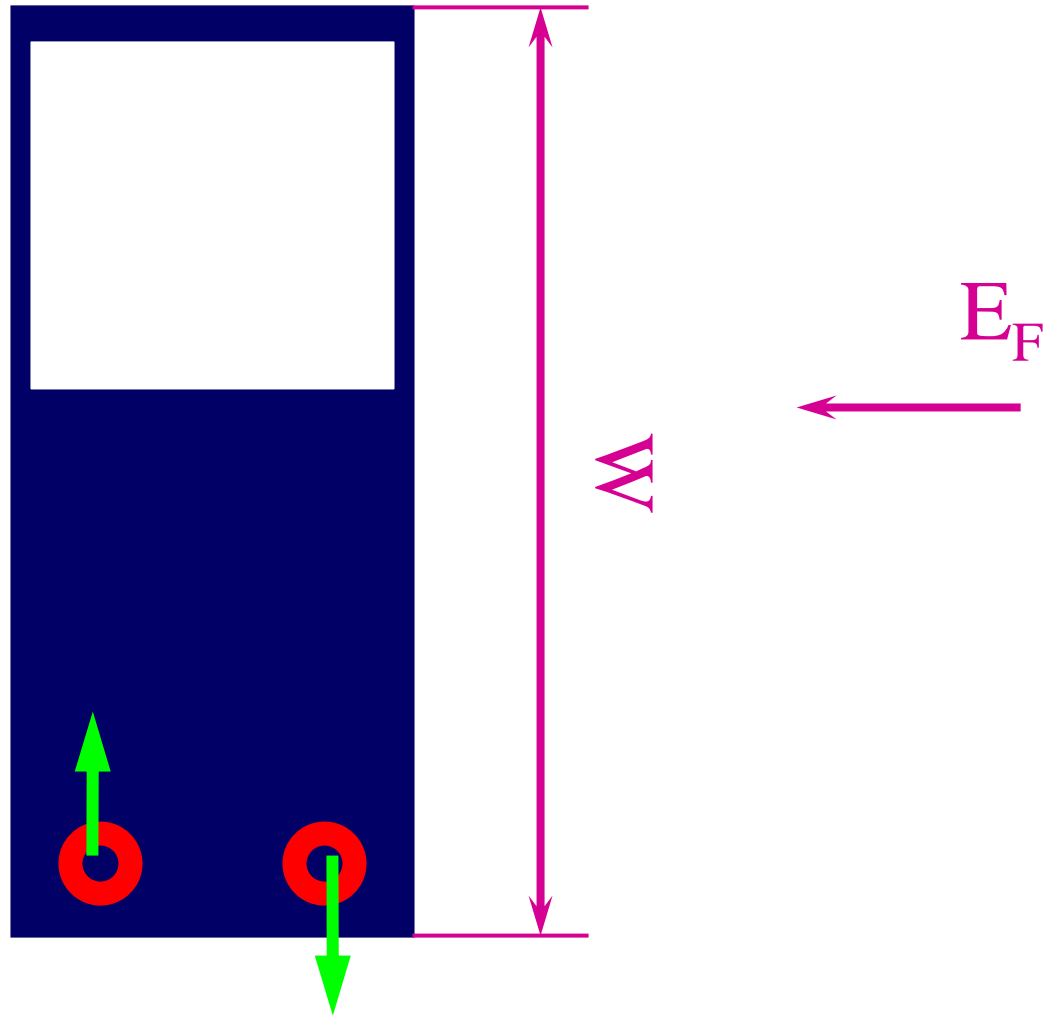


E_F

$$E_{BOND} = E_{tot}^{crystal} - E_{tot}^{atom} = \int_{valence}^{E_F} (E - E_{valence}^{atom}) n(E) dE$$

