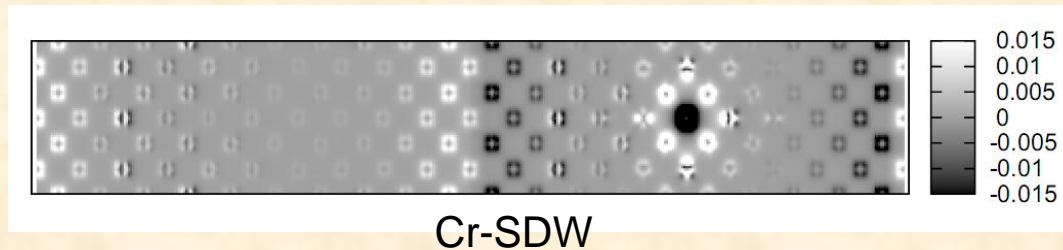
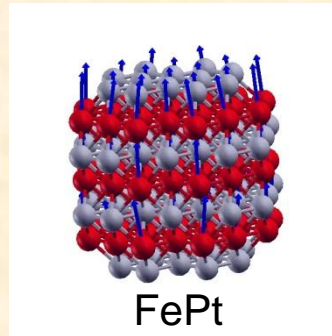


Correlation Between Magnetism and Structure in Fe alloys: the case of Fe-Cr and Fe-Pt

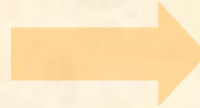
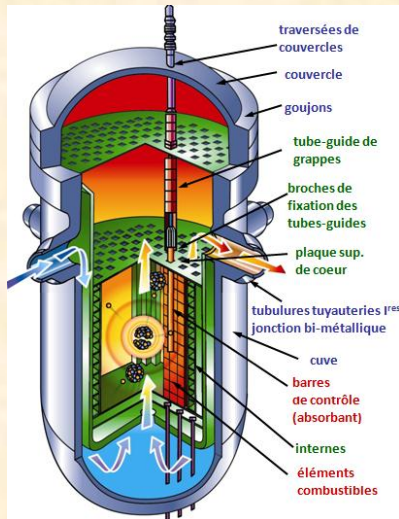
Cyrille Barreteau (SPCSI)
Chu Chun Fun (SRMP)
Romain Soulairol (SRMP)
Daniel Spanjaard (LPS)



Service de Physique et Chimie des Surfaces et Interfaces (**SPSCI**)
Service de Recherche sur la Métallurgie Physique (**SRMP**)
Laboratoire de Physique des Solides (**LPS**)

Why Fe-Cr

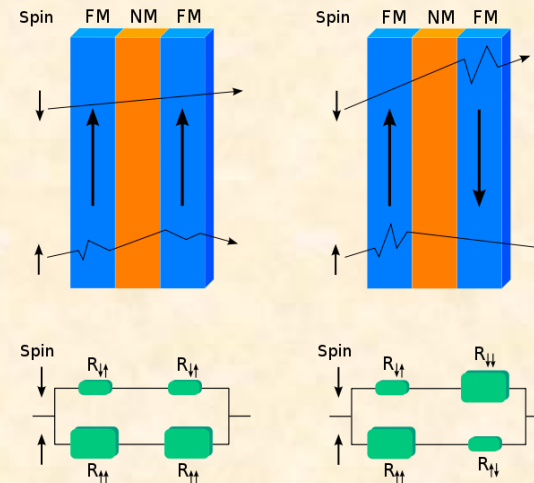
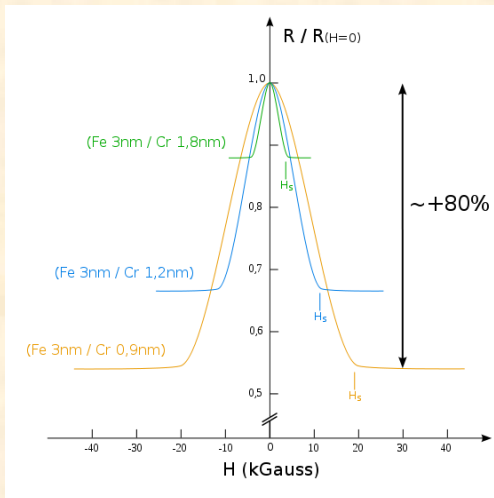
Material for nuclear industry



- FeCr alloy: resistance to corrosion, and irradiation. Decrease of swelling

- Structure material for future nuclear reactors (fission, fusion).

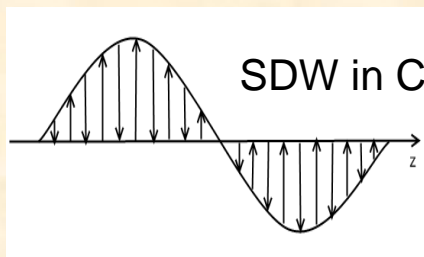
GMR in FeCr multilayers



Why Fe-Cr

Complex magnetic order

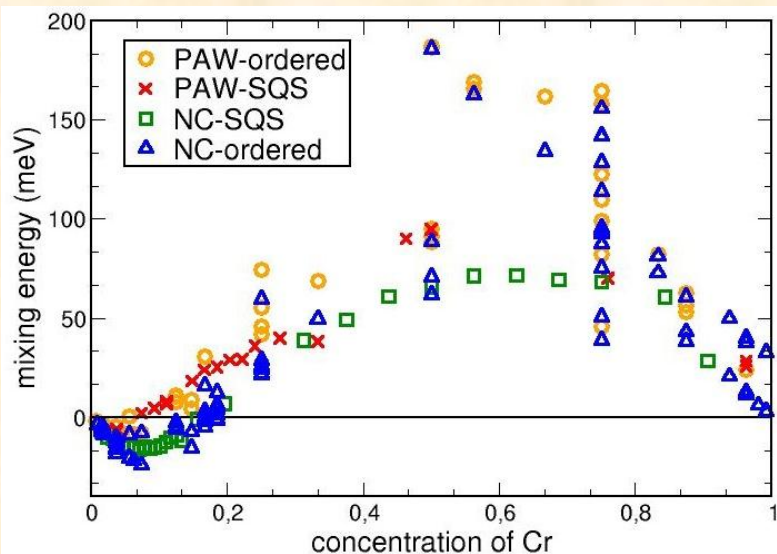
Spin spiral in Fe fcc



SDW in Cr bcc

How magnetic order modifies the energetic of defects and vice versa..

Mixing energy



Influence of magnetism on alloy properties

Methods

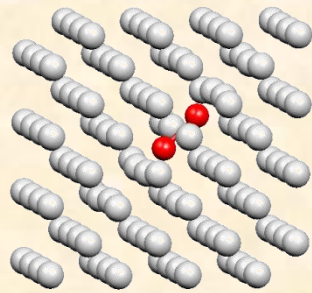
Simulated time

1h

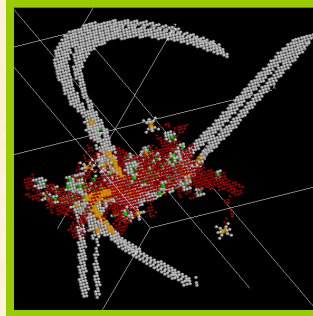
1ns

0

DFT (SIESTA, PWSCF)
Tight-Binding (TB)

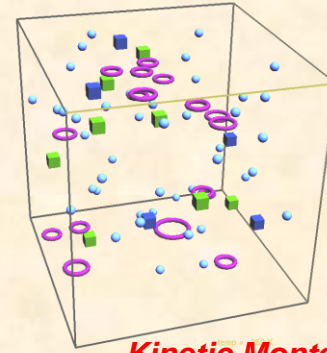


100 at. 1000 at.



Molecular dynamics with semi-empirical potentials

10nm



Kinetic Monte Carlo

Size of the system

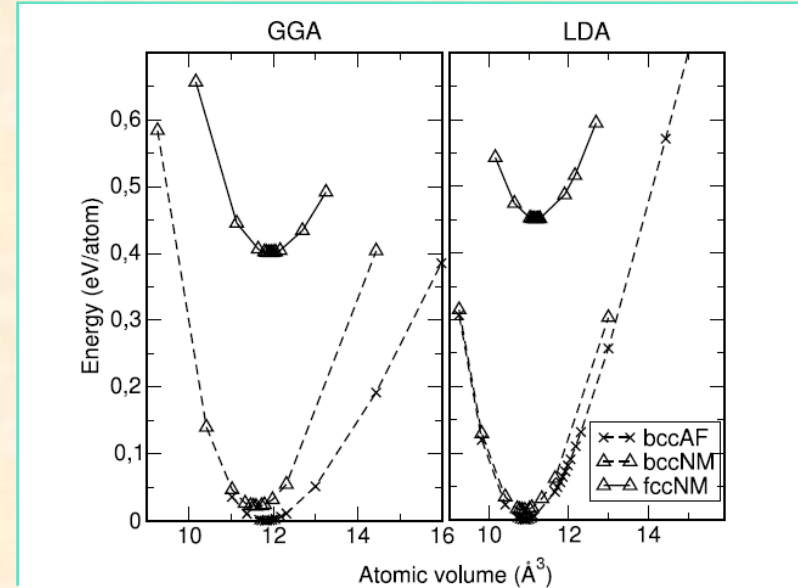
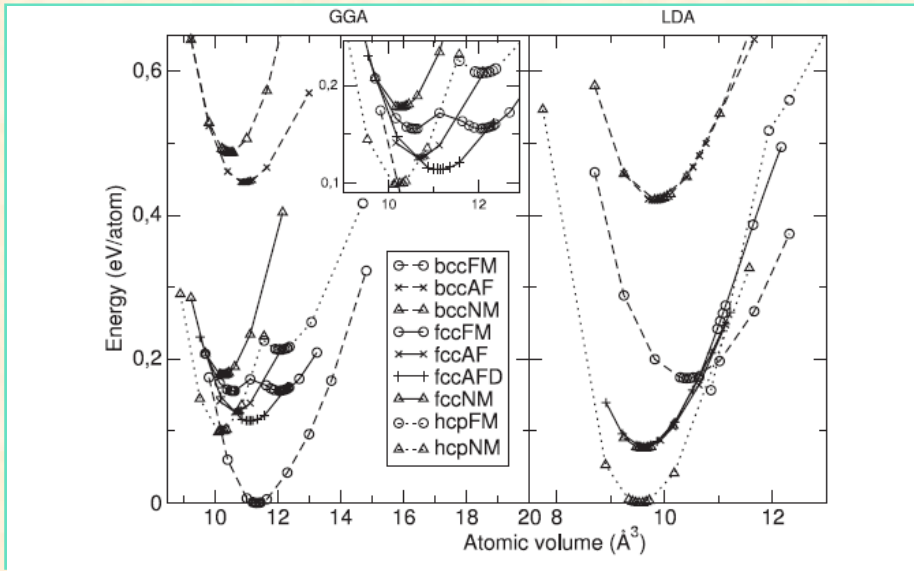
1 μ m