Transition metal

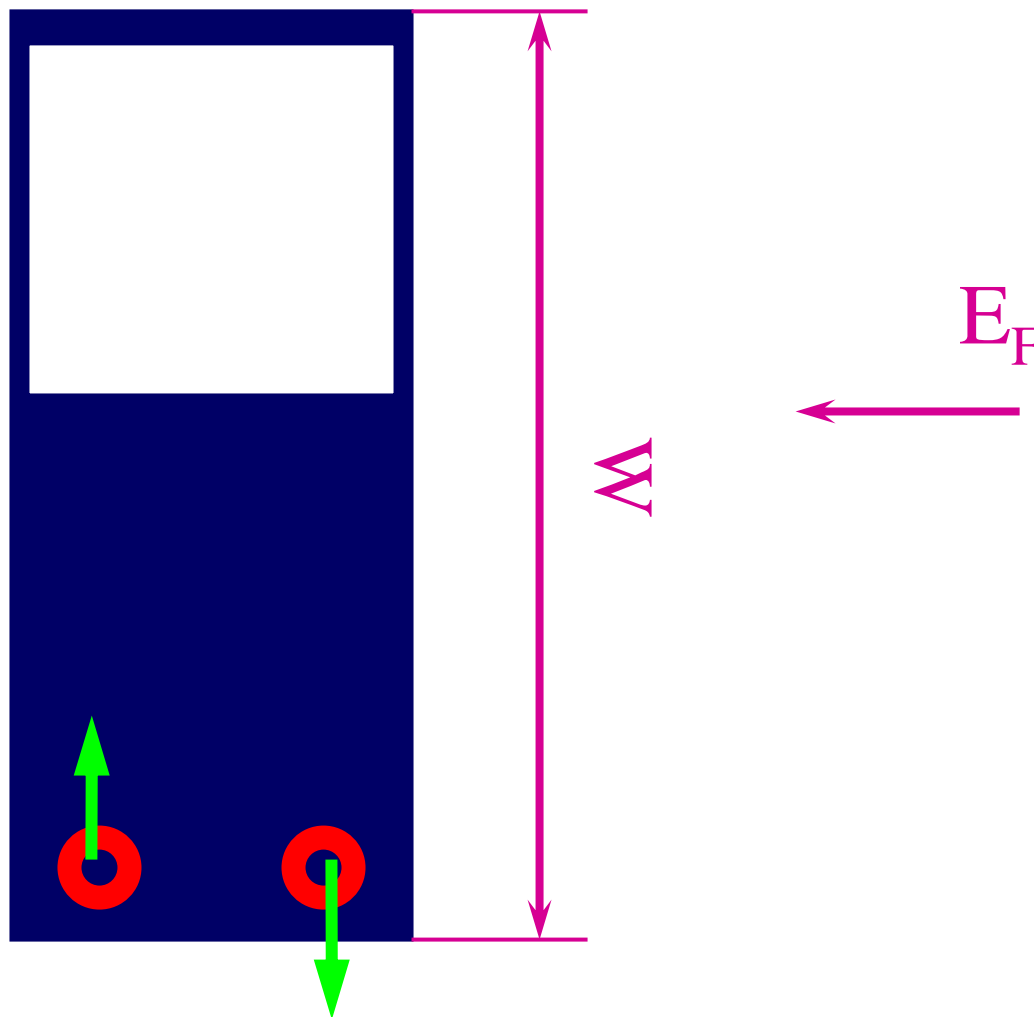
Ru



$$E_{BOND} = E_{tot}^{crystal} - E_{tot}^{atom} = \int_{valence}^{E_F} (E - E_{valence}^{atom}) n(E) dE$$