Code	DFT - PWscf	DFT - SIESTA	TB-Stoner
E_{xc}	GGA	GGA	-
Pseudopot.	NC, US, PAW	NC	-
Basis	Plane waves	Localized	Localized (spd)
Efficiency/precision	Very precise	Precise and efficient	Very efficient
Size of the system	< 500 at	< 1000 at	> 1000 at

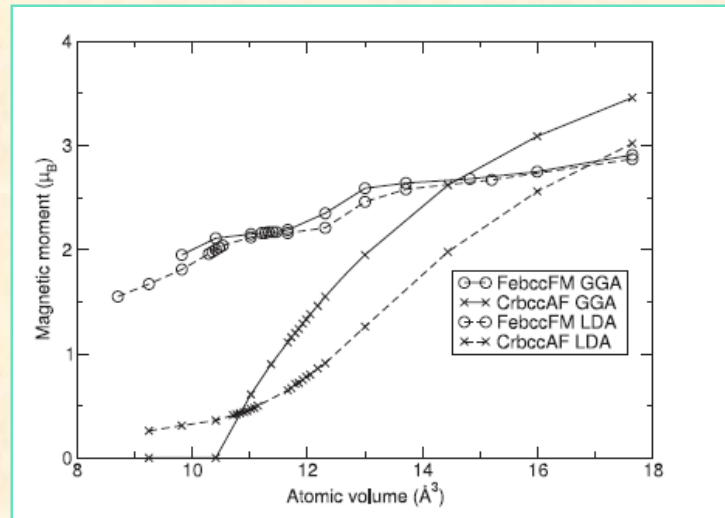
Testing the methods

Functional Effect: LDA vs GGA (PWscf PAW)

E(d)



M(d)



Testing the methods

Pseudopotential Effect NC, US, PAW

 $\Delta E(\text{eV})$

	GGA				LDA			
	NC/LCAO	PAW	US	NC	NC/LCAO	PAW	US	NC
FebccFM	0	0	0	0	0.13	0.17	0.13	0.16
FefccFM-HS	0.11	0.16	0.13	0.11	—	—	—	—
FefccFM-LS	0.20	0.16	0.18	0.22	—	—	—	—
FefccAF	0.15	0.13	0.14	0.17	0.07	—	—	—
FefccAFD	0.10	0.11	0.34	0.12	—	—	—	—
FehcpNM	0.17	0.10	0.11	0.19	0	0	0	0
CrbccAF	0	0	0	0	—	0	—	—
CrbccNM	0.03	0.02	0.02	0.04	0	0.01	0	0

 $M(\mu_B)$

	GGA				LDA				Exp. [16]
	NC/LCAO	PAW	US	NC	NC/LCAO	PAW	US	NC	
FebccFM									
Exp. (2.86)	2.22	2.19	2.25	2.23	2.16	2.16	2.19	2.16	2.22
Calc.	2.25	2.17	2.23	2.27	2.08	2.00	2.12	2.06	2.22
(a)	(2.87)	(2.83)	(2.85)	(2.87)	(2.79)	(2.75)	(2.77)	(2.78)	

	GGA				LDA				Exp. [31]
	NC/LCAO	PAW	US	NC	NC/LCAO	PAW	US	NC	
CrbccAF									
Exp. (2.88)	1.39	1.29	1.25	1.66	0.59	0.75	0	0.78	0.5
Calc.	1.39	1.20	1.30	1.70	0	0.44	0	0	0.5
(a)	(2.88)	(2.87)	(2.88)	(2.89)	(—)	(2.79)	(—)	(—)	

Testing the methods

Basis Effect

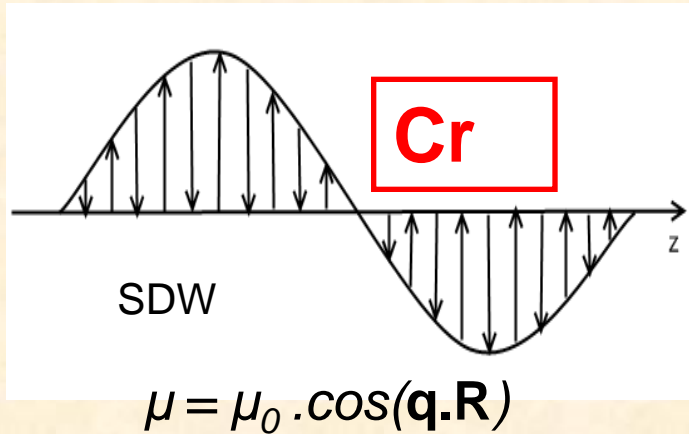
NC (PWscf)~NC(Siesta) if the localized basis is well optimized

Basis 1: DZ(2s), SZ(3p) SZ(5d)= 10 orbitals

Basis 2: DZ(2s), SZ(3p) SDZP(10d)= 15 orbitals

Minimal Basis seems accurate enough!

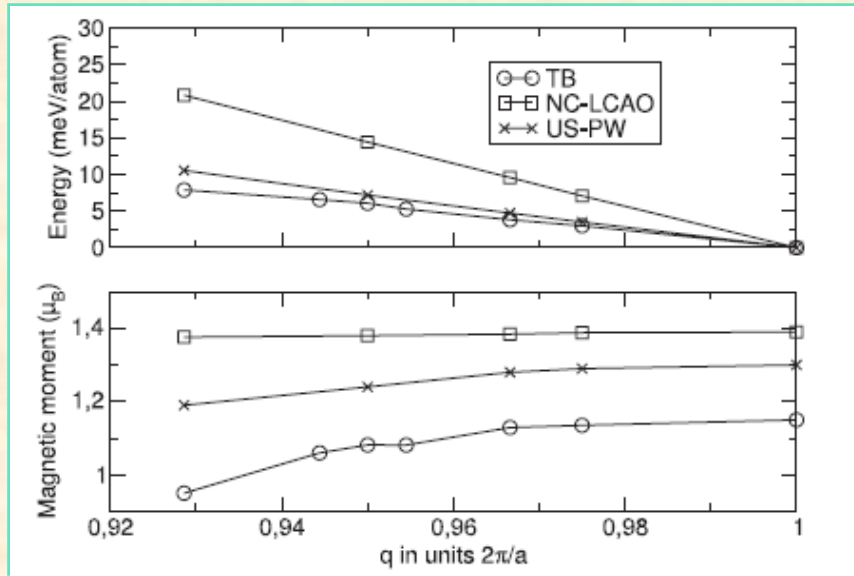
Ground state of Cr: SDW!!



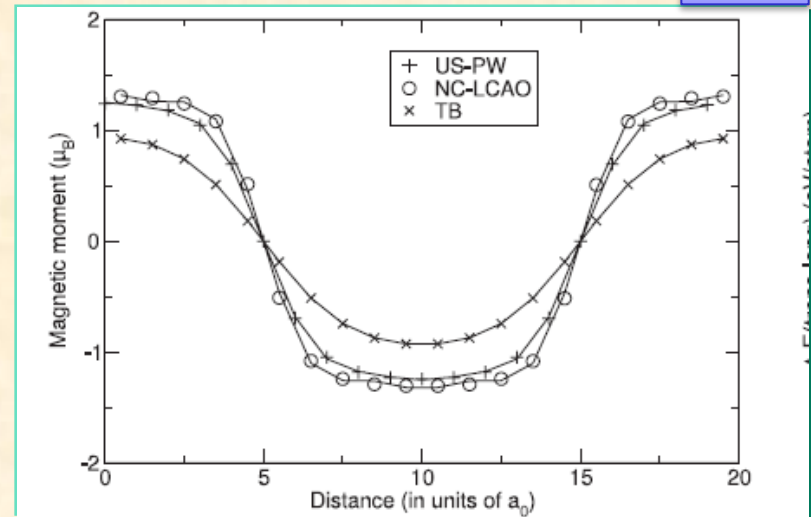
Experimental observation by neutron scattering (Corliss, 1959) : direction (001) et $q \approx 0,953$.

SDW is never stabilized

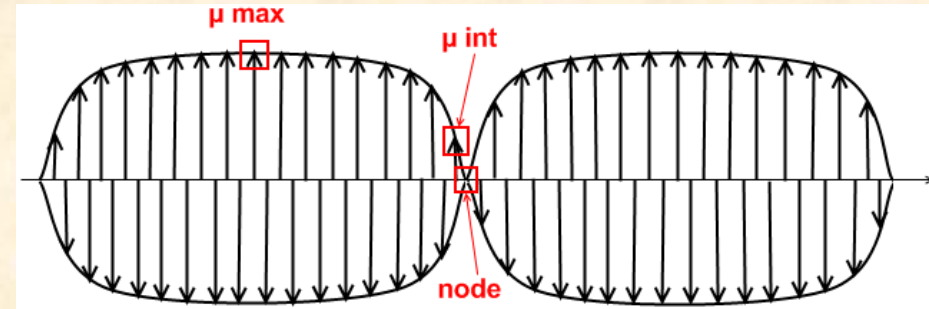
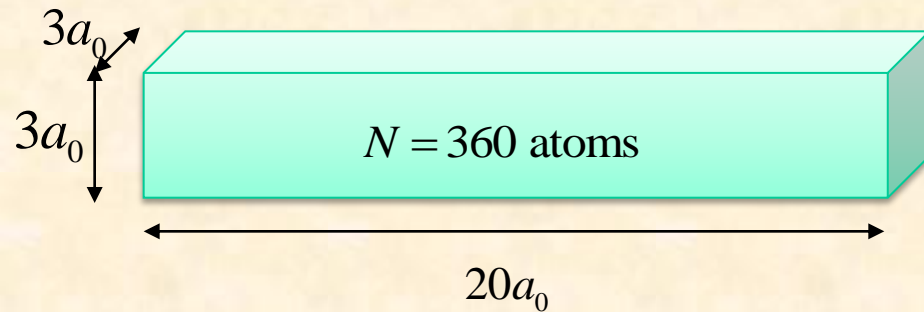
$\Delta E(q)$



$M(z)$



Vacancy formation energy in Cr SDW



$$E_V^f(\text{Cr}) = E(n-1)\text{Cr},V - \frac{(n-1)}{n} E(n\text{Cr})$$

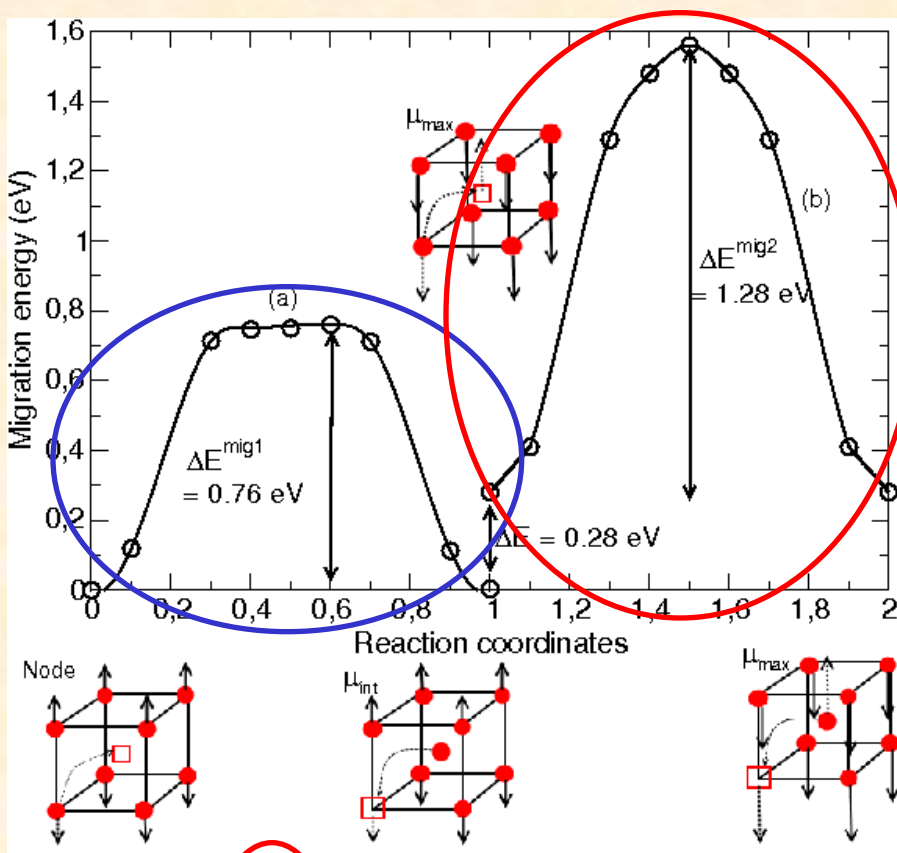
Position V	AF	NM	SDW-nœud	SDW- site inter.	SDW- site max.	Exp.*
E_V^f (eV)	2.41	2.32	2.09	2.10	2.38	2.0

* Landolt-Börnstein, PAS experiments (1985)

SIESTA

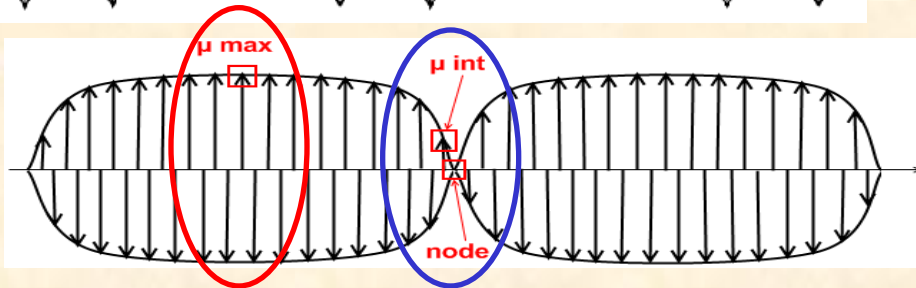
Vacancy formation is easier in a SDW node

Vacancy migration energy in Cr SDW



Cr	AF	NM	SDW-site μ_{max}	SDW-nœud
$E_{\text{Vac}}^{\text{mig}}$ (eV)	1.32	0.82	1.28	0.76

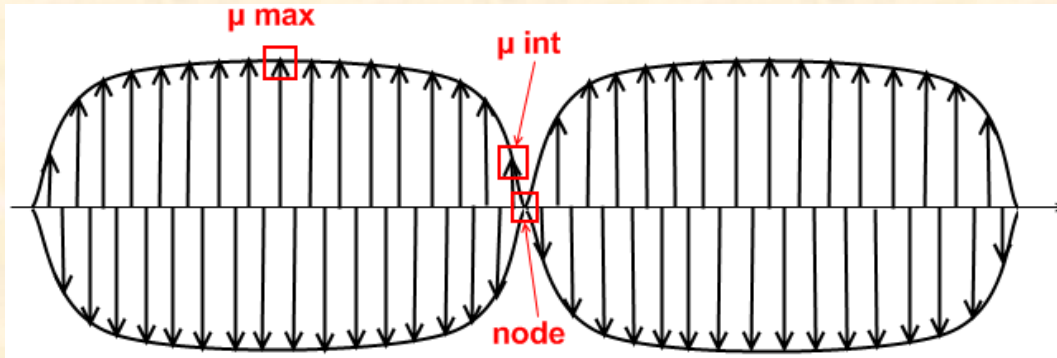
- Migration energy lower in SDW :
 $E_{\text{Vac}}^{\text{mig}}(\text{SDW}) < E_{\text{Vac}}^{\text{mig}}(\text{AF, NM})$
- Anisotropy of migration energy in SDW:
 $E_{\text{Vac}}^{\text{mig}}(\text{SDW node}) < E_{\text{Vac}}^{\text{mig}}(\text{SDW } \mu_{\text{max}})$



Soulaïrol, Fu and Barreateau, PRB 83, 214103 (2011)

Solution energy of Fe in Cr SDW

Strongly magnetic impurity: Fe



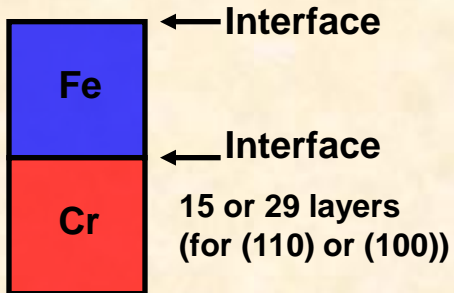
Position	AF	NM	SDW- nœud	SDW- site inter.	SDW- site max
$E_{\text{Fe}}^{\text{sol}}$ (eV)	0.38	0.29	0.27	0.28	0.39

SIESTA

- The solution energies are lower in the SDW.
- But $\Delta E(\text{AF-SDW})$ is rather low for Fe ($< \Delta E(\text{AF-SDW})$ for Cu)
 - Magnetic frustration of Fe in Cr : 2 possible spin states, $\mu_{\text{Fe}} = 0$ ou $2 \mu_{\text{B}}$. Multiples metastable solution when %Fe \uparrow .
- SDW experimentally stable for %Fe $< 1.6\%$ → understanding the destruction mechanism of the SDW

FeCr Interfaces

Interface energies



$E^f_{\text{interface}}$ (J/ m ²)	AF Cr	NM Cr	NCol.	Cr (100) SDW
Fe/Cr (100)	0.108	0.135	0.171	0.102
Fe/Cr (110)	0.189	0.091	0.174	0.115
Fe/Cr (111)	0.134	0.123	0.194	-

SIESTA

- Fe/Cr (100) interface is stabilized by magnetic effect contrary to the (110) interface

→ Magnetic frustration

Fe-Fe	Fe-Cr	Cr-Cr
↑ ↑	↑ ↓	↑ ↓

- Two possible ways of relaxing the magnetic frustration:

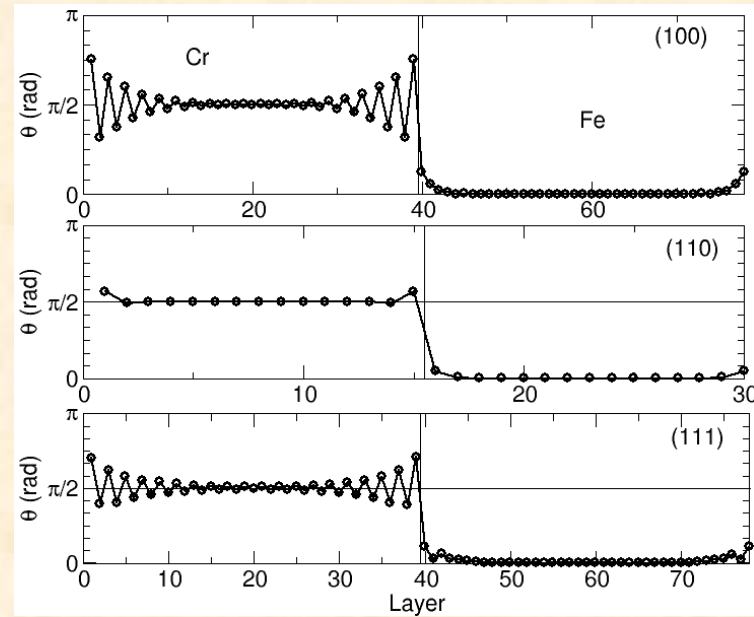
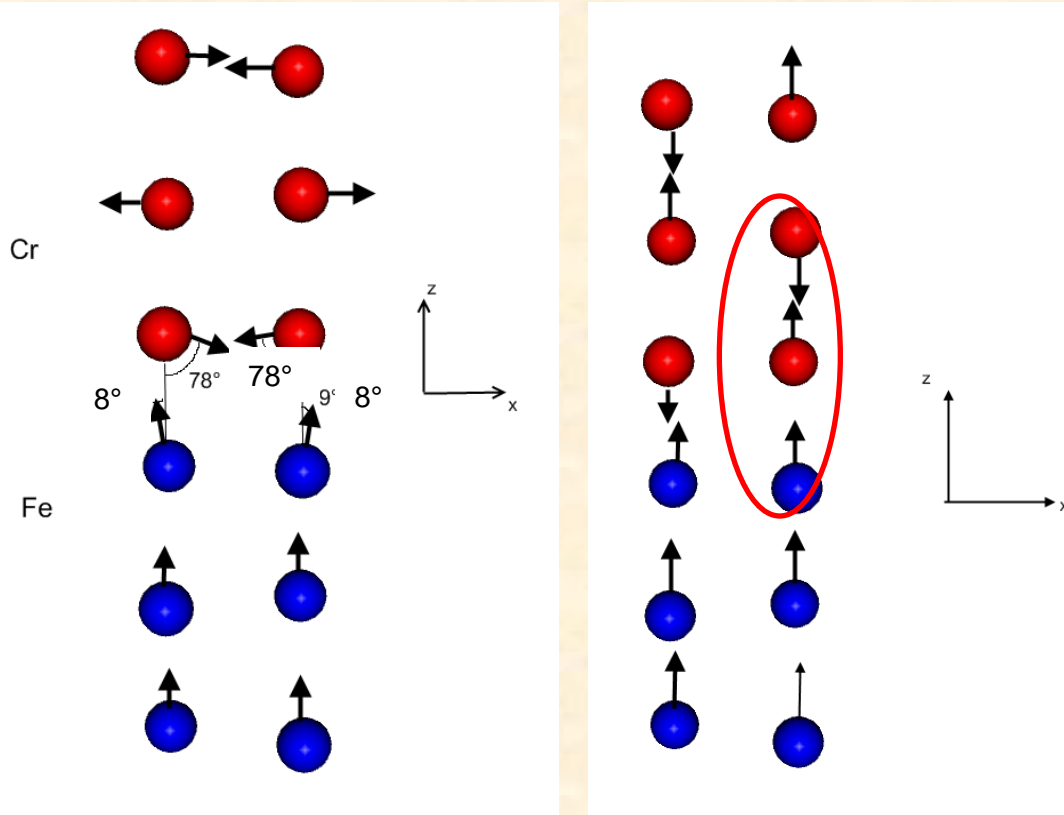
→ SDW with NM nodes near the interface (exp. : Bödeker *et al.* PRL, **81**, 914)