Transition metal

Ru



$$E_{BOND} = -\frac{1}{20} W N_d (10 - N_d)$$

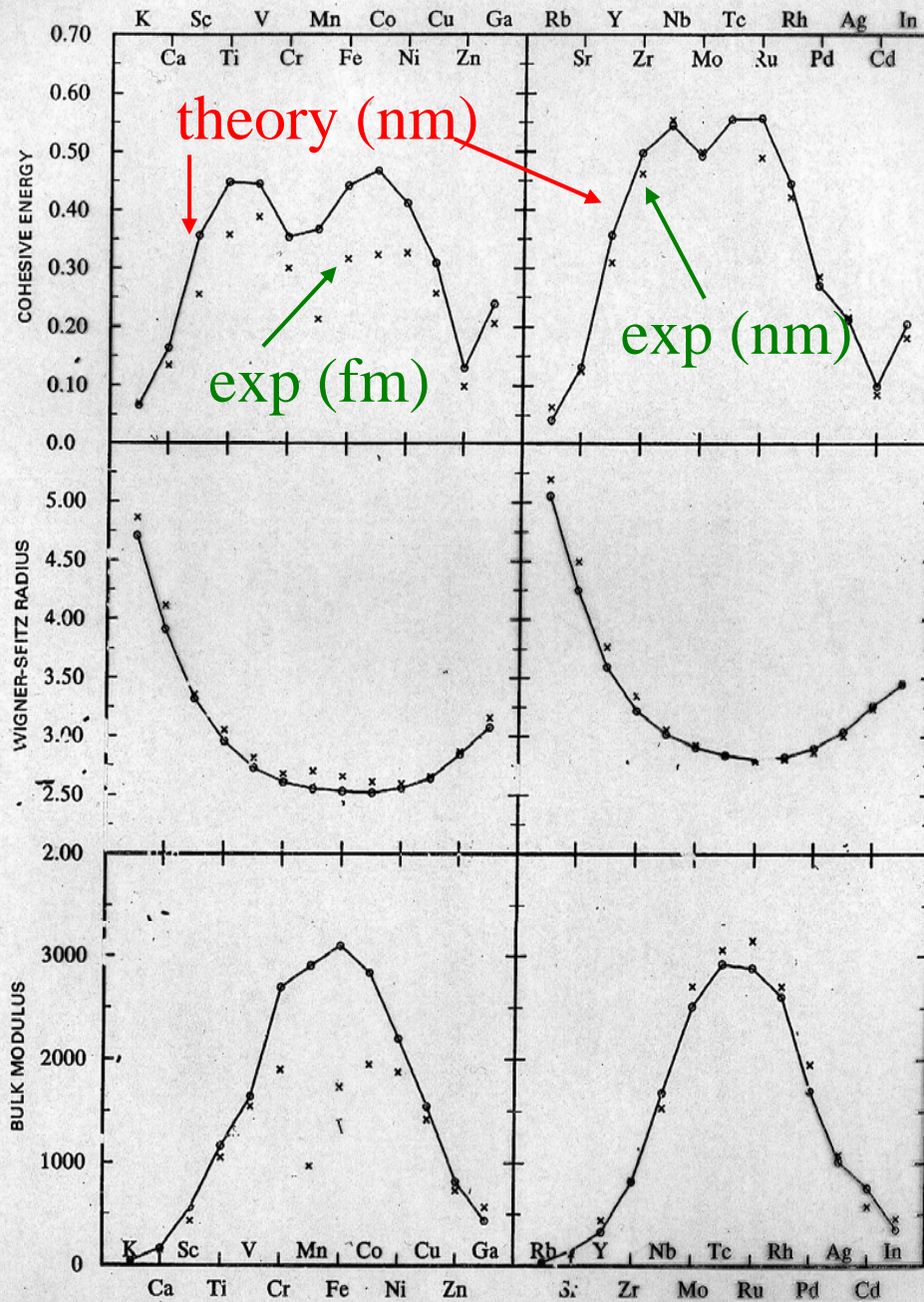
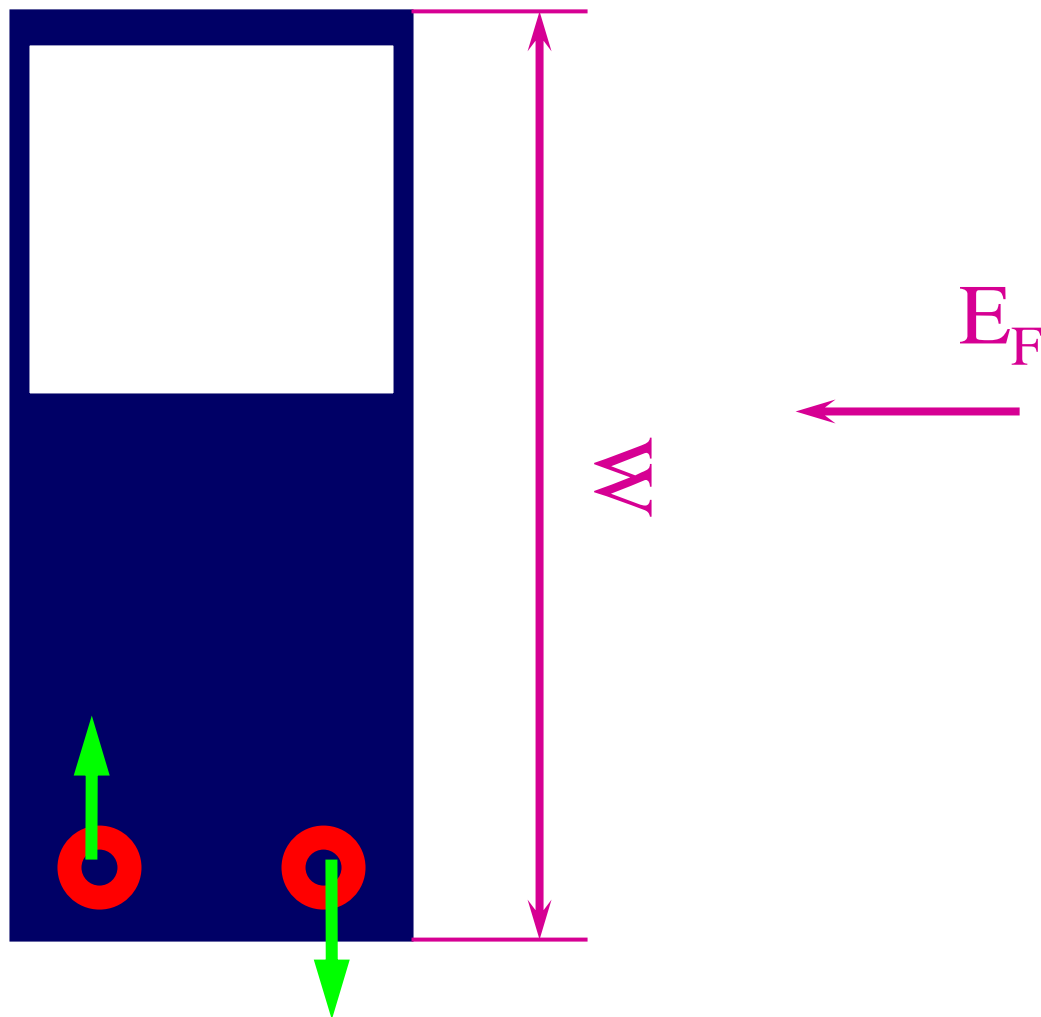