→ Non collinear configurations (exp. : Fritzsche *et al.* PRB, **65**, 144408)

FeCr Interfaces

SIESTA

Interface magnetic configurations



Non collinear

Collinear

E^f (NCol.)
0.174 J/m²

<

E^f (Col.)
0.189 J/m²

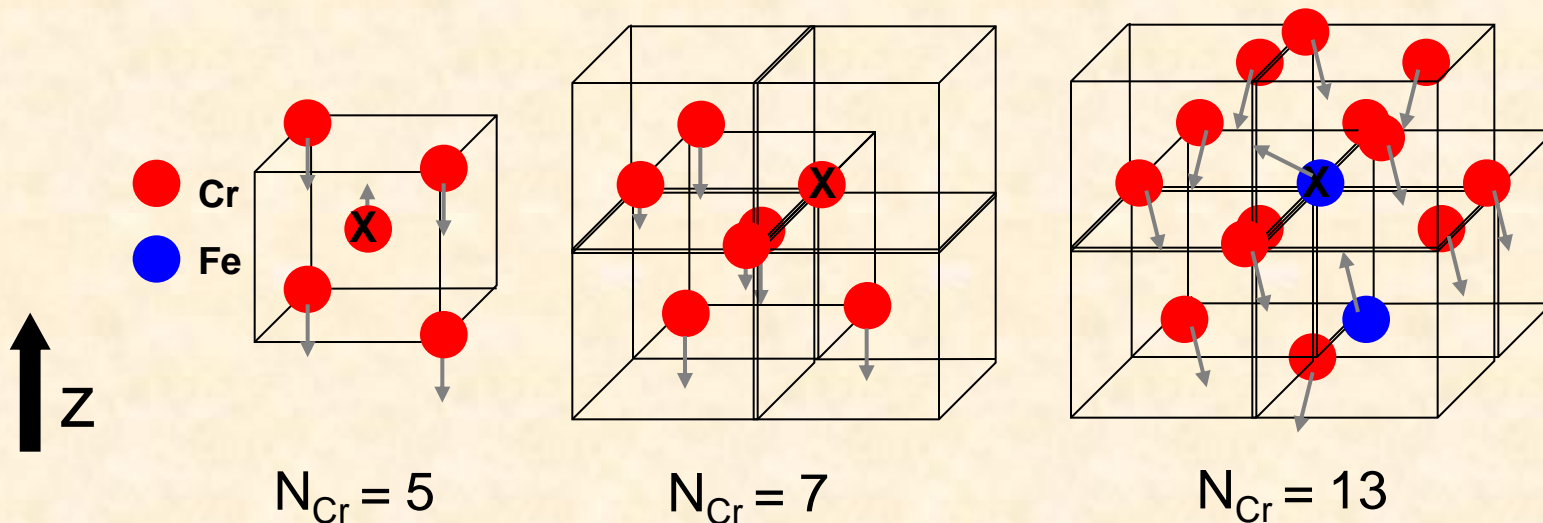
- Non collinearity lowers interface energy in Fe/Cr (110).
- Perpendicular magnetic coupling between Fe and Cr

Cr clusters in Fe matrix

Fe_xCr_y	$\text{Fe}_{123}\text{Cr}_5$	$\text{Fe}_{121}\text{Cr}_7$	$\text{Fe}_{115}\text{Cr}_{13}$
$\Delta E(\text{Col} - \text{NCol})$ (meV/Cr ou Fe)	0	0	7

SIESTA

- Collinear configurations for small Cr clusters in an Fe matrix ($N_{\text{Cr}} = 5$ and 7)
- Possible non collinear configurations of slightly lower energies for clusters of intermediate sizes ($N_{\text{Cr}} = 13$) [Longo *et al.* PRB 77, 212406 (2008) and Robles *et al.* PRB 74, 094403 (2006)]
- Possible non collinear configurations of slightly lower energies for clusters with (110) facets



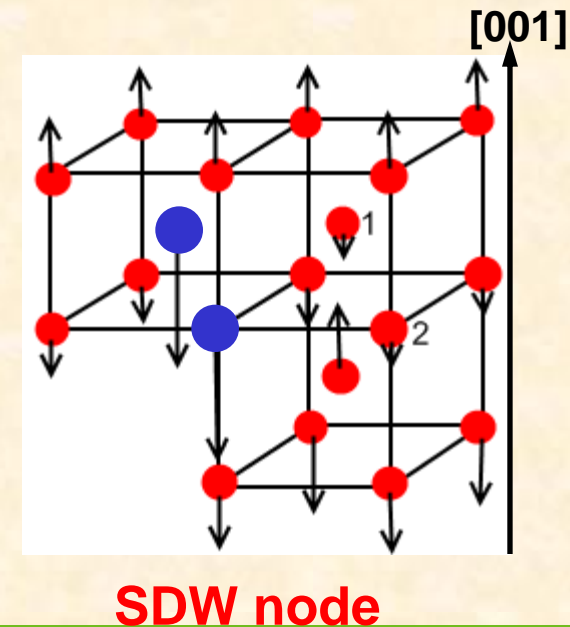
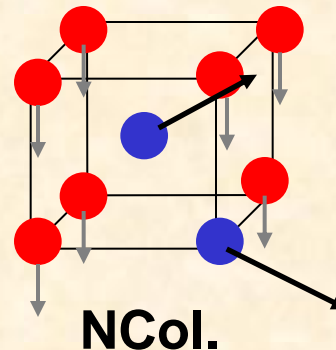
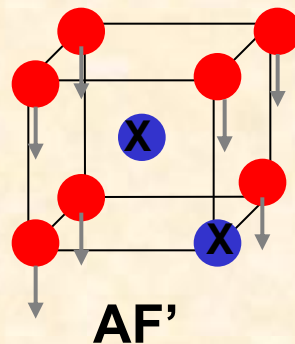
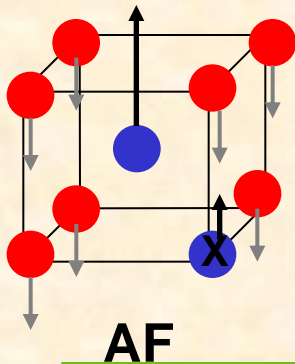
+ attraction
- répulsion

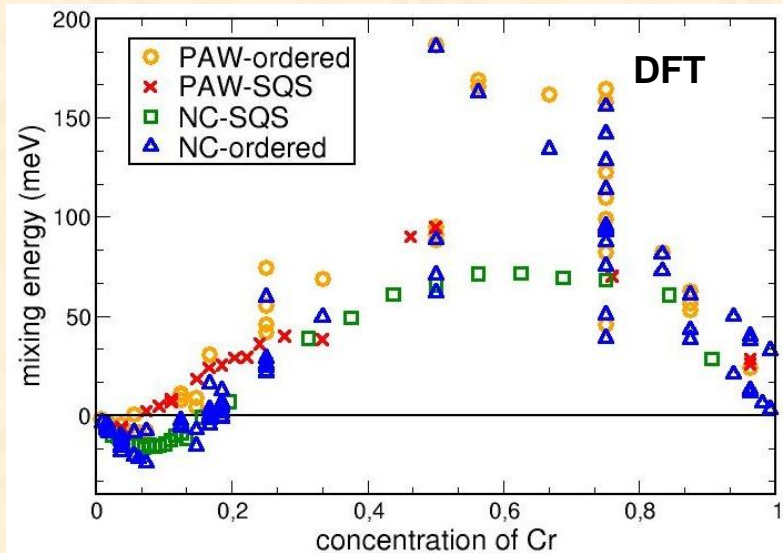
Fe-Fe interaction energy

SIESTA

Environnement Cr	AF	AF'	NCol.	SDW- node
$E_{\text{Fe-Fe}}^b$ (eV)	- 0.01	- 0.04	+ 0.07	+ 0.12
μ_{Fe1} (μ_B)	2.26	0.07	1.89	2.06
μ_{Fe2} (μ_B)	0.67	0.06	1.88	2.04

- Multiple magnetic states for Fe dimer in Cr.
- Precipitation of Fe is favored in Cr SDW.
- Configurations non colinéaires de faible énergie pour les clusters avec facettes (110).





Mixing energy calculated at P=0

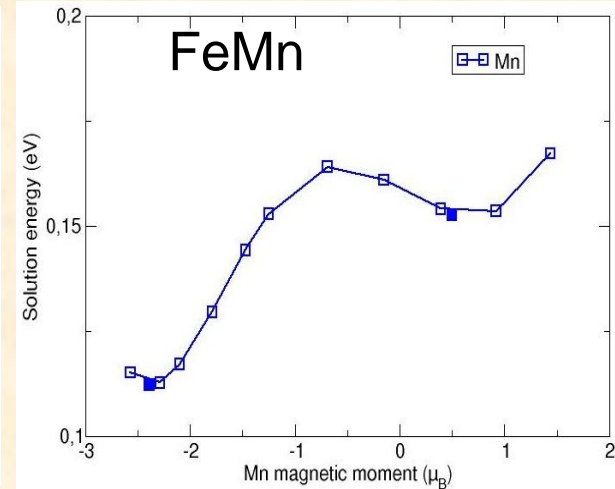
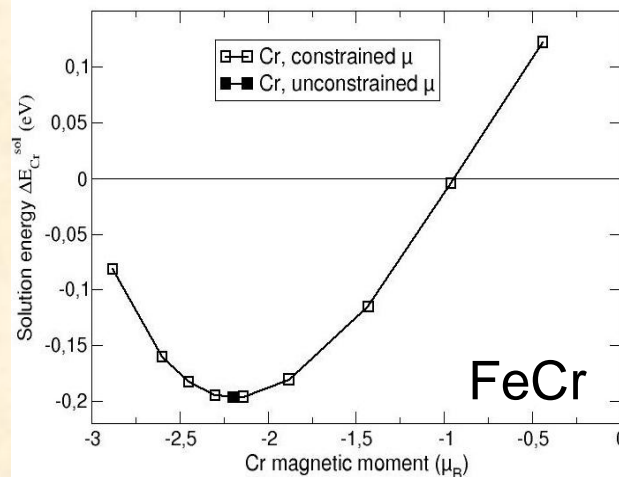
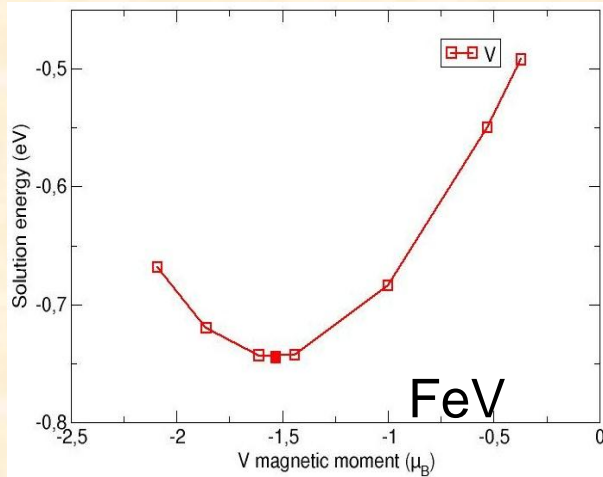
(Cr exp. [Mirebeau *et al.*, PRL 53, 687 (1984)])

Some trends in the periodic table

5	6	7	8	9	10
VB	VIB		VIIIB		
23 V Vanadium 50.9415	24 Cr Chrome 51.9961	25 Mn Manganèse 54.938049	26 Fe Fer 55.8457	27 Co Cobalt 58.933200	28 Ni Nickel 58.6934
41 Nb Niobium 92.90638	42 Mo Molybdène 95.94	43 Tc Technétium (98)	44 Ru Ruthénium 101.07	45 Rh Rhodium 102.90550	46 Pd Palladium 106.42

	FeV	FeCr	FeMn
ΔE^{sol}	-	- puis +	+
$N_{\text{Fe}}^{\text{d}} - N_{\text{X}}^{\text{d}}$	3	2	1
$\Delta\mu = \mu_{\text{sol}} - \mu_{\text{bulk}} (\mu_{\text{B}})$	1,53	1.00	0,85

$$\Delta E_{sol} = E(n-1)Fe, Cr - (n-1)E(Fe) - E(Cr) \quad V, Cr, Mn = 1.85\% at$$



PWscf

- Effet of d band filling on the solubility of V, Cr et Mn in Fe
- AF interaction between Fe and Cr, Mn or V favors AF solutions

FeCr: intermédiaire case : magnétisme is the driving force

$\mu_{Cr} < -0.8\mu_B \rightarrow$ mixing

$\mu_{Cr} > -0.8\mu_B \rightarrow$ demixing

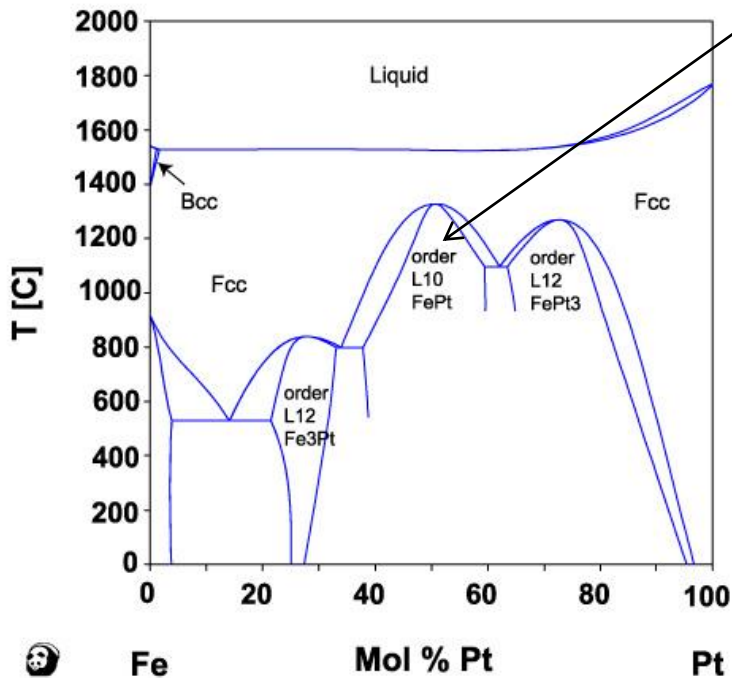
Partial Conclusion

- Fe-Cr is a particularly complex system where magnetism plays a crucial role in the energetics.

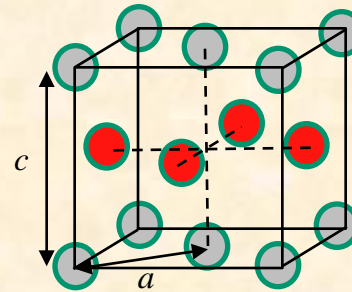
Questions and Comments

- Can we stabilize the SDW (Fermi Surface Nesting?)
- How could we introduce (spin and ion) temperature effects?
- Need for simpler models?

FePt phase diagram



L10 structure



$$a_{\text{exp}} = \frac{3.86}{\sqrt{2}} = 2.73 \text{ \AA}$$

$$c_{\text{exp}} = 3.72 \text{ \AA}$$

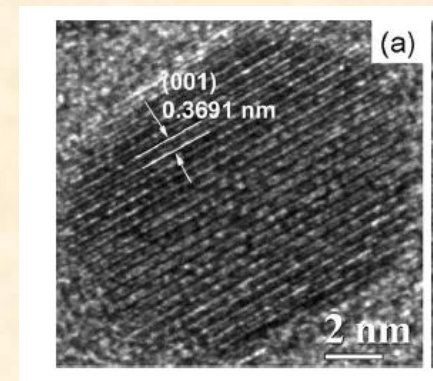
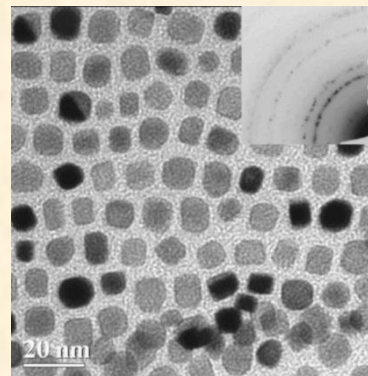
$$\frac{c_{\text{exp}}}{a_{\text{exp}}} = 1.36$$

$$V_{\text{exp}} = 27.7 \text{ \AA}^3$$

very high magnetic uniaxial anisotropy

MAE=1.4meV/atom (exp.)

Good control of nanocrystal growth



Magnetic TB model

$$H = H_0 + H_{mag} + H_{LCN} + H_{SOC}$$

- H_0 : spd Tight-Binding (non magnetic) Hamiltonian
- H_{mag} : Stoner Hamiltonian

$$H_{mag} = -\frac{1}{2} \sum_{i\lambda} I_{i\lambda} \vec{m}_{i\lambda} \cdot \vec{\sigma}$$

- H_{LCN} : local charge neutrality constraint

$$H_{LCN} = \sum_{i\lambda} U_{LCN} (n_i - n_i^0) |i\lambda\rangle \langle i\lambda| + \sum_{i\lambda \neq d} U_d (n_{i,d} - n_{i,d}^0) |i\lambda\rangle \langle i\lambda|$$

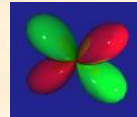
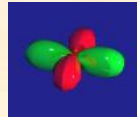
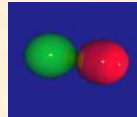
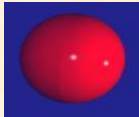
- H_{SOC} : Spin Orbit Coupling

$$H_{SOC} = \sum_i \xi_i (r - R_i) \vec{L}_i \cdot \vec{S} \quad \xi_{d,i} = \int_0^\infty R_{d,i}^2(r) r^2 dr$$

Determination of parameters

$$H_0 = \sum_{ij\lambda\mu} |i\lambda\rangle \langle i\lambda| H |j\mu\rangle \langle j\mu|$$

i=atom
λ=orbital



λ = s

p_x

p_y

p_z

d_{xy}

d_{xz}

d_{yz}

d_{x²-y²}

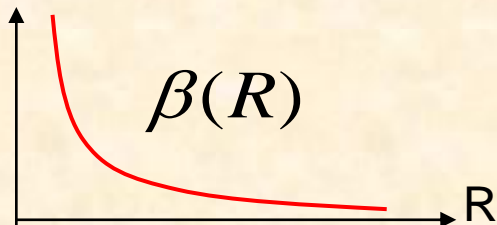
d_{3z²-r²}

Hopping integral

Onsite term

$$\beta_{ij}^{\lambda\mu} = \langle i\lambda | H | j\mu \rangle$$

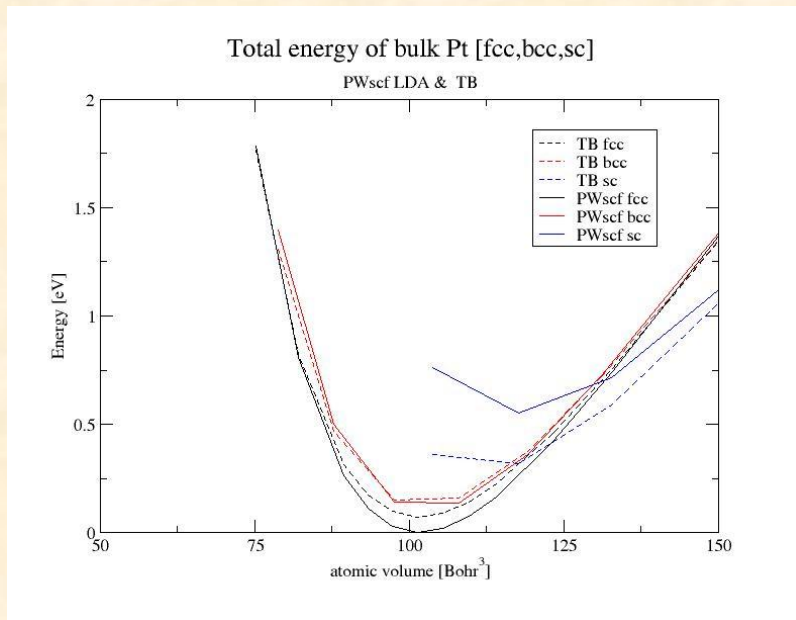
$$\varepsilon_{i\lambda} = \langle i\lambda | H | i\lambda \rangle$$