Figure 1.1 Cohesive properties. Top row- cohesive energy (Ry/atom). Middle row- Wigner-Seitz radius (a.u.). Bottom row- bulk modulus (Kbar). Measured values are indicated by crosses.

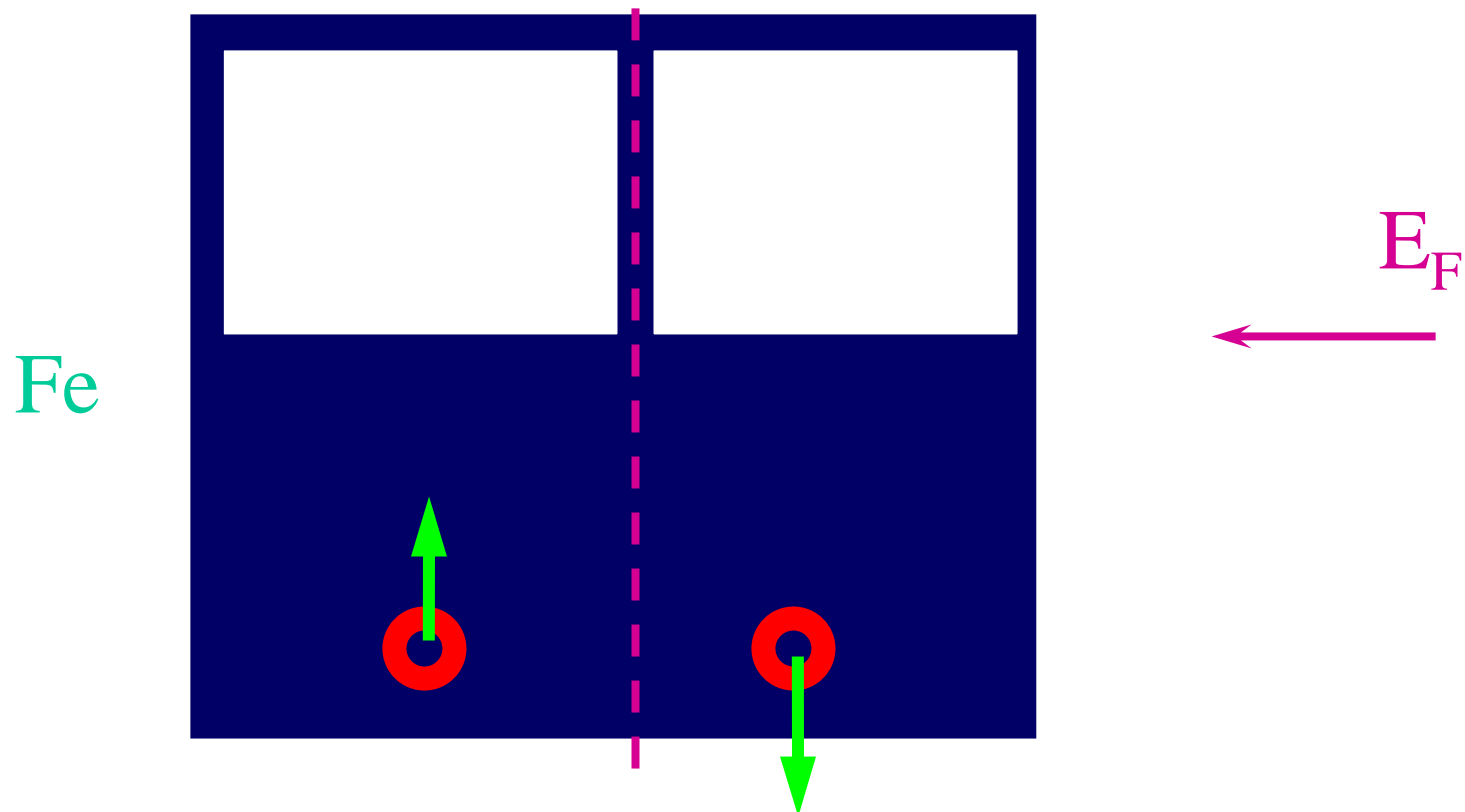
Transition metal

Ru



$$E_{BOND} = -\frac{1}{20} W N_d (10 - N_d)$$

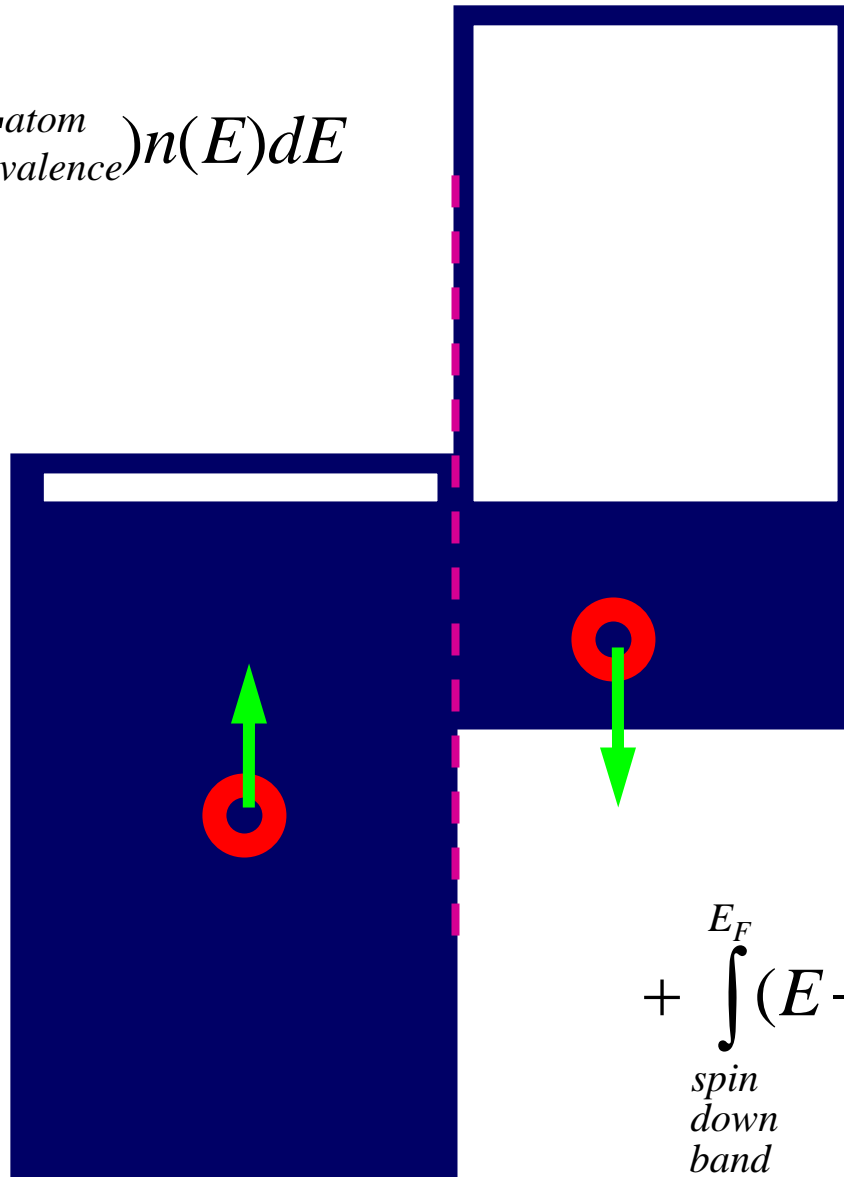
Transition metal



$$E_{BOND} = -\frac{1}{20} W N_d (10 - N_d)$$

$$E_{BOND} = \int_{\substack{spin \\ up \\ band}}^{E_F} (E - E_{valence}^{atom}) n(E) dE$$