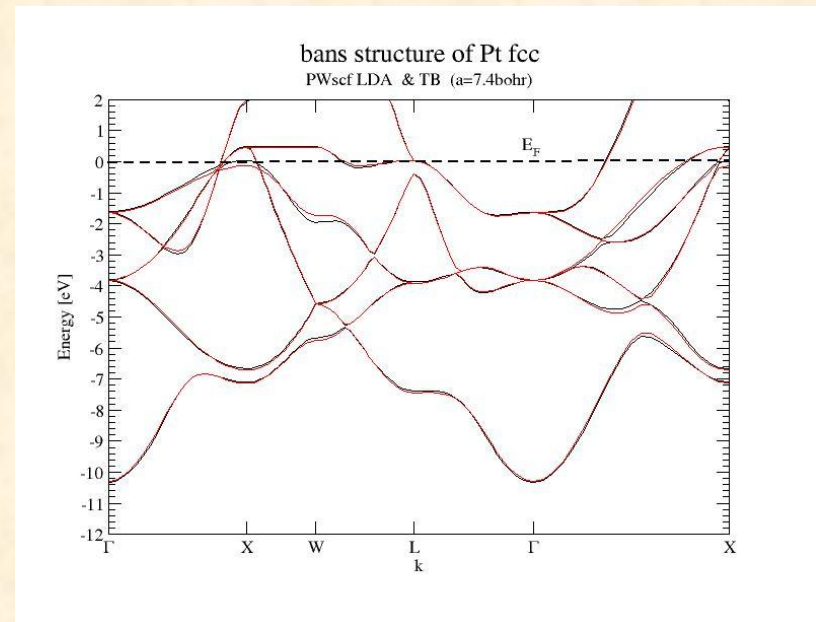
Determination of parameters

- H_0 : Hopping integrals and onsite elements obtained from simultaneous fit of ab-initio **band structure** and **total energy** curves of **bulk non magnetic Fe and Pt**

Total energy



Band structure



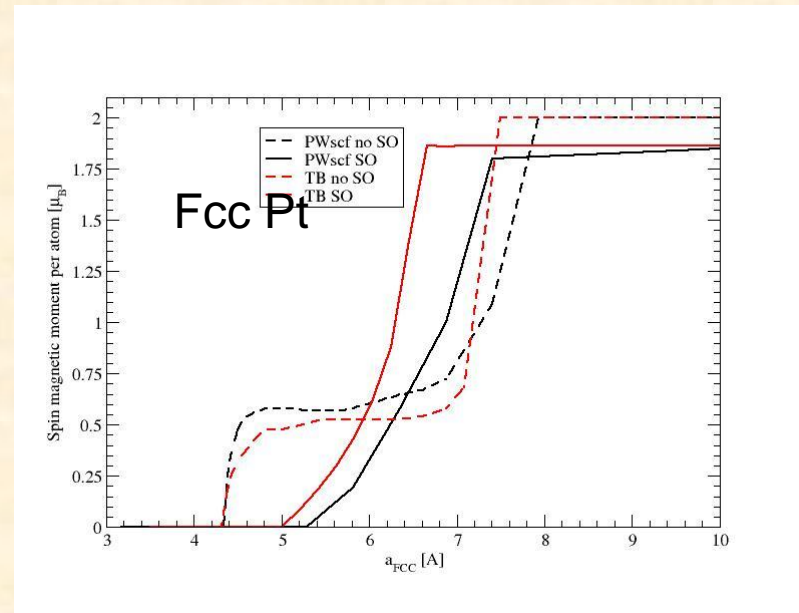
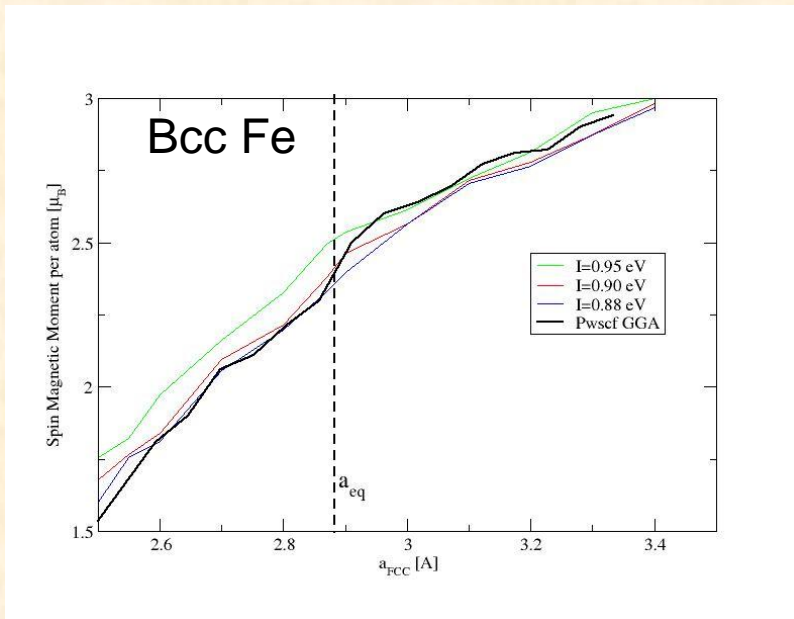
Determination of parameters

- H_{mag} : Stoner parameter I adjusted to reproduce ab-initio $M(d)$ of bulk Fe and Pt

$$H_{mag} = -\frac{1}{2} \sum_{i\lambda} I_{i\lambda} \vec{m}_{i\lambda} \cdot \vec{\sigma}$$

$$I_{Fe} \in [0.88, 0.95] eV$$

$$I_{Pt} = 0.60 eV$$

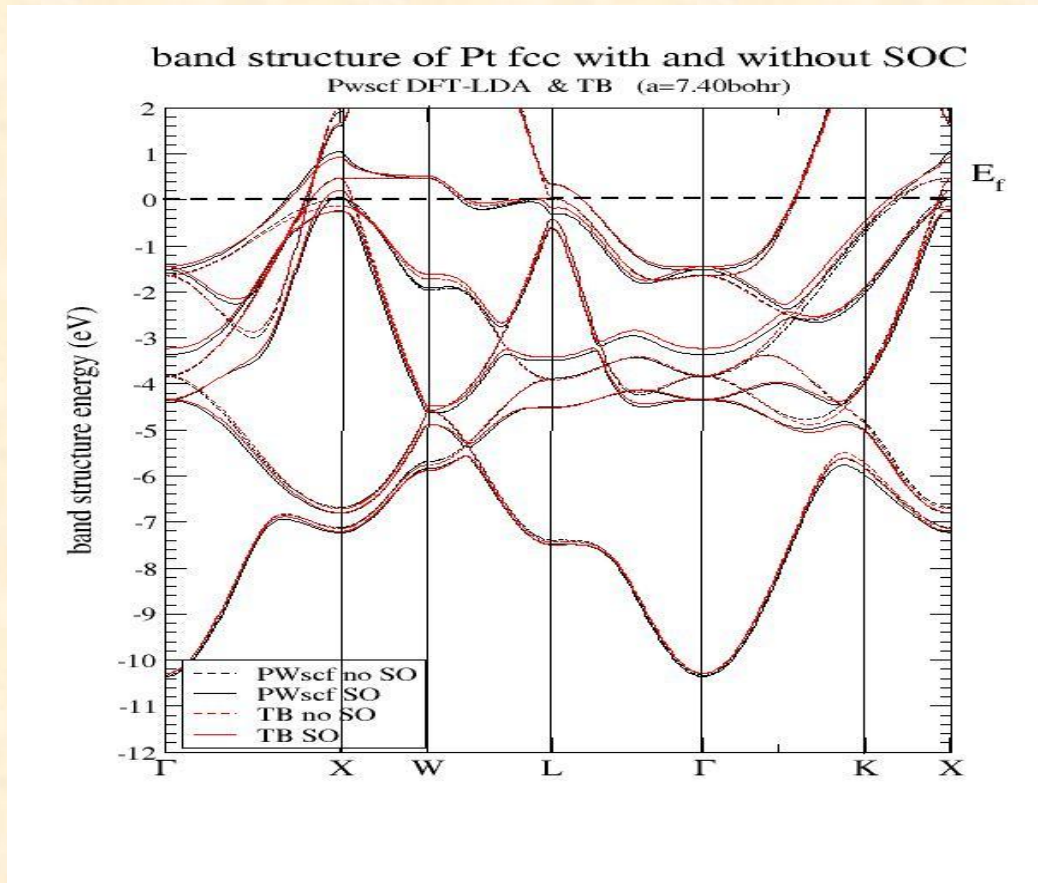


Determination of parameters

- H_{SOC} : Spin Orbit Coupling adjusted to reproduce ab-initio band structure

$$H_{\text{SOC}} = \sum_i \xi_i (r - R_i) \vec{L}_i \cdot \vec{S}$$

$$\xi_{d,i} = \int_0^{\infty} R_{d,i}^2(r) r^2 dr$$



$$\xi_{\text{Fe}} = 0.06 \text{ eV}$$

$$\xi_{\text{Pt}} = 0.57 \text{ eV}$$

Determination of parameters

- H_{LCN} : local charge neutrality

$$H_{LCN} = \sum_{i\lambda} U_{LCN} (n_i - n_i^0) |i\lambda\rangle \langle i\lambda| + \sum_{i\lambda \neq d} U_d (n_{i,d} - n_{i,d}^0) |i\lambda\rangle \langle i\lambda|$$

↑
Charge neutrality

↑
“d” orbital filling

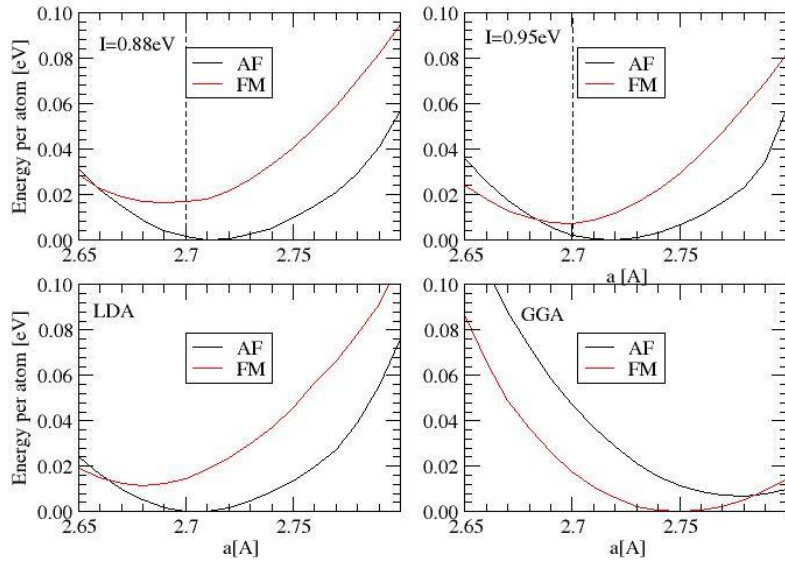
$$U = U_d = 20eV$$

$n_{i,d}^0$ adjusted to reproduce electronic and magnetic properties of FePt L10

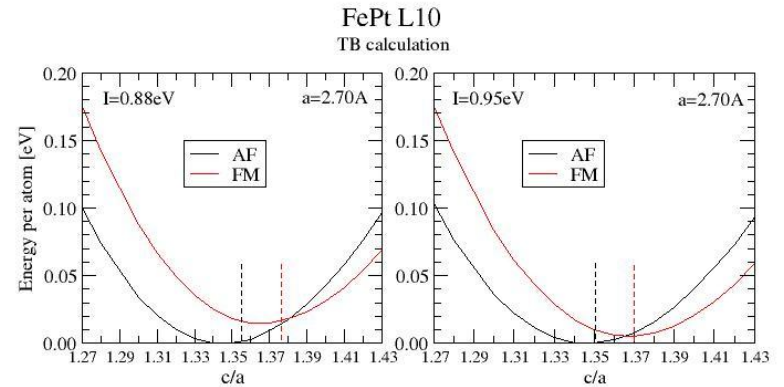
$$n_{Fe,d}^0 = 6.6 \quad n_{Pt,d}^0 = 8.8 \quad \longrightarrow \quad M_{Fe} \sim 3\mu_B \quad M_{Pt} \sim 0.35\mu_B$$