Fe



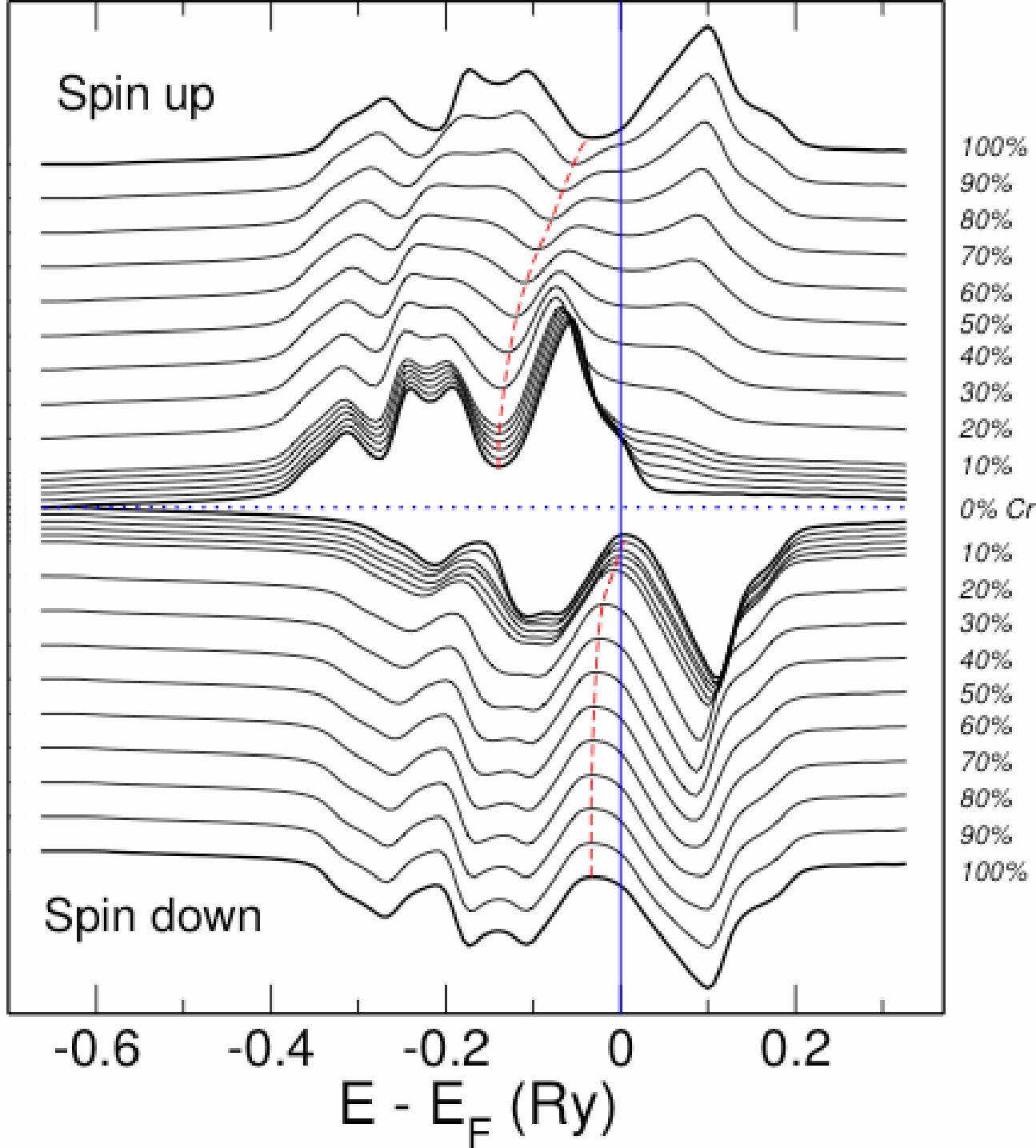
E_F



$$+ \int_{\substack{spin \\ down \\ band}}^{E_F} (E - E_{valence}^{atom}) n(E) dE$$

DIFFERENT PROPERTIES !!!