FM vs AF

Functionnal effect

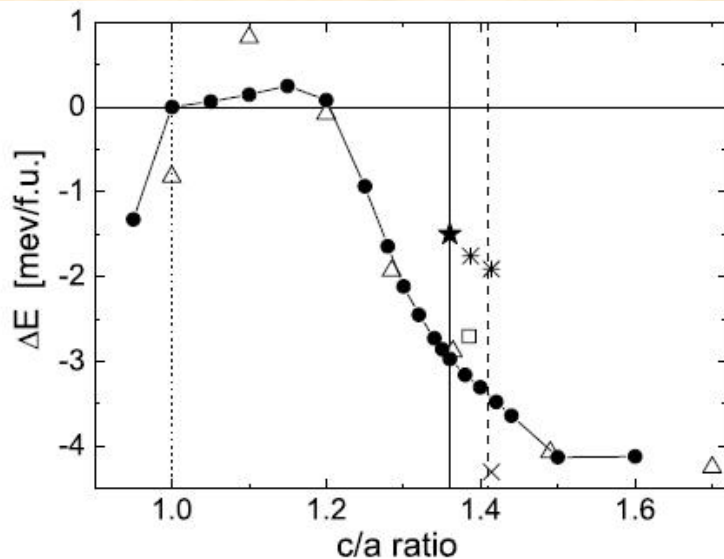


Structural effect

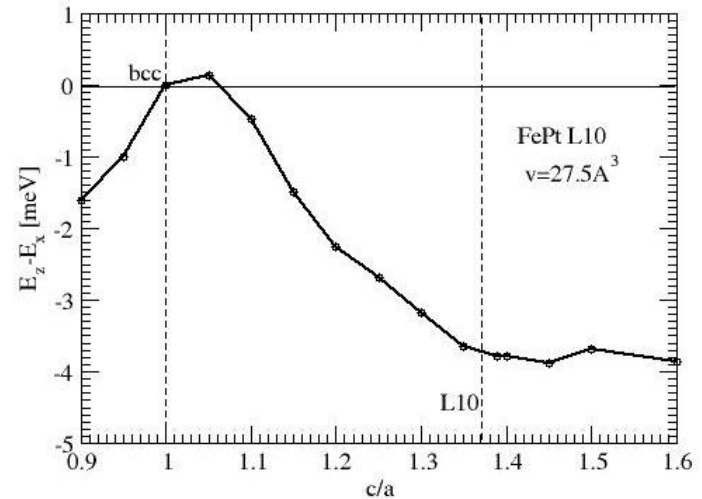


Magnetic Anisotropy Energy

J.Phys.: Condens. Matter 17 (2005) 4157



Tight-Binding



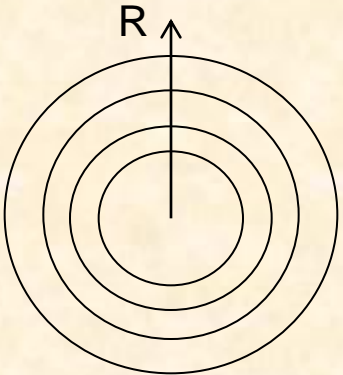
Remark: MAE too large!

LSDA+U+SOC

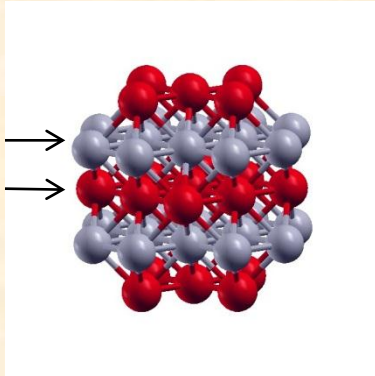
Disorder?

FePt L10 clusters

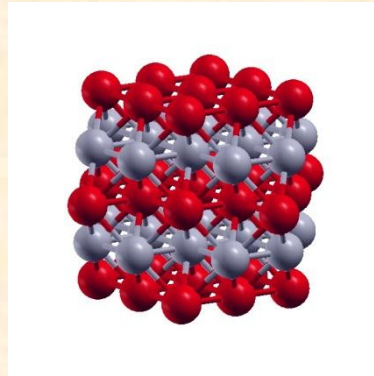
Clusters of increasing size



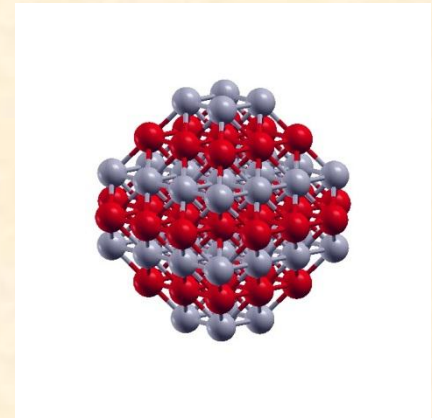
Pt →
Fe →



$N = 43$



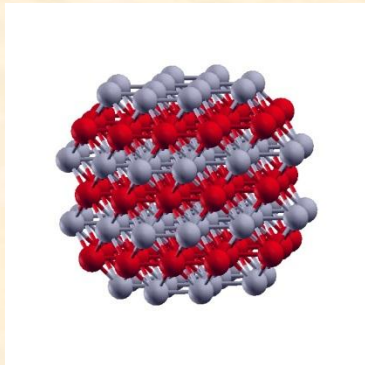
$N = 55$
cuboctahedron



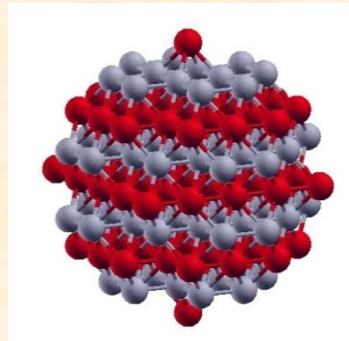
$N = 79$

R N

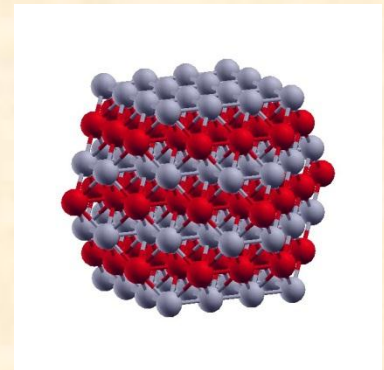
- 147
- 141
- 135
- 87
- 79
- 55
- 43
- 19
- 13
- 1