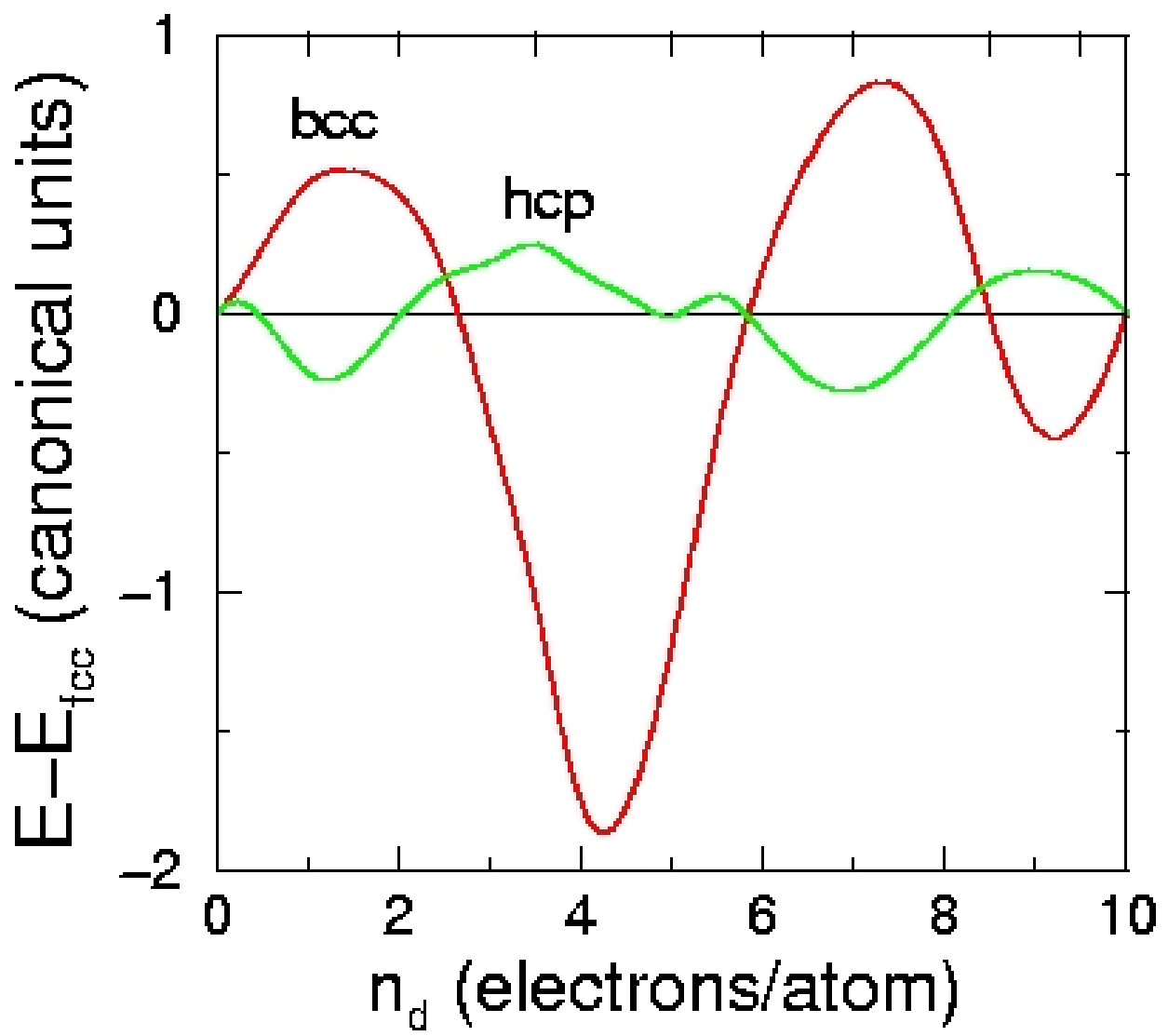
FM density of states (arbitrary units)

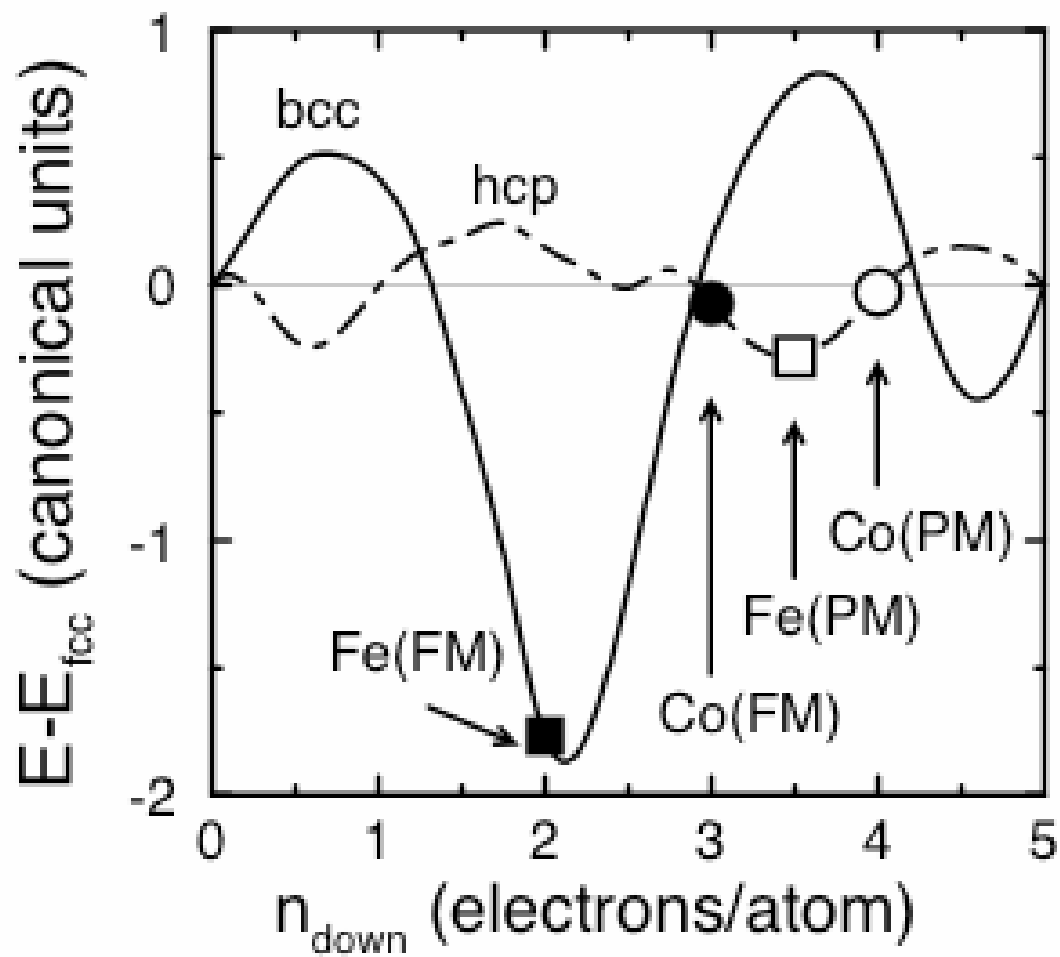


| | | | | | | | | | | | | | | | | | | |
|----|----|----|----|----|----|----|----|----|-----|-----|-----|----|----|----|----|----|----|--|
| H | | | | | | | | | | | | | | | | | He | |
| Li | Be | | | | | | | | | | | B | C | N | O | F | Ne | |
| Na | Mg | | | | | | | | | | | Al | Si | P | S | Cl | Ar | |
| K | Ca | Sc | Ti | V | Cr | Mn | Fe | Co | Ni | Cu | Zn | Ga | Ge | As | Se | Br | Kr | |
| Rb | Sr | Y | Zr | Nb | Mo | Tc | Ru | Rh | Pd | Ag | Cd | In | Sn | Sb | Te | I | Xe | |
| Cs | Ba | | Hf | Ta | W | Re | Os | Ir | Pt | Au | Hg | Tl | Pb | Bi | Po | At | Rn | |
| Fr | Ra | | Rf | Db | Sg | Bh | Hs | Mt | Uun | Uuu | Uub | | | | | | | |
| | | La | Ce | Pr | Nd | Pm | Sm | Eu | Gd | Tb | Dy | Ho | Er | Tm | Yb | Lu | | |
| | | Ac | Th | Pa | U | Np | Pu | Am | Cm | Bk | Cf | Es | Fm | Md | No | Lr | | |

4d-element

Y Zr Nb Mo Tc Ru Rh Pd





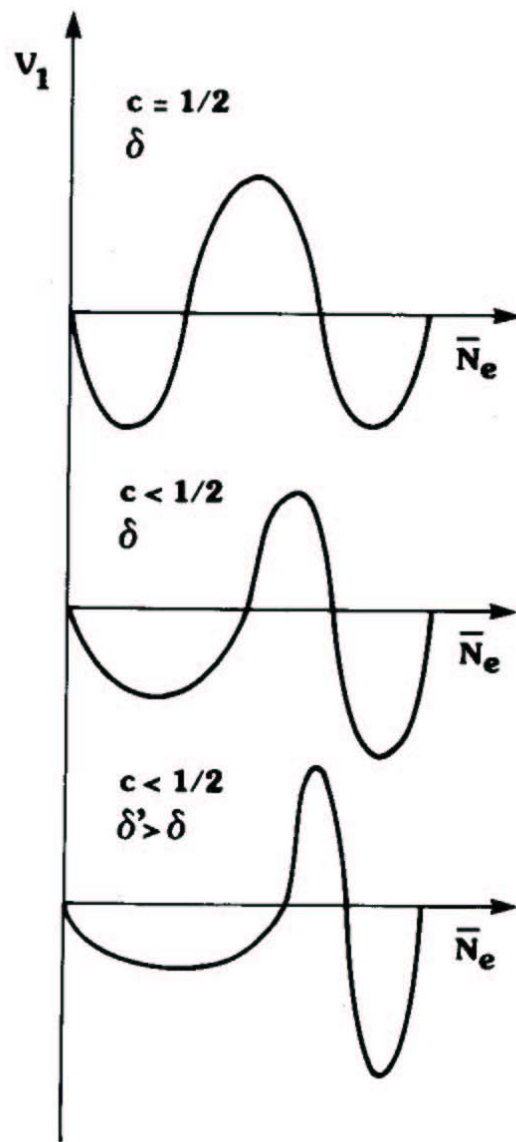


Fig. 7.5. Typical variations of $V_1(\bar{N}_e)$ for different values of c and of δ .

Chemical sro effects in ferromagnetic Fe alloys in relation to electronic band structure