$N = 135$



$N = 141$

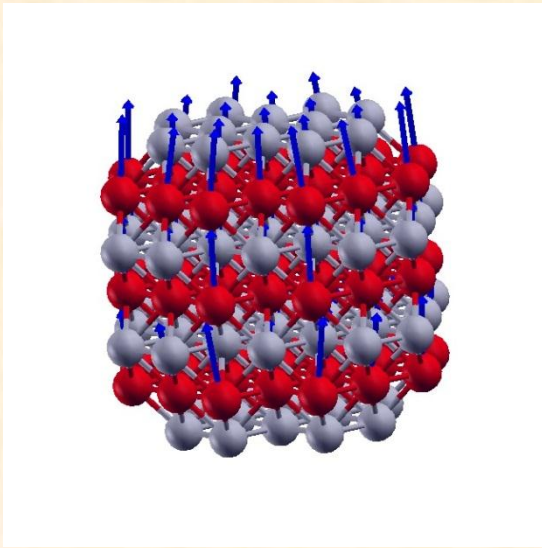


$N = 147$
cuboctahedron

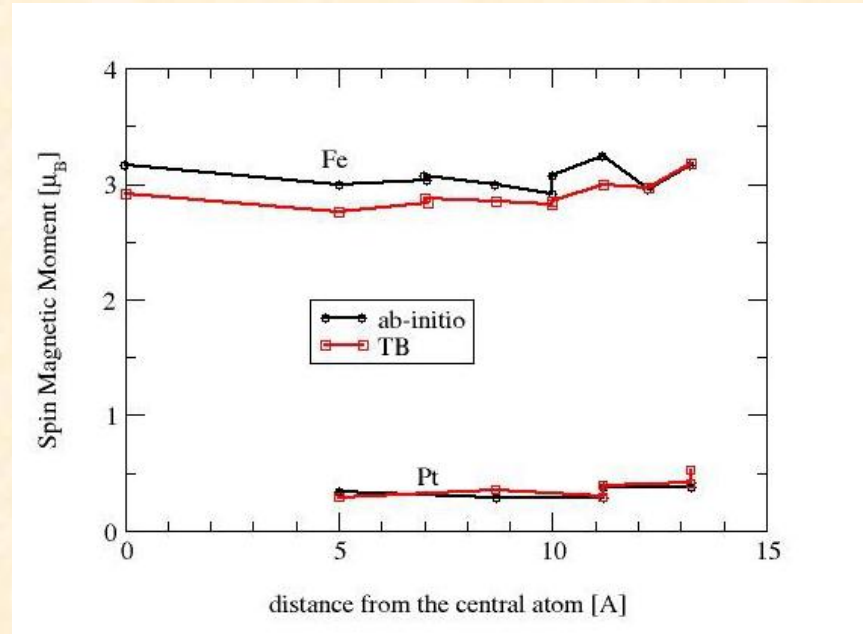
Not a spherical shell

Repartition of spin magnetic moment in the cluster

FM // z

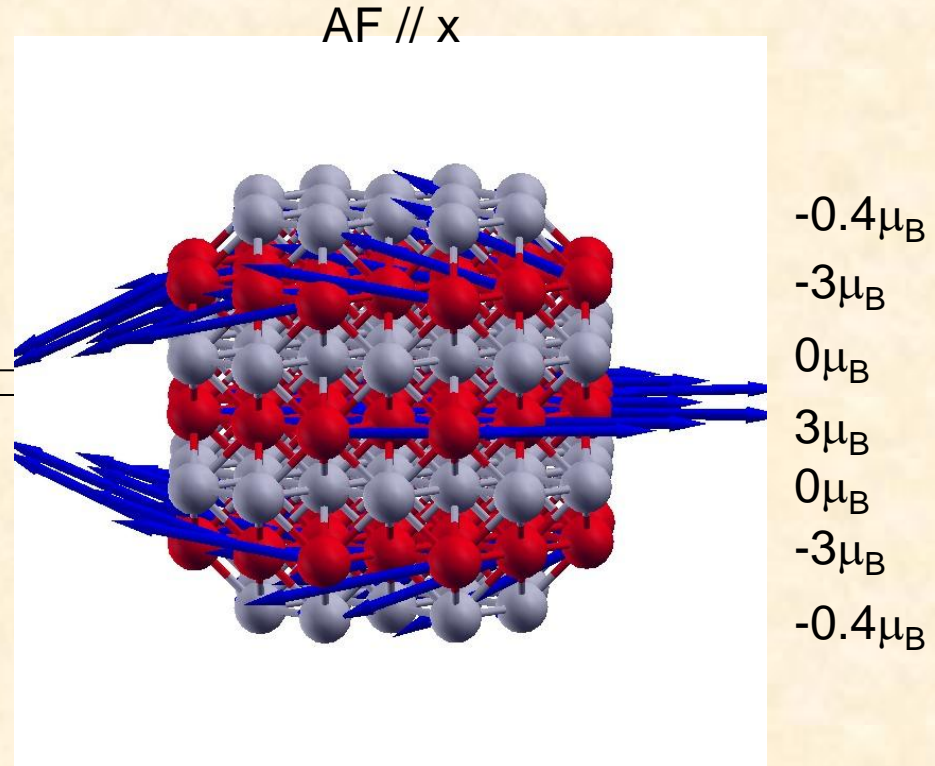
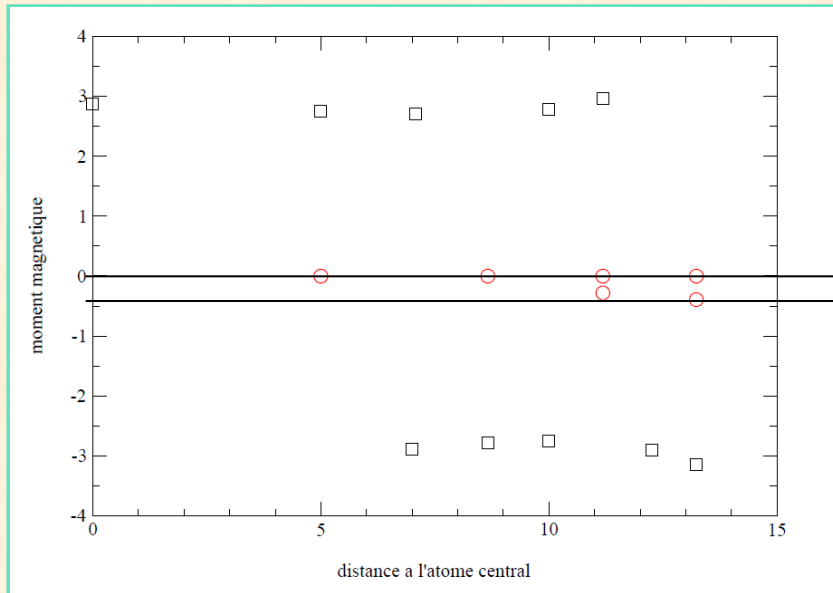


$N = 135$



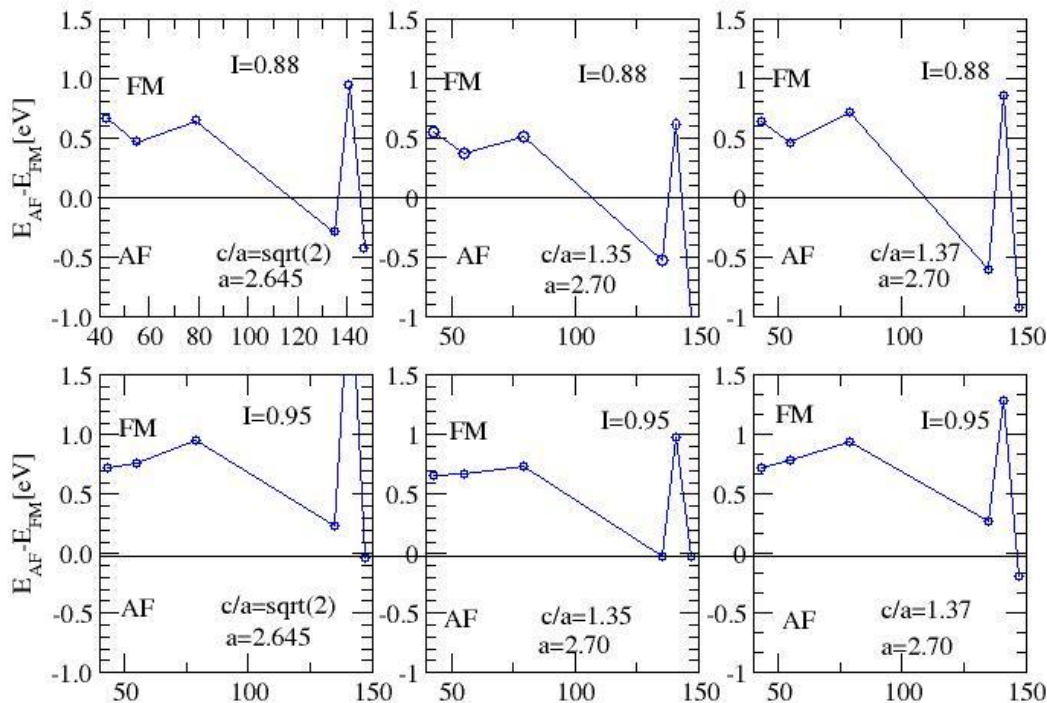
Ab-initio= Comput. Matt. Sci. 35 (2006) 279

Repartition of spin magnetic moment in the cluster

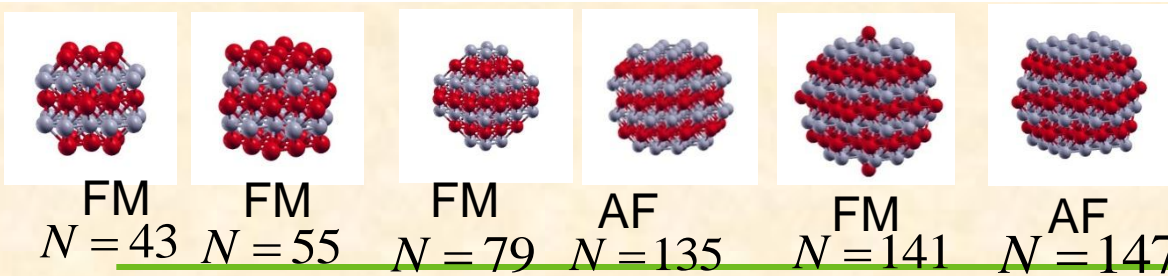


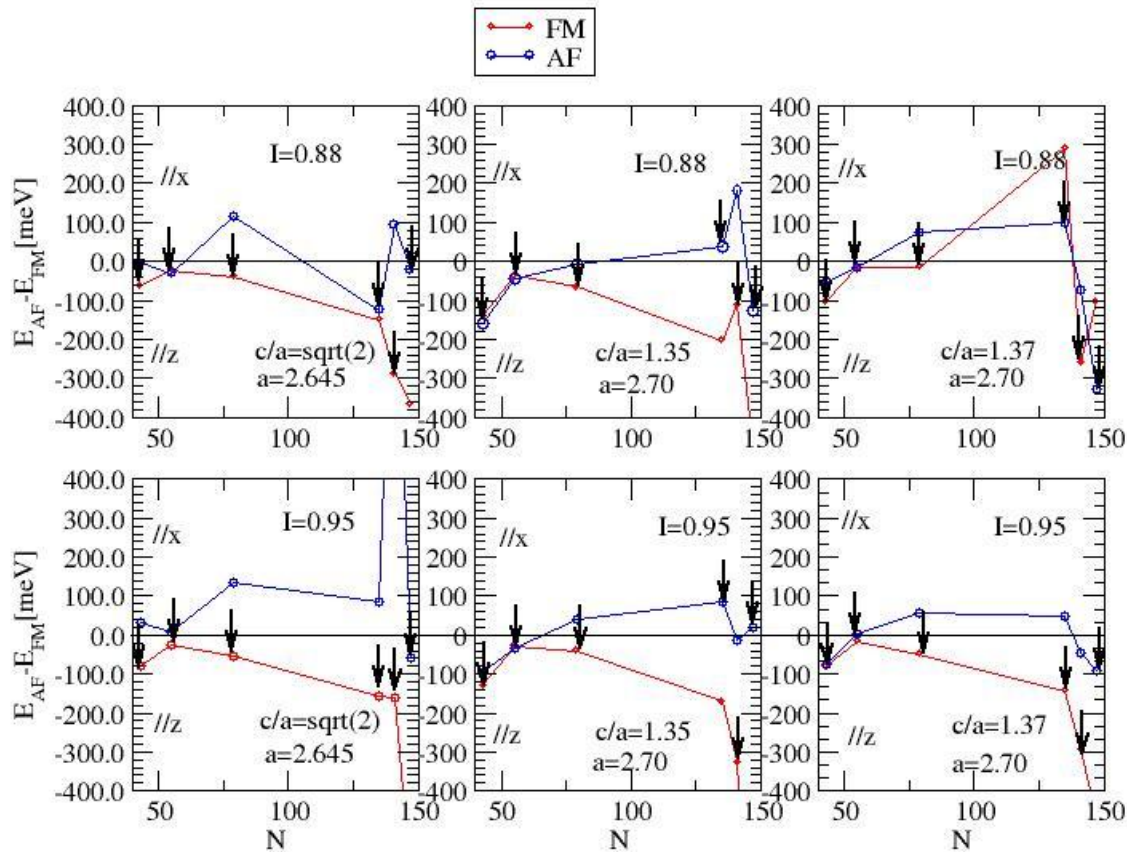
AF vs FM order in FePt clusters

FePt L10 clusters: FM vs AF



Pt termination favors AF
 Fe termination favors FM
 Large Stoner parameter favors FM
 Large c/a favors FM





Easy axis along z
 Except for N=135,147
 (AF ordering)

Partial Conclusion

- Efficient and quantitative TB method for electronic and magnetic properties of metals and their alloys.
- Complex magnetic behavior of FePt clusters: FM vs AF, oscillating MAE etc..
 - influence of surface termination
 - influence of c/a

Questions and Comments

Why is there no experimental evidence of AF order?

LSDA+U

Disorder?

Strain effect?

THANK YOU FOR YOUR ATTENTION