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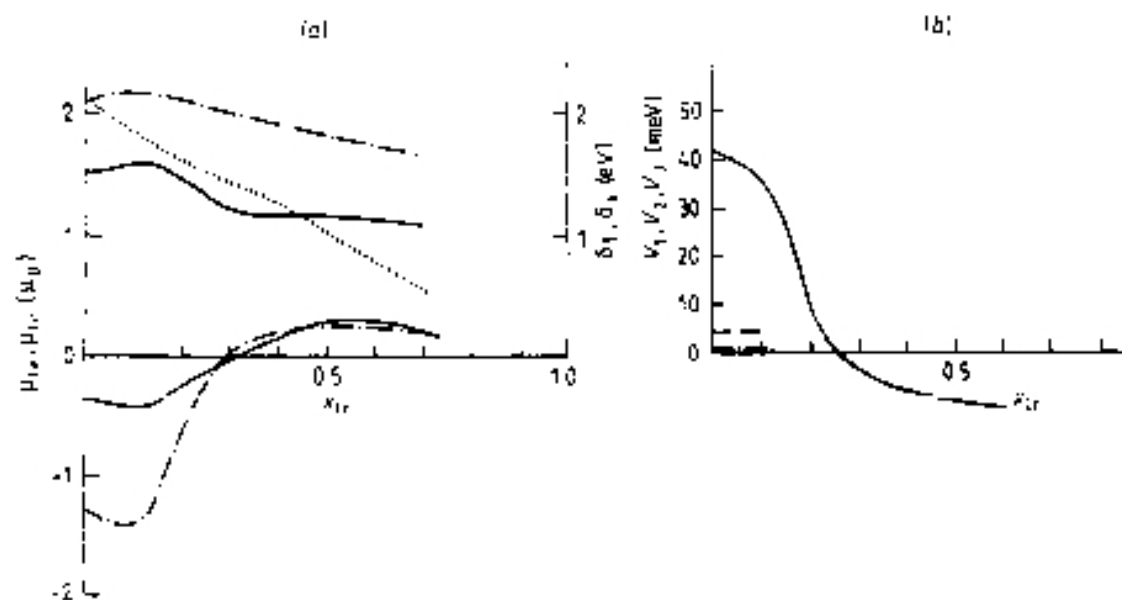
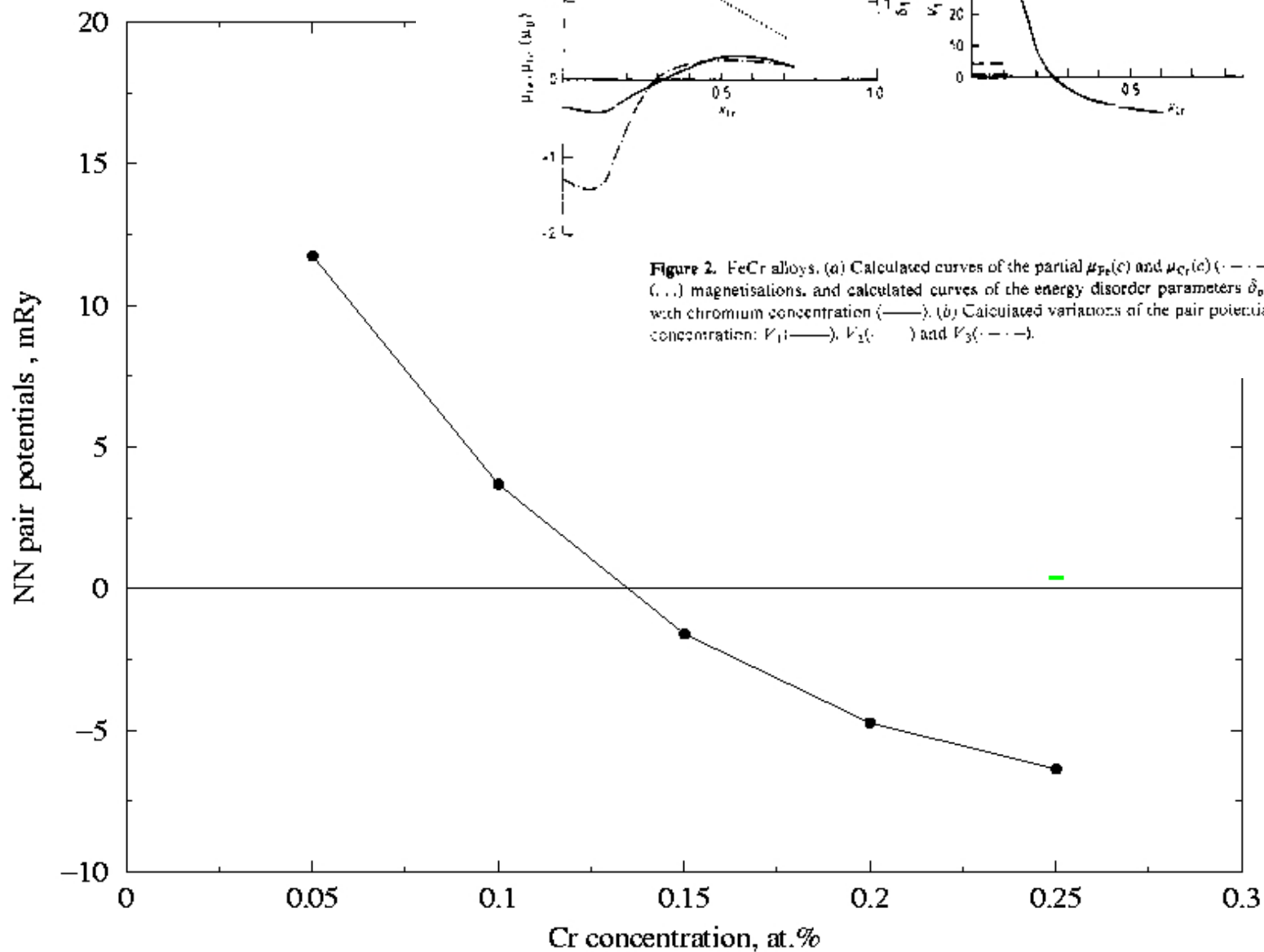


Figure 2. FeCr alloys. (a) Calculated curves of the partial $\mu_{Fe}(c)$ and $\mu_{Cr}(c)$ (---) and total (—) magnetisations, and calculated curves of the energy disorder parameters δ_{σ} ($\sigma = \uparrow, \downarrow$) with chromium concentration (—). (b) Calculated variations of the pair potentials with Cr concentration: V_1 (—), V_2 (---) and V_3 (— · —).



CONCLUSIONS :



- First-principles simulations can be carried out for real materials of technological importance.
- The results allow for the cautious optimism, and the choice of methodology should depend on the problem at hand.
- The mixing enthalpy for paramagnetic alloys is positive at all concentrations, in excellent agreement with experiment.
- On the contrary, ferromagnetic bcc Fe-Cr alloys are anomalously stable at low Cr concentrations. Therefore, magnetic effects **must** be taken into account in simulations.
- The stability of the ferromagnetic alloys is determined by the strong concentration dependence of interatomic interactions in this system, which therefore **must** be taken into account in simulations